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IMPROVED ANALYTICAL SHAPED CHARGE CODE: BASC

John T. Harrison

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US ARMY ARMAMENT RESEARCH AND DEVELOPMENT COMMAND  
BALLISTIC RESEARCH LABORATORY  
ABERDEEN PROVING GROUND, MARYLAND

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20. ABSTRACT (Continue on reverse side if necessary and identify by block number) (ddg) The work of many researchers has been combined to produce a simplified, analytical, computer code to address the shaped charge problem. This code named BASC (BRL Analytical Shaped Charge) is programmed in the FORTRAN language. BASC utilizes a modified version of the metal acceleration model of M. Defourneaux of France to account for both the final liner velocity and its acceleration history. The BASC code accurately predicts jet-tip velocity and treats the buildup of the massive lead pellet, for both heavy confined and unconfined charge geometries. Calculation of the massive lead pellet has traditionally		

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been difficult and is accomplished in BASC by including time dependent acceleration and material compressibility in the region of the liner near the apex of the conical liner. The BASC code includes the jet formation theory of Pugh, Eichelberger, and Rostoker; the shaped-charge penetration theory of DePersio, Simon and Merendino for penetration-standoff curves; and the piece wise penetration of Defourneaux for hole profiles. BASC enables parametric investigations of shaped charge problems with relatively small amounts of computer time since the code is basically analytic. Results from the BASC code are compared to experiments as well as to more sophisticated hydrodynamic computer codes. The report documents BASC, including the equations utilized, necessary empirical relations, FORTRAN listing of the program, selected problem examples and comparison with experiments.

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## I. INTRODUCTION

The Ballistic Research Laboratory (BRL) of the U.S. Army Armament Research and Development Command (ARRADCOM) has a wide interest in the shaped-charge problems, ranging from detailed studies of the flow characteristics of the collapsing liner, to designing practical devices for future warhead applications. For these efforts several experimental and theoretical techniques are employed. Often, it is necessary to determine parametric relationships in the design application. Since it is more feasible to utilize theoretical calculations that are economically sound for parametric studies, the BRL has several finite-difference, hydrodynamic, computer codes that have been applied to shaped-charge problems.<sup>1,2</sup> Although these codes are adaptable to various geometrical considerations, they require operator experience and a seasoned analyst to insure proper application. Further, large, high-speed computers, and long calculational times are necessary. Quite often, it is desirable to have a simplified procedure for addressing parametric design studies quickly and economically. The BRL computer code named BASC (BRL Analytical Shaped Charge)<sup>3</sup> was formulated from analytic expressions to provide such capability. Although several advantages occur with the original BASC approach, several areas of difficulty were experienced, particularly those relating to accurate calculation of jet tip or lead pellet behavior and confined charges. Extensive semi-empirical functions, regarding liner acceleration and confinement effects, have been included to provide more accurate representations. This improved, simplified procedure, hereinafter referred to as the BASC code also, together with additional refinements are presented and discussed in this report.

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<sup>1</sup>J. T. Harrison and R. R. Karpp, "Terminal Ballistic Applications of Hydrodynamic Computer Code Calculations," BRLR 1964, April 1977. (AD #A041065)

<sup>2</sup>J. T. Harrison, "A Comparison Between the Eulerian, Hydrodynamic Computer Code (BRLSC) and Experimental Collapse of a Shaped Charge Liner," ARBRL-MR-02841, June 1978 (AD A059711)

<sup>3</sup>J. Harrison, R. DiPersio, R. Karpp and R. Jameson, "A Simplified Shaped Charge Computer Code: BASC," DEA-AF-F/G-7304 Technical Meeting: Physics of Explosives, Vol II, April-May 1974, Paper 13 presented at the Naval Ordnance Laboratory, Silver Spring, MD.

The BASC code is an assembly of various theoretical and empirical techniques. Central to the procedure utilized in BASC is the Defourneaux model<sup>4</sup> for final plate velocity resulting from the shock of an adjacent detonating explosive. This assumption is adequate for portions of the liner which are initially removed (remote) from the collapse (jet formation) region or cone axis. In actuality, several shock reverberations are required to achieve a final liner metal velocity. Material near to the apex of the cone can enter the collapse process long before the liner is accelerated fully and, hence, does not achieve its ultimate velocity. This leads to the well known phenomena referred to as "the inverse velocity gradient" which forms the massive jet tip. This phenomena has been observed<sup>5</sup> and calculated earlier.<sup>1,2,6</sup> The author has modified the Defourneaux model to account for the time-dependent acceleration of the liner resulting in a gradual build-up to the ultimate collapse velocity. Since liner elements near the apex region of a cone will not achieve this ultimate velocity, the first element of the jet will move more slowly than the following elements of the jet. The faster elements collide and are compressed into the massive lead pellet. The final jet-tip velocity will become the mass-weighted average of these inverse velocity elements. This is the jet-tip velocity, which will be used in jet penetration theory. BASC uses a combination of the shaped-charge penetration of DiPersio, Simon and Merendino (DSM)<sup>7</sup> and the piece wise penetration of Defourneaux<sup>4</sup>. The DSM model is used to calculate total penetration-standoff curves and the piecewise penetration model is used to calculate whole profiles.

In addition to the inverse velocity gradient, flow into the stagnation region during the collapse process from particles near the apex of a conical liner may be supersonic and fail to form a jet or form a so-called incoherent jet. This is the jet-no-jet criteria. The criteria used in BASC is a modified version of the supersonic limitations of Chou, et. al.<sup>8</sup> The final theoretical technique used in BASC is the

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<sup>4</sup>M. Defourneaux, "Hydrodynamic Theory of Shaped Charges and of Jet Penetration," *Memorial DeL'art Ille'rie Francasise-T*, 44, 1970.

<sup>5</sup>R. DiPersio, C. W. Whiteford, and J. Simon, "An Experimental Method of Obtaining Collapse Velocities of the Liner Walls of a Linear Shaped-Charge Liner," BRL-MR-1696, September 1965. (AD #478326)

<sup>6</sup>A. Kiwan and H. Wisniewski, "Theory and Computations of Collapse and Jet Velocities of Metallic Shaped Charge Liners," BRLR-1620, November 1972. (AD #907161)

<sup>7</sup>R. DiPersio, J. Simon, and A. Merendino, "Penetration of Shaped-Charge Jets Into Metallic Targets," BRLR 1296, September 1965. (AD #476717)

<sup>8</sup>P. Chou, J. Carleone, R. Karpp, "Criteria for Jet Formation from Impinging Shells and Plates," *J. Appl. Physics*, Vol. 47, No. 7, July 1976.

Pugh, Eichelberger, and Rostoker (PER) theory of jet formation.<sup>9</sup>

This report includes a description of analytical equations, the procedures for their application and a FORTRAN listing of BASC. Experimental results for selected shaped charge problems are presented in a comparison to the calculations. The range of useful applications is discussed as well as the limitations of the approach.

BASC is a "living code" which has provided insights to the jet-formation process and is a very good tool for parametric design for a selected class of problems.

A simplified flow chart and the FORTRAN code listing is presented in Appendices A and B, respectively.

## II. GOVERNING EQUATIONS

The initial equation used in the liner acceleration portion of the BASC code determines the angle of liner bending,  $\phi$ , produced by a detonation wave traveling with a velocity,  $D$ , and inclined to the liner wall at an angle,  $i$  (the angle of incidence). This relationship is illustrated in Figure 1. The author has modified the original platepush relationship of Defourneaux to be

$$\frac{1}{\phi} = \frac{1}{\phi_0} + K \frac{\rho \epsilon}{e B}, \quad (1)$$

$$B = 1 + A/\rho_c e_c,$$

where  $\rho$  and  $\epsilon$  are the density and thickness of the liner wall, and  $e$  is the explosive thickness. Added are  $\rho_c$ ,  $e_c$ , and  $A$  which are, respectively, density, thickness, and a constant, which is determined from the experimental data for the confinement casing around the charge<sup>10</sup>. The constant,  $A$ , when set to zero, represents an unconfined explosive charge and the Defourneaux relationship, as illustrated in Figure 1.  $\phi_0$  and  $K$  are functions of the angle of incidence,  $i$ , and

<sup>9</sup>E. M. Pugh, R. J. Eichelberger and N. Rostoker, "Theory of Jet Formation by Charges with Lined Conical Cavities," *J. Appl. Physics*, Vol. 23, No. 5, May 1952.

<sup>10</sup>R. DiPersio, J. Simon, and T. Martin, "A Study of Jets From Scaled Conical Shaped-Charge Liners," *BRLMR-1298*, August 1960. (AD #246352)

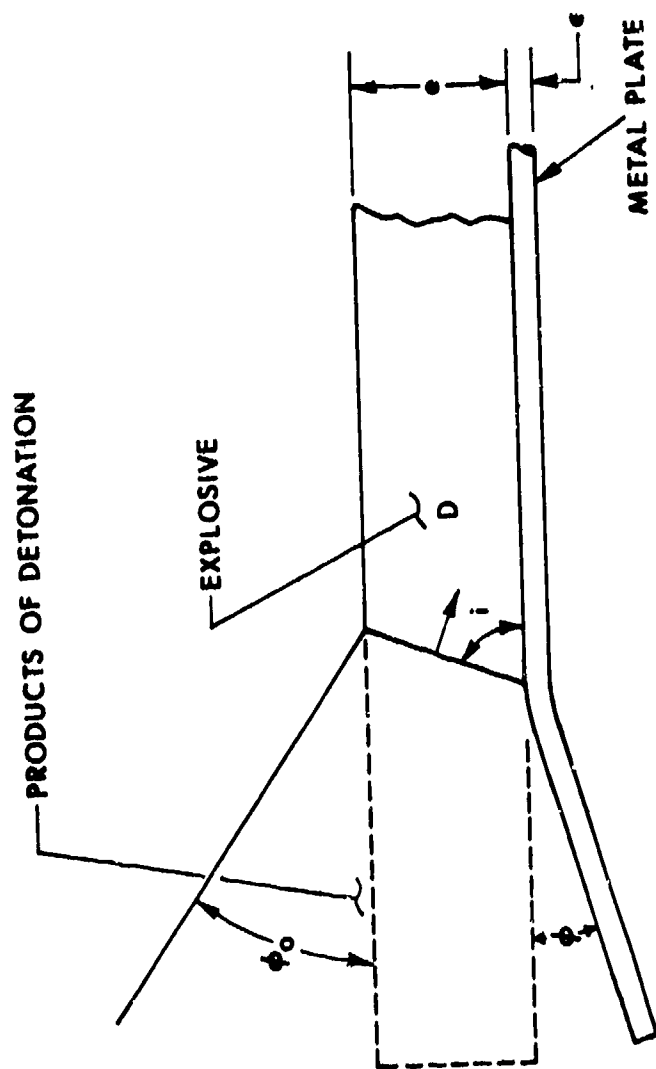


Figure 1. Projection of a Metal Plate by High Explosive

are determined for certain types of explosives.<sup>11</sup> Figure 2 illustrates the functional relationship between  $\phi_0$ ,  $K$ , and  $i$ .  $K$  is a constant over the range of  $i$  considered for typical shaped-charge designs, i.e. a conical liner contained in a cylinder of explosive. Figure 3 illustrates the linear relationship, Equation 1, that  $1/\phi$  has with the ratio of liner mass per unit volume to explosive mass per unit volume, given as

$$\mu = \frac{\rho_e}{\rho_{H.E.}} e,$$

where  $\rho_{H.E.}$  is the density of the explosive. The values of  $1/\phi_0$  and  $K$  used in Equation 1 are the Y-intercept and the ratio of the slope of the line to the density of explosive respectively. Equation 1 along with the Taylor formula,

$$v_0 = 2D \sin \frac{\phi}{2},$$

where  $D$  is the explosive detonation rate, will result in collapse velocities,  $v_0$ , obtained by Gurney.<sup>12</sup> Two types of explosive compositions are shown in Figure 3 (data taken from reference 11) which represent the linear function at a constant grazing (parallel) incidence,  $i$ , of the detonation from the normal to the metal surface (see Figure 1).

From the theory of Defourneaux, as the detonation wave sweeps toward the base of a typical shaped charge,  $\phi$  decreases due to the decrease in the explosive thickness,  $e$ , shown in Figure 3. This assumption that  $\phi$  decreased monotonically with a decrease in  $e$  is justified for most of the liner collapse since there is sufficient time for the liner to undergo several shock reverberations and achieve a bending angle close to its maximum before entering the flow of jet formation. However, the region near the apex of the cone

<sup>11</sup>M. DeFourneaux and L. Jacques, "Explosive Deflection of a Liner as a Diagnostic of Detonation Flows," *Proceedings Fifth Symposium (International) on Detonation*, ACR-184 Office of Naval Research-Department of Navy, pp. 457-466, Pasadena, California, August 18-21, 1970.

<sup>12</sup>W. Gurney, "The Initial Velocity of Fragments from Bombs, Shells, and Grenades," BRL Report No. 405, Sept. 1943. (AD #ATI36218)

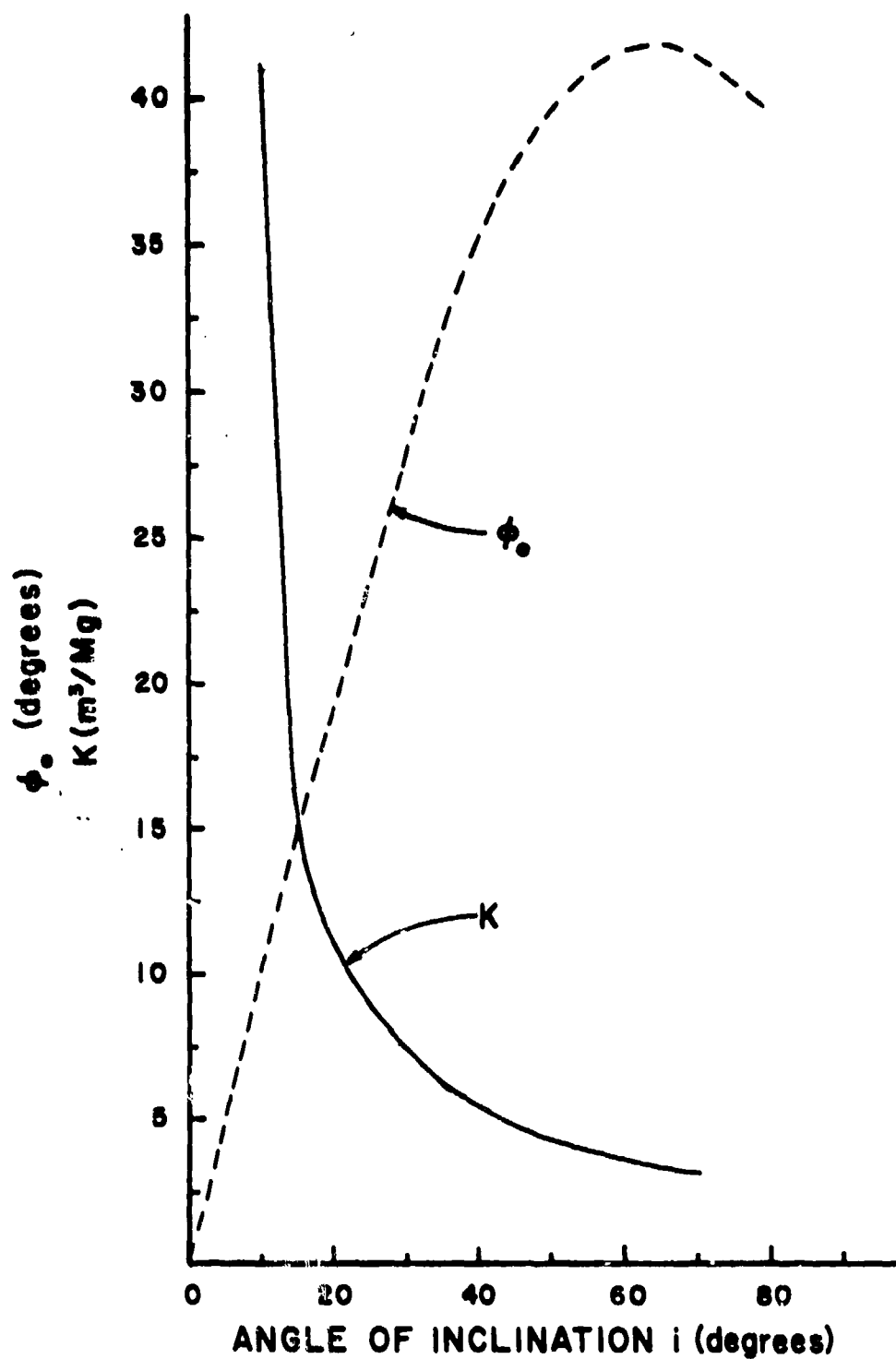


Figure 2.  $\phi_0$  and  $K$  are functions of the detonation wave angle to the liner.

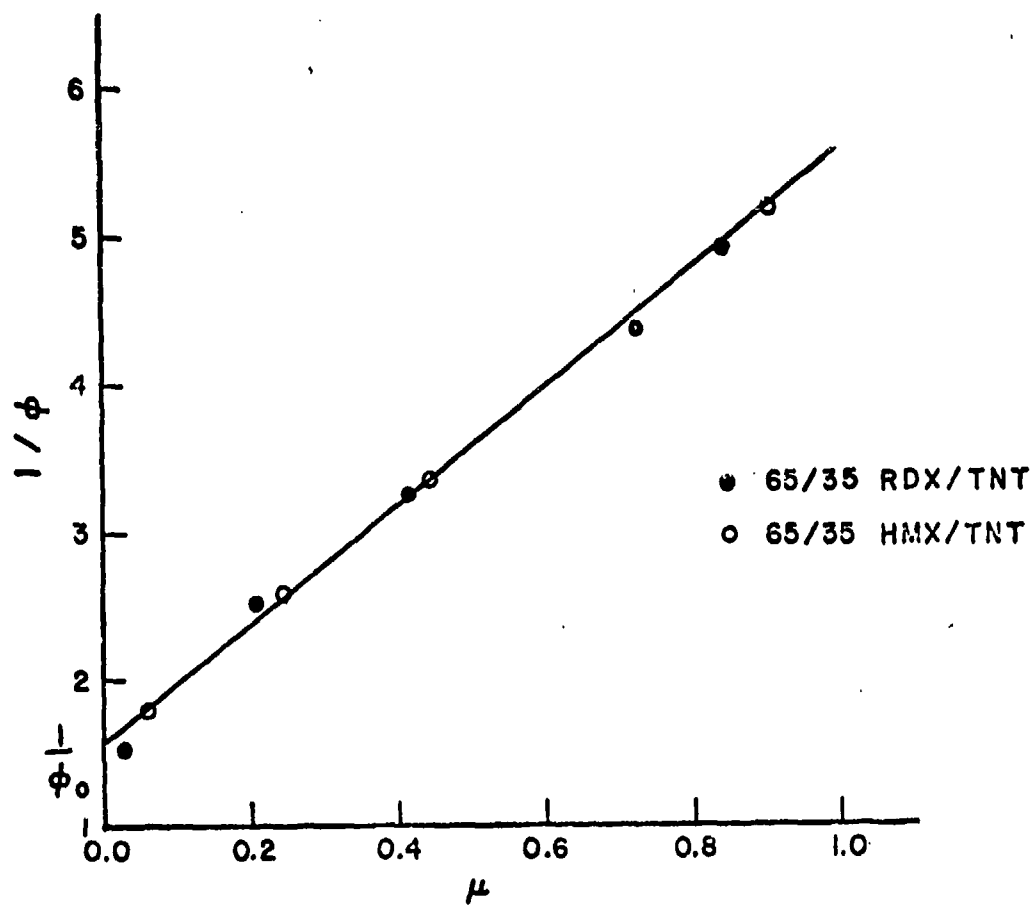


Figure 3.  $1/\phi$  vs  $\mu$  as determined from experiments. Data is taken from Reference 11.

enters the flow shortly after its initial acceleration and hence does not achieve its ultimate bending angle. Therefore, for material close to the liner's apex,  $\phi$  will increase to a maximum. This maximum will be located near liner elements that originate from a position approximately 40% of the liner height measured from the apex of the cone. After this point,  $\phi$  will decrease according to theory. We call this "the inverse collapse process." Associated with this process is the total time that a liner element takes to reach the axis. The equation for the collapse time is

$$\tau = (r \sin i) / (D(\sin(\alpha + \phi) - \sin \alpha)), \quad (2)$$

where several new variables are introduced.  $r$  is the instantaneous distance of a liner element from the axis before its collapse.  $\alpha$  is the half angle of the cone of a conical liner or, in the general case, the instantaneous angle with the axis made by a tangent line to the liner.  $\tau$  will be affected by the inverse collapse process. Material close to the liner apex enters the flow of jet formation sooner than would be predicted by the collapse of Defourneaux.<sup>4</sup>

In order to calculate the inverse collapse process, a new equation for the bending angle,  $\phi$  and an iterative scheme were added to the code. The new equation for  $\phi$  is

$$\phi_N = \phi e^{\sqrt{\frac{b}{\tau}}}, \quad (3)$$

where

$$b = C_1 \left[ \frac{k\rho e}{eB} \right],$$

$C_1$  is a constant determined from shape-charge collapse data.<sup>13</sup> The iterative scheme, between Equations (2) and (3), continues until the criteria of  $\phi_N$  approaching  $\phi$  within an epsilon is satisfied.

Having determined  $\phi$  and  $\tau$ , we can proceed with the collapse process.

<sup>13</sup>F.E. Allison and R. Vitali, "An Application of Jet Formation Theory to the 105-mm Shaped-Charge," BRLR-1165, March 1962. (AD #277458)



The velocity of collapse of the liner walls,  $v_o$ , toward the charge axis to form the jet is given by

$$v_o = 2D(\sin(\phi/2)) / (\sin i) \quad (4)$$

The apparent explosive detonation velocity, with respect to liner wall is given by

$$Da = D / \sin i \quad (5)$$

The substitution of Equation 5 into Equation 4 yields

$$v_o = 2Da \sin(\phi/2) \quad (6)$$

This is the so-called Taylor formula utilized in the code. The angle between the collapse direction and the charge axis is given by

$$\gamma = (\pi/2) - (\alpha + (\phi/2)). \quad (7)$$

The liner element first hits the axis at a distance,  $\overline{sp}$ , from the liner apex, given by

$$\overline{sp} = (z \sin \phi) / (\cos \alpha (\sin(\alpha + \phi) - \sin \alpha)), \quad (8)$$

where  $z$  is the axial component of the liner element position before its collapse.

While the liner is collapsing, the angle formed by a tangent to the collapsed portion on the axis and the axis itself (called the collapse angle) is computed from

$$\tan(\beta - \alpha) = \frac{\Delta z [\sin(\alpha + \phi) - \sin \alpha] \tan \phi + r \Delta \phi \cos \alpha}{\Delta z [\sin(\alpha + \phi) - \sin \alpha] - r \Delta \phi \cos \alpha \tan \phi}, \quad (9)$$

where  $\Delta z$  is the axial increment chosen in the computational scheme, and  $\Delta \phi$  is the incremental change in  $\phi$  between adjacent liner elements. The cartesian coordinates of a liner element (which originates at position  $z, r$ ) during collapse are given by the pair of equations:

$$\begin{aligned} x(z, t) &= z + v_o n \Delta t \sin(\alpha + (\phi/2)) \\ y(z, t) &= r - v_o n \Delta t \cos(\alpha + (\phi/2)), \end{aligned} \quad (10)$$

where  $\Delta t$  is the time interval taken by the detonation wave between successive liner elements and  $n$  is a positive integer. These equations apply in the time interval given by

$$0 < n\Delta t \leq \tau.$$

Figure 4 is an illustration of the relationship of the variables employed in the BASC code of a generalized axisymmetric collapse of a typical shaped charge. Shown on Figure 5 are drawings giving a detailed description of the collapse process. Figure 5A shows the velocity vectors of an element at point, P, on the collapsing liner. The element is projected toward the axis of symmetry with a collapse velocity,  $v_o$ , and a bending angle,  $\phi$ . When the detonation wave with velocity,  $D$ , has progressed a distance,  $P'$ , (i.e. from point P to point Q) during the time interval,  $\tau$ , then the element initially at point P will collide with the cone axis, producing the geometrical relations at the collision or stagnation point,  $sp$ , shown in Figure 5B. This relationship at the stagnation point is with respect to a coordinate system moving at the stagnation point velocity,  $v_{sp}$ . The velocity of the liner wall flowing into the stagnation point is  $v_f$  and the angle between it and the cone axis is the collapse angle,  $\beta$ . Figure 5C shows a cross-section of the collapsing liner depicting the variables employed. Applying Bernoulli's equations at the stagnation point, we find that the flow velocity,  $v_f$ , separates into two equal but directionally opposite velocities. One is called the jet velocity,  $v_j$ , and the other is called the slug velocity,  $v_N$ . This relationship is shown on Figure 6. Resolving the flow velocity,  $v_f$ , at the stagnation point in the laboratory coordinate system (Figure 6A), the following set of equations are obtained:

$$v_j = v_f + v_{sp} , \quad (11)$$

$$v_N = v_{sp} - v_f . \quad (12)$$

In accordance with the laws of conservation of mass and momentum, when the liner material reaches the cone axis after its collapse, it proceeds either as a fast-moving jet or as the more massive but slower-moving slug (Figure 6A). The jet velocity equation that results is

$$v_j = v_o \cos(\alpha + (\phi/2) - (\beta/2)) / \sin(\beta/2) , \quad (13)$$

and the equation for the slug velocity is

$$v_N = v_o \sin(\alpha + (\phi/2) - (\beta/2)) / \cos(\beta/2) . \quad (14)$$

The relative distribution of mass (Figure 6B) that results in jet and

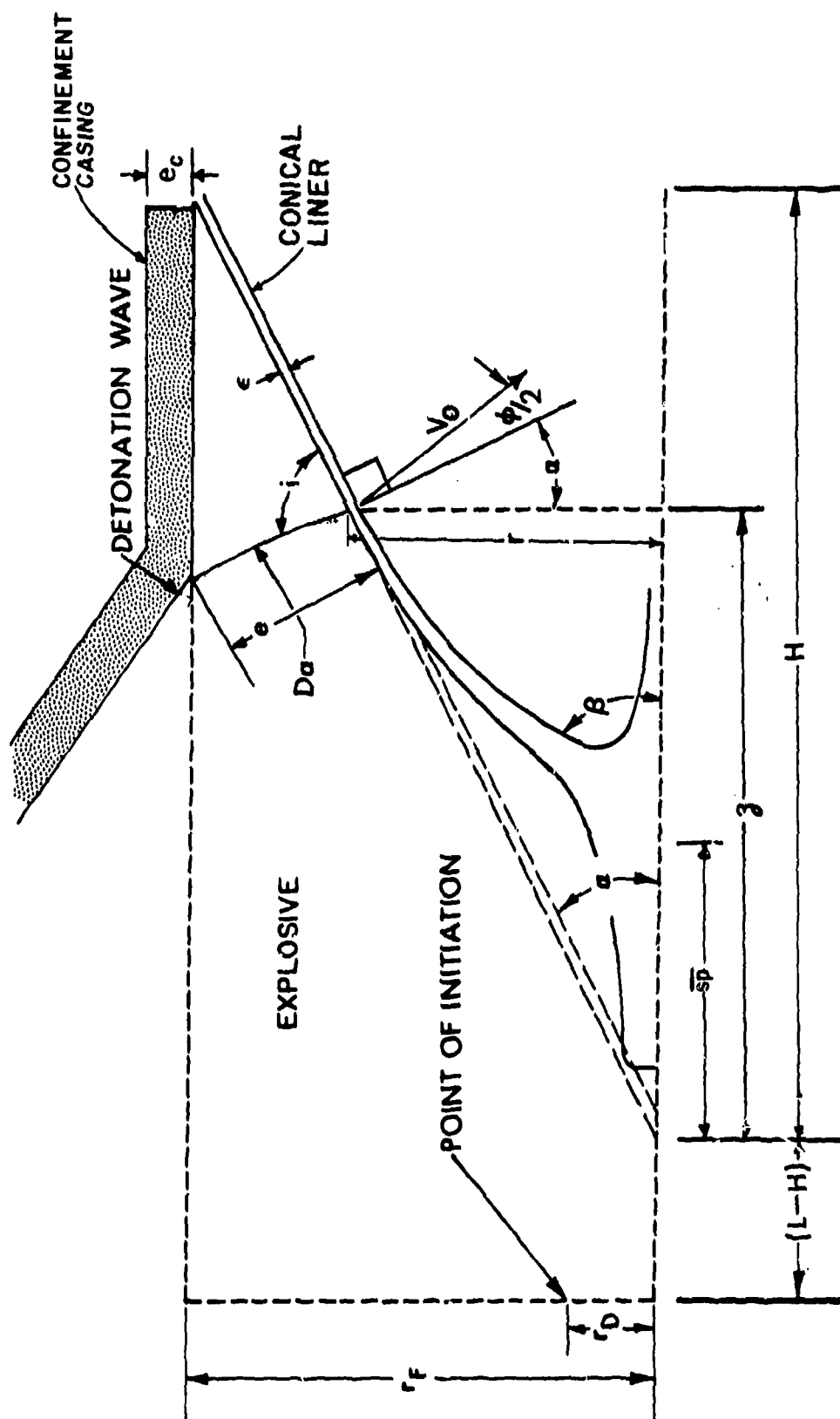
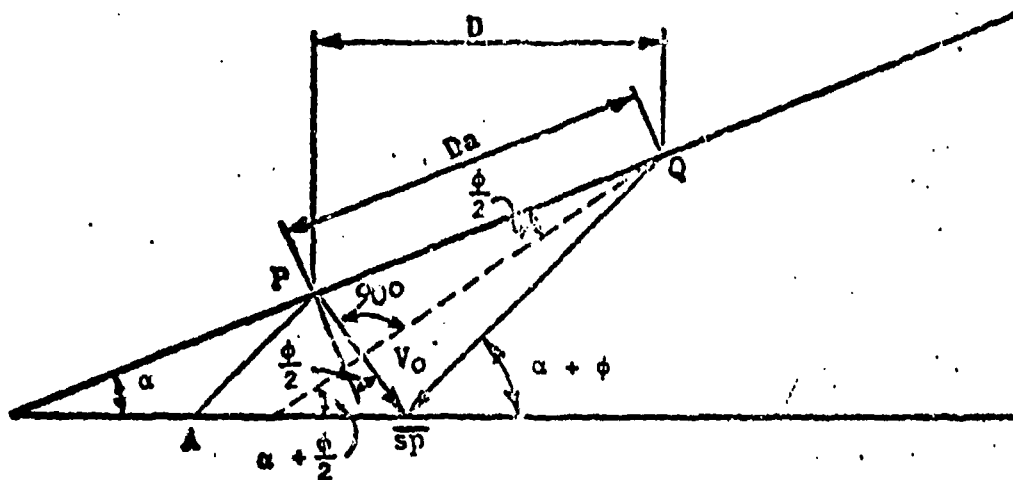
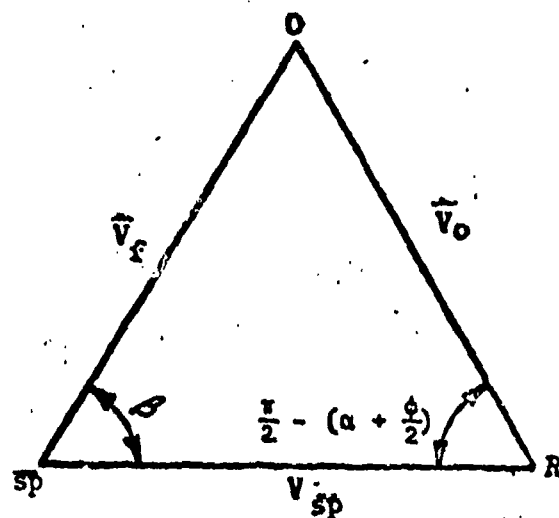


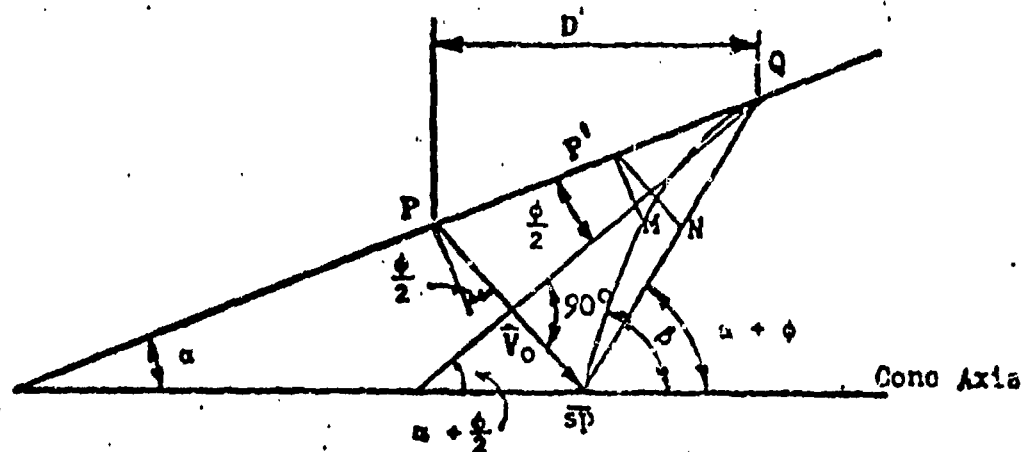
Figure 4. Line drawing illustrating the quantities employed in the BASC code of a generalized axisymmetric collapse of a shaped-charge.



A. Velocity vectors of an element of the collapsing liner.



B. Geometrical relations at the stagnation point.



C. Cross-section of collapsing liner.

Figure 5. Illustration of the relationship between variables and the collapse process.

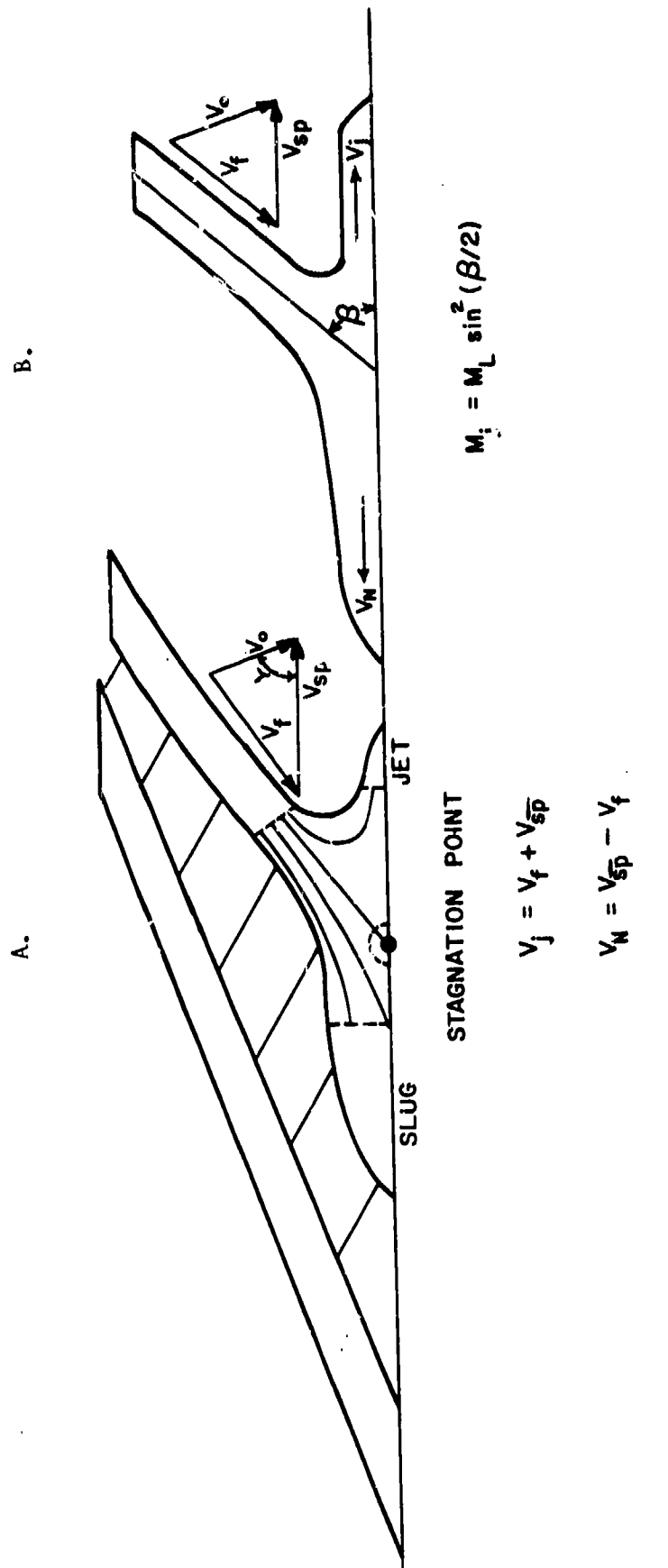


Figure 6. Illustration of the resolved variables in the laboratory coordinate system.

slug material are calculated, respectively, by

$$\frac{dm_j}{dm_L} = \sin^2(\beta/2)$$

(15)

and

$$\frac{dm_N}{dm_L} = \cos^2(\beta/2) .$$

Also, the relative distribution of the kinetic energy for the jet and slug, respectively, are

$$\frac{dE_j}{dE_L} = \cos^2(\alpha + (\phi/2) - (\beta/2))$$

(16)

and

$$\frac{dE_N}{dE_L} = \sin^2(\alpha + (\phi/2) - (\beta/2)) .$$

The variables  $dm_L$  and  $dE_L$  are the incremental change in the liner's mass and kinetic energy, respectively.

The impinging flow velocity,  $v_f$ , has been considered by shaped-charge researchers such as Walsh, et. al.<sup>14</sup>; Cowan, et. al.<sup>15</sup>; and Chou, et. al.<sup>8</sup> to be critical in the jet-formation theory. As illustrated in Figure 7, when  $v_f$  is less than the material sound speed,  $c$ , the jet formed is coherent or a good jet (Figure 7A). Even when  $v_f$  is slightly greater than  $c$ , this too forms a coherent jet (Figure 7B). But, when  $v_f$  is sufficiently greater than  $c$ , the jet will be incoherent or bifurcated. We call it a no-jet condition (Figure 7C). From equations 11 and 12, solving for  $v_f$  yields

$$v_f = .5(v_j - v_N) .$$

(17)

We then use the following relationship as the jet limiting criteria for a coherent jet:

$$M_R = v_f/c < 1.23 .$$

(18)

<sup>14</sup>J.M. Walsh, R.G. Shreffler, and F.J. Willig, "Limiting Velocity Conditions for Jet Formation in High Velocity Collisions," *Journal of Applied Physics*, Vol. 24, No. 3, pp. 349-359, March 1957.

<sup>15</sup>G.R. Cowan and A.H. Holtzman, "Flow Configurations in Colliding Plates: Explosive Bonding," *Journal of Applied Physics*, Vol. 34, No. 4, pp. 928-939, April 1963.

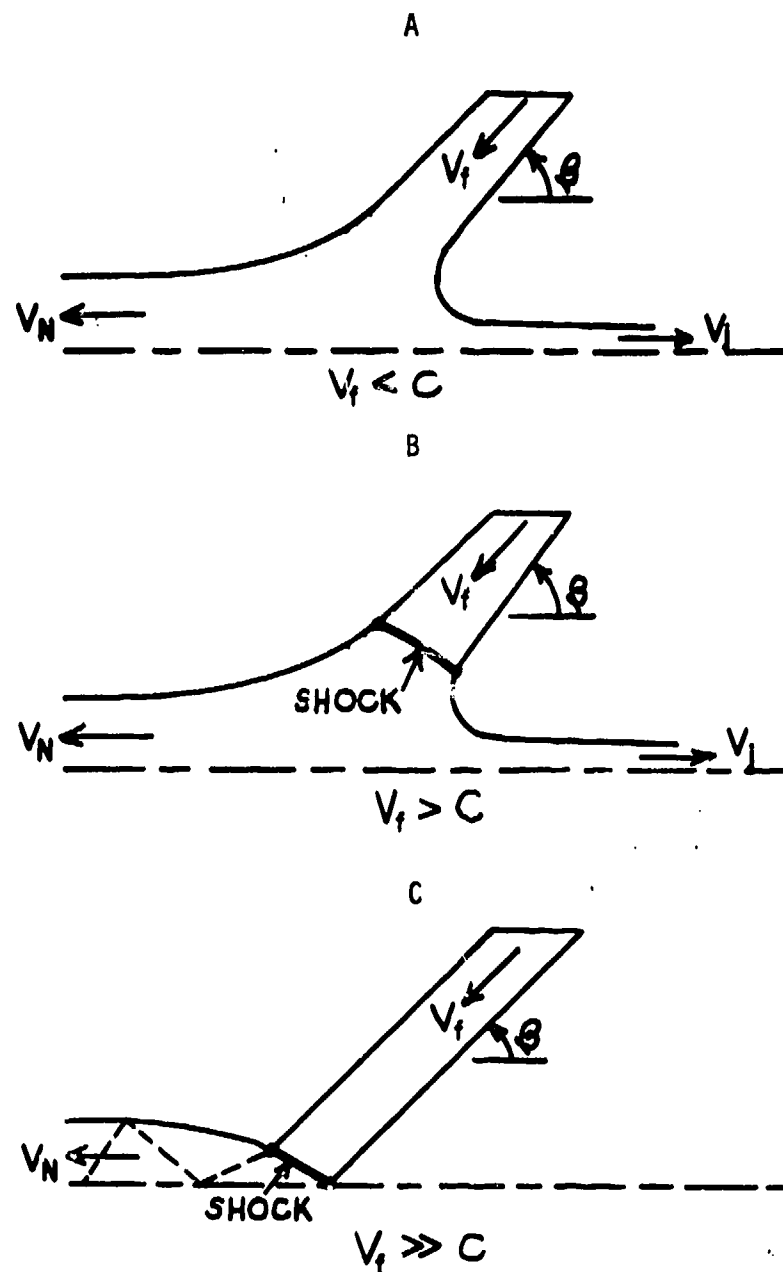


Figure 7. Illustration of the variables employed in the BASC code to limit jet formation. The jet limiting criteria for a coherent jet (A and B) is  $M_R = v_f/c < 1.23$ .

The threshold factor for a coherent jet,  $M_c = 1.23$ , is determined by comparing radiographic, jet-particle data measured<sup>16</sup> from experiments<sup>17</sup> with BASC results. The threshold factor holds for several materials considered in BASC calculations.

In order to determine the jet-tip velocity and the mass of the lead pellet to be used in the shaped-charge penetration theory, the inverse velocity gradient has to be equilibrated and the jet mass in this region to be compressed into a single zonal element, LP. This zone will then contain the so-called steady-state, lead pellet. The equilibrated jet-tip velocity,  $v_{jo}$ , is obtained by the following process:

$$v_j^{i+1} = \frac{v_j^{i+1} \left[ \frac{dm_j}{dm_L} \right]^{i+1} dm_L^{i+1} + v_j^i \left[ \frac{dm_j}{dm_L} \right]^i dm_L^i}{\left[ \frac{dm_j}{dm_L} \right]^{i+1} dm_L^{i+1} + \left[ \frac{dm_j}{dm_L} \right]^i dm_L^i}$$

for  $v_j^{i+1} > v_j^i$  and  $1 \leq i \leq LP$ , where  $i$  is the  $i^{th}$  zonal element. The equilibrated jet-tip velocity is

$$v_{jo} = v_j^{LP} \quad (18)$$

The steady-state, lead pellet mass is

$$dm_{jo} = \sum_{i=1}^{LP} \left[ \frac{dm_j}{dm_L} \right]^i dm_L^i \quad (19)$$

In the theory of shaped-charge jet penetration into a target used in BASC, an important parameter is the time that a given liner element, which enters into the jet, arrives at the bottom of the target hole when penetration is in progress. Time is usually started from the moment the detonation wave reaches the apex of the liner. A time parameter,  $\theta$ , is defined by

$$\theta = t_z + \tau, \quad (20)$$

<sup>16</sup> Private Communication from J. Blische at BRL.

<sup>17</sup> R. DiPersio, W.H. Jones, A.B. Merendino, and J. Simon, "Characteristics of Jets from Small Caliber Shaped Charges with Copper and Aluminum Liners," BRLMR No. 1866, September 1967. (AD #823839)



where  $t_z$  is the time at which the detonation wave reaches the liner element<sup>z</sup> at an initial axial distance,  $z$ , from the apex and  $\tau$  is the time taken by this element to collapse onto the charge axis. The sum of standoff distance from the liner base to the initial surface of the target,  $h$ , and the liner height,  $H$ , defines  $Z_0$ :

$$Z_0 = h + H. \quad (21)$$

The time for the tip of the jet to reach the target surface is

$$T_0 = Z_0 / v_{j0} \quad (22)$$

It is assumed that the jet tip originates from the zonal element, LP, and contains the highest velocity of the jet,  $v_{j0}$ . An element from the liner between LP and base results in a jet velocity,  $v_j$ , which is less than  $v_{j0}$  and is a function of its initial axial distance,  $z$ . A parameter,  $\mu$  is defined by

$$\mu = v_{j0} / v_j \quad (23)$$

The time for the portion of jet formed from a given element of the liner to reach the bottom of the target hole is

$$T = \mu^{(1+k)} \left[ T_0 + (1+k) \sum_0^i \mu^{-k} (\Delta f)_i \right] \quad (24)$$

where  $(\Delta f)_i = f_i - f_{i-1}$  and denotes the  $i^{\text{th}}$  zonal element. In

this equation,  $k = \sqrt{\rho_j / \rho_c}$ , where  $\rho_j$  is the jet density (assumed equal to the liner density), and  $\rho_c$  is the target density.  $f$  is a time parameter given by

$$f = ((\theta / \mu) - (\overline{sp} / v_{j0})) , \quad (25)$$

where  $\overline{sp}$  was previously defined by Equation 8, and  $v_{j0}$ ,  $\theta$ , and  $\mu$  were defined by Equations 18, 20, and 23. The value<sup>j0</sup> of  $\mu$  to be used outside the bracket in Equation 24 is that which applies to the original elemental liner position.  $\Delta f$  is the difference in  $f$  values between adjacent elemental positions, and the summation applies to all adjacent elemental positions up to the point in question on the liner.

While the jet is still continuous in nature, it stretches due to its velocity gradient and, therefore, decreases in diameter with increase in time. The equation for calculating the jet radius,  $r_j$ , is

$$r_j^2 = - \frac{2 r \epsilon D \sin^2 (\beta/2)}{\left. \frac{\Delta z}{\Delta t_z} \right|_T \sin i} \quad (26)$$

In this equation,  $r$  is the initial radial position of the liner element from the charge axis,  $\epsilon$  is the liner thickness at this position, and  $\beta$  is the collapse angle that is formed when this

element reaches the charge axis. The factor,  $\left. \frac{\Delta z}{\Delta t_z} \right|_T$ , is given by

$$\left. \frac{\Delta z}{\Delta t_z} \right|_T = (T - t_z) \frac{\Delta v_j}{\Delta t_z} - v_j + \frac{\Delta(z - r \cot(\beta/2))}{\Delta t_z} \quad (27)$$

In this equation,  $t_z$  is defined following Equation 20 and  $\Delta t$  is the incremental time interval between arrival of the detonation wave between successive liner elements (a constant). The value computed by Equation 27 is a function of time,  $T$ , which starts at zero when the detonation wave first reaches the liner apex. It is a negative value which increases in absolute value as  $T$  increases. Therefore, it can be seen from Equation 26 that the jet radius, which originates from material at any point on the liner, decreases with an increase in  $T$ . The minus sign in Equation 26 is necessary to make  $r_j^2$  a positive quantity.

When the jet cannot sustain any further stretching, it breaks up into individual axial particles. It is assumed that this occurs throughout the whole jet at the same time. The breakup time for the jet is designated as  $T_1$  and, at present, must be obtained from experimental observations. At times greater than the jet breakup time, the individual jet particles do not stretch in length or decrease in diameter. The constant jet particle radius that one obtains for material originating from a given element of the liner

is calculated by Equation 26 in which the factor  $\left. \frac{\Delta z}{\Delta t_z} \right|_T$  is calculated

at the breakup time,  $T_1$ , in Equation 27. However, radii of different particles are different due to the variability of the initial liner element position in Equations 26 and 27.

The equations that are used for jet penetration theory are dependent upon the standoff distance between the charge and the target. If the target is placed close enough to the charge so that the jet tip reaches the target before the time of jet breakup, one set of equations are used. On the other hand, if the jet particulates before reaching the target surface, a different set of equations must be used. In the former case, even though the jet

starts penetrating in the continuous state, it becomes discontinuous before the end of penetration and different equations are required after time  $T_1$ . The variables and the two states of penetration are illustrated in Figure 8.

The jet penetrates the target in both the continuous and discontinuous state whenever the following criteria is satisfied:

$$Z_0 < v_{jo} T_1 \quad (28)$$

The depth of penetration into the target,  $p$ , for any  $T$ , such that  $T_0 \leq T \leq T_1$ ,

$$p = Z_0 \left[ \left( \frac{T}{T_0} \right)^{(k/(k+1))} - 1 \right] \quad (29)$$

where all factors have been previously defined. For times greater than  $T$ , the equation used is

$$p = Z_0 \left[ (k+1) \left( \frac{T_1}{T_0} \right)^{(k/(k+1))} \frac{T}{T_1 + kT_1} - 1 \right] \quad (30)$$

The total penetration depth into the target is calculated by

$$P_T = (k+1)(v_{jo} T_1)^{\frac{k}{k+1}} Z_0^{k+1} - \sqrt{k(k+1) U^{\min}_{T_1} (v_{jo} T_1)^{\frac{k}{k+1}} Z_0^{\frac{k}{k+1}} - Z_0} \quad (31)$$

The only hitherto undefined term in Equation 31 is the factor  $U^{\min}$ . This is called the minimum penetration velocity. According to theory, the velocity of penetration into the target monotonically decreases with increase in penetration depth until it reaches the

value  $U^{\min}$ . When  $U^{\min}$  is reached, penetration stops and all remaining jet material just piles up at the bottom of the target

hole.  $U^{\min}$  is an empirical constant whose value depends upon the target material and its hardness. It is invariant with standoff distance. With a given charge, target, and standoff distance, the computer calculates the constant total penetration depth from Equation 31. This value is used by the computer as a signal to stop calculating  $p$  in Equation 30 and also to determine the time at the end of the penetration process. The radius of the hole in the target before the time of jet breakup is given by

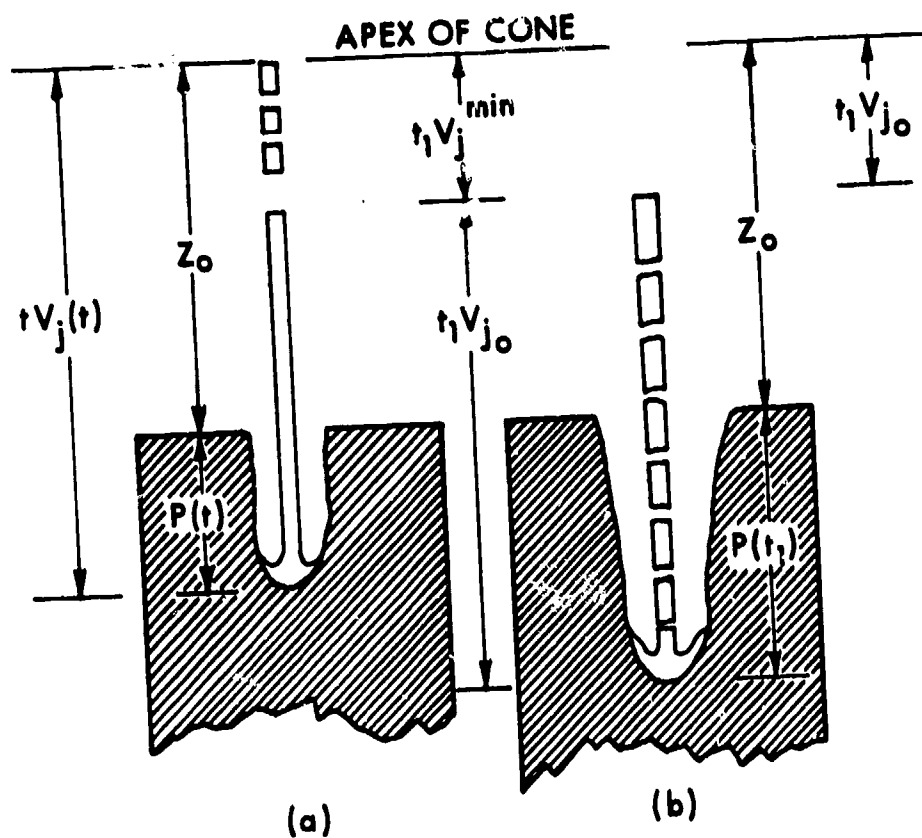


Figure 8. Schematic cross section of a shaped-charged jet penetrating a target in two states:

- (a) Continuous and discontinuous,
- (b) Discontinuous.

$$r_c = \sqrt{\frac{\rho_j}{2kc_k}} \sqrt{\frac{T_1}{T_0}} \left(\frac{Z_0}{Z_0+p}\right)^{\frac{3+k}{2k}} v_{jo} r_j . \quad (32)$$

This equation contains one last empirical constant of the target,  $c_k$ . This is the jet kinetic energy needed to produce a unit of hole volume in the target. Its value depends upon the target used and its hardness. The value of  $r_j$  in Equation 32 is obtained by Equations 26 and 27. After the breakup time,  $T_1$ , the hole radius is obtained by

$$r_c = \sqrt{\frac{\rho_j}{2kc_k}} \frac{1}{k} \left[ (k+1) \left(\frac{Z_0}{v_{jo} T_1}\right)^{\frac{1}{k+1}} - \frac{Z_0+p}{v_{jo} T_1} \right] v_{jo} r_j , \quad (33)$$

where  $r_j$  is now obtained by Equations 26 and 27 with the restriction that the factor  $\frac{\Delta Z}{\Delta t_z}|_T$  must be evaluated at  $T = T_1$ . The hole profile, as computed by Equation 33, is terminated when the penetration depth,  $p$ , reaches the value  $P_T$  as calculated by Equation 31.

The last set of penetration equations is used when the standoff distance between charge and target is so large that the following condition is satisfied:

$$Z_0 > v_{jo} T_1 . \quad (34)$$

In this case, all penetration is accomplished by the jet while it is particulated in nature. The equation for penetration depth is then

$$p = v_{jo} (T - T_0) T_1 / (T_1 + (T/k)) \quad (35)$$

where the time factor,  $T$ , varies between the time of first target impact,  $T_0$ , and the time of last jet penetration,  $T_p$ . The total penetration depth is given by

$$P_T = [v_{jo} T_1 - \sqrt{U^{\min} T_1 (v_{jo} T_1 + (Z_0/k))}] k . \quad (36)$$

The radius of the hole produced in the target as a function of its depth is given by

$$r_c = v_{jo} r_j (1 - (p/k v_{jo} T_1)) \sqrt{\rho_j / 2k c_k} , \quad (37)$$

where the jet radius,  $r_j$ , is obtained by Equations 26 and 27 with the restriction that the factor  $\frac{\Delta Z}{\Delta t_z} \Big|_T$  is evaluated at  $T = T_1$ .

The variation of total penetration depth with standoff distance is computed by means of Equations 31 and 36. The factor  $Z_0$  is the only variable in these equations. Equation 31 is used first until  $Z_0$  increases to the value  $v_{jo} T_1$ , then Equation 36 is used.

### III. CALCULATIONAL SCHEME

The BASIC code enables one to perform parametric studies for designing warheads. The variables employed in the code of the generalized, axisymmetric collapse of a shaped charge were illustrated previously in Figure 4. The parameters which can vary include the following:

$\alpha$	(ALPHA)	The half angle of the liner (degrees)
K	(CON)	The empirical constant for the detonation products. Value known for Composition B explosive.
$\epsilon$	(EPS)	The thickness of the liner (cm)
$\rho_j$	(RHOJ)	The density of the liner (gm/cm <sup>3</sup> )
$\rho_c$	(RHOC)	The density of the target (gm/cm <sup>3</sup> )
$r_F$	(RF)	Radius of the base of large (cm).
H	(H)	Height of the liner (cm). If H is zero, H will be calculated by $H = \frac{r_F}{\tan \alpha}$ .

D	(D)	Detonation velocity (cm/ $\mu$ s)
SO	(SO)	Standoff distance between base of liner and target, (cm)
$C_K$	(CK)	Constant for determining hole volume
$T_1$	(T1)	Jet breakup time ( $\mu$ s)
$U^{MIN}$	(UMIN)	Velocity cutoff for the penetration of jet into a target (cm/ $\mu$ s)
L	(DPOINT)	Initially the total length of the charge (cm) but is converted to be the initiation point of the explosive (i.e. DPOINT = L-H).
RDPT	(RDPT)	Radius above axis where explosive is initiated (cm).
JOHN	(JOHNI)	If JOHNI greater than zero, $\phi_0$ will vary; if not, $\phi_0$ will be a calculated constant.
N	(N)	The number of zones into which the grid is subdivided. If N is zero, default is seventy zones.

The code will set up the grid based upon Figure 4 and the equation for increment in Z is

$$DZ = H/N \quad (38)$$

and from this the time increment is

$$DTZ = DZ/D \quad (39)$$

Therefore, at each increment in Z we have also the corresponding time increment.

The code then marches sequentially through the governing equations (1-19), calculating and storing jet formation information to be printed and plotted at the end of the iteration. Those variables include:

i	I	$i^{th}$ zone
$\alpha$	ALPHA(I)	Tangent angle of liner wall to the axis (degrees)
z	Z(I)	z distance from apex to base of liner

$t_z$	TZ(I)	Time at the $i^{th}$ zone
$\gamma$	GAMMA(I)	Angle (degrees)
E	E(I)	Thickness of explosive above $i^{th}$ zone
$\phi$	PHI(I)	Angle (degrees)
$\beta$	BETA(I)	Angle (degrees)
$\Delta\phi$	DPHI(I)	Increment of angle $\phi$
$\frac{1}{\phi}$	RPHI(I)	Reciprocal of angle $\phi$
$v_o$	V(I)	Collapse velocity (cm/ $\mu$ s)
r	R(I)	Radius of the $i^{th}$ zone (cm)
$\tau$	TAU(I)	Time to collapse $i^{th}$ zone ( $\mu$ s)
$\overline{sp}$	C(I)	Distance from the apex of the liner where $i^{th}$ zone will hit the axis (cm)
$v_j$	VJ(I)	Velocity of the jet in (cm/ $\mu$ s)
$v_N$	VN(I)	Velocity of the slug in (cm/ $\mu$ s)
$\frac{dm_j}{dm_L}$	DMJ(I)	Relative mass of the jet, dimensionless
$\frac{dm_N}{dm_L}$	DMN(I)	Relative mass of the slug, dimensionless
$\frac{dE_j}{dE_L}$	DEJ(I)	Relative energy of the jet, dimensionless
$\frac{dE_N}{dE_L}$	DEN(I)	Relative energy of the slug, dimensionless
$v_f$	RV(I)	Relative velocity between the jet and slug



In the penetration portion of BASC, the constants listed below must be calculated:

$$VJO = \text{jet tip velocity from Equation 18} \quad (40)$$

$$ZO = H + SO \quad (41)$$

$$TO = \frac{ZO}{VJO} \quad (42)$$

$$AKAY = \sqrt{\frac{\rho_j}{\rho_c}} \quad (43)$$

$$PT = \text{Total penetration (cm) from Equation 31 or 36} \quad (44)$$

The next set of outputs are for the penetration of the jet into the target. The code also marches through the governing Equations 20-37, calculating and storing penetration information to be printed and plotted at the end of the iteration. The variables calculated and stored are listed below:

i	I	i <sup>th</sup> zone
$\mu$	AMU(I)	Relative velocity between the jet tip velocity and jet velocity of the i <sup>th</sup> zone
$\theta$	THETA(I)	Time parameter in microseconds
f	F(I)	Time parameter in microseconds
$\Delta f$	DF(I)	Time increment of f in microseconds.
T	T(I)	Time that the i <sup>th</sup> element reaches the bottom of the target hole in microseconds.
$\Delta t$	DT(I)	Time increment of T
r	RSQ(I)	is $\sqrt{r_j^2}$ which is the radius of the i <sup>th</sup> element of the jet.
A	A(I)	is equal to $z - r \cot \beta/2$ .
$\Delta A$	DELA(I)	is the increment of A(I).
$\Delta v_j$	DVJ(I)	is the increment of jet velocity.

$\frac{\Delta z}{\Delta t_z}$	DZODT(I)	Equation 27 in governing equations
$r_c$	RC(I)	Radius of the hole in target of the $i^{th}$ zone in centimeters
$p$	p(I)	Depth of penetration in centimeters

After these parameters have been printed and plotted, the code then returns to start and begins another case. This is continued until the end of file (i.e. next problem) is encountered causing the run to terminate.

#### IV. COMPARISON OF BASC CODE RESULTS

The performance of the BASC code is best illustrated by presenting results of a calculation and comparing these results, where possible, with experimental data or with results from other numerical techniques. The first set of comparisons will be with experimental results from the following:

- a. 105-mm, unconfined,  $42^\circ$ , copper-lined shaped charge with a Composition B explosive fill tested by Allison and Vitali<sup>13</sup>.
- b. A study of jets from scaled,  $42^\circ$ , copper-lined, conical shaped charges filled with Composition B explosive (test by DiPersio, et. al.<sup>10</sup>).
- c. 3.2-inch, BRL precision, shaped charge with a copper liner  $42^\circ$  apex angle, and Composition B explosive fill, having its jet characteristics measured<sup>16</sup> from radiographic data recorded at the BRL.
- d. A study of jet characteristics from small-caliber shaped charges with copper and aluminum liners and variation in apex angle from  $20^\circ$  to  $90^\circ$ . All charges were filled with Composition B explosive. The tests conducted at the BRL by DiPersio, et. al.<sup>17</sup>

The second set of comparisons will be another numerical technique. The other technique is the two-dimensional, hydrodynamic computer program based upon the Eulerian numerical scheme called the BRLSC (Ballistic Research Laboratory Shaped Charge) code.<sup>18</sup> The BRLSC

<sup>18</sup>M.L. Gittings, "BRLSC: An Advanced Eulerian Code for Predicting Shaped Charges," Vol. I, BRL CR 279, Prepared by System, Science and Software, December 1975. (AD #A023962)

code is a modified version of the HELP (Hydrodynamic Elastic-Plastic) code<sup>19</sup> developed by System, Science and Software.

Figure 9 and 10 are experimental comparisons for the 105-mm, unconfined, shaped charge. Collapse velocity,  $v_o$ , and jet velocity,  $v_j$ , as a function of the relative distance from the apex of the cone are shown on Figure 9. The dashed line is the jet velocity after jet compression and illustrates exact agreement with both the jet tip velocity and position on the axis between BASC and experiment. This illustrates that jet tip is compressed into one element at a position which is approximately 40% of the distance from the apex of the cone. The collapse velocity is slightly higher than that calculated by Allison and Vitali, but the overall agreement is good. Figure 10 shows a comparison of the predictions of jet mass versus cone mass between the BASC code and Allison and Vitali. Those two predictions are in slight disagreement with one another because Allison and Vitali were unable to recover 100% of all the slug material. They used essentially the same theory as the BASC code to predict jet mass, and all of the material is required to accurately predict the true conservation laws.

Figure 11 is the second experimental comparison of the scaled, shaped charge with a heavy confinement casing. This figure shows collapse velocity,  $v_o$ , and jet velocity,  $v_j$ , as a function of the relative distance from the apex of the cone. The dashed line is the compressed jet velocity distribution. This again illustrates exact agreement for both jet tip velocity and its position on the axis between BASC and experiment. The open squares at approximately the 48% position from the apex of the cone shows the spread in the experimental data from the experimental scaled rounds. The BASC code results are identical for the same scaled rounds (see Reference 10). Shown at the base of the cone, for values greater than 80% of the distance from the apex of the cone, is a change in the slope of the jet velocity curve. This is due to gas leakage or breaking of the confinement casing. This phenomenon is modeled in BASC as a gradual change in the confinement factor,  $A$  in Equation 1, until it reaches zero, i.e. unconfined. This comparison shows very good agreement with the experimental results.

Table I is a tabulation of a comparison between measured, radiographic experimental data and BASC code results from the 3.2-inch, BRL, precision shaped charge. The jet tip velocity and mass at jet particle number one are in agreement, but the accumulated total jet mass from jet particle number one to jet particle number

<sup>19</sup> L.J. Hageman and J.M. Walsh, "HELP, A Multi-Material Eulerian Program For Compressible Fluid and Elastic-Plastic Flows in Two Space Dimensions and Time," BRL CR 39, Prepared by System, Science and Software, May 1971. (AD Nos. 726459 and 726460)

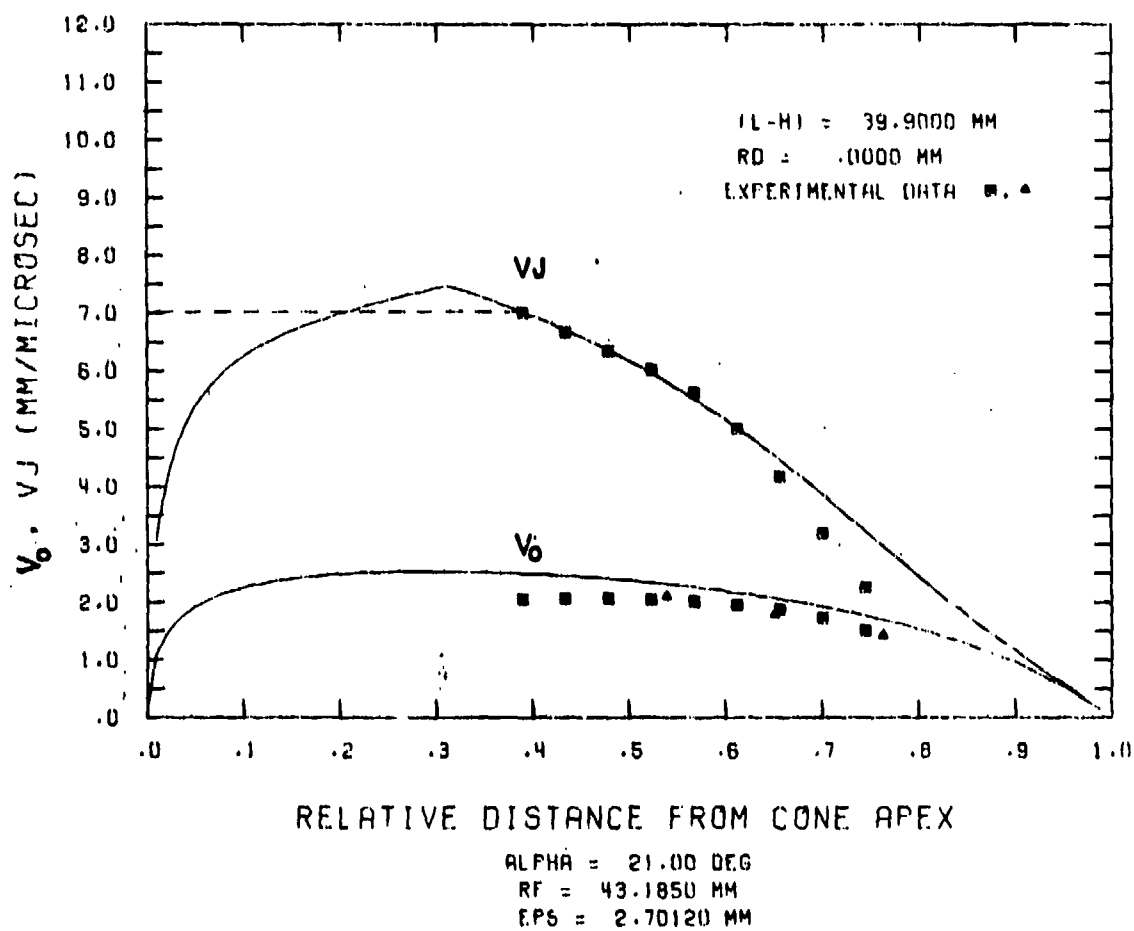


Figure 9. Comparison between experimental and BASC Code results of jet and collapse velocities from the 105-mm, unconfined, shaped charge.

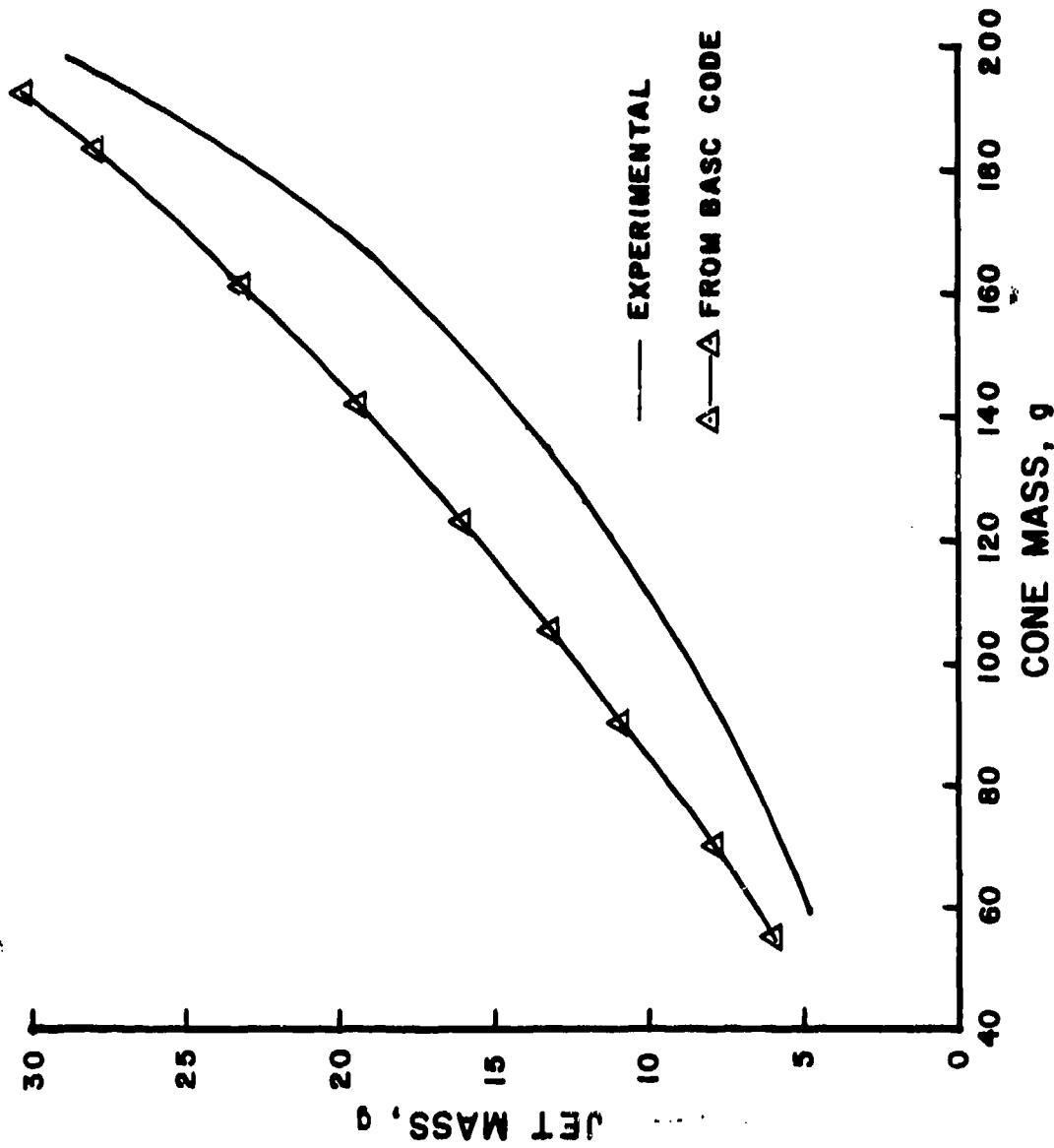


Figure 10. Comparison of calculations of jet mass vs cone mass between BASC code and experimental results from the 105-mm, unconfined, shaped-charge (Reference 13).

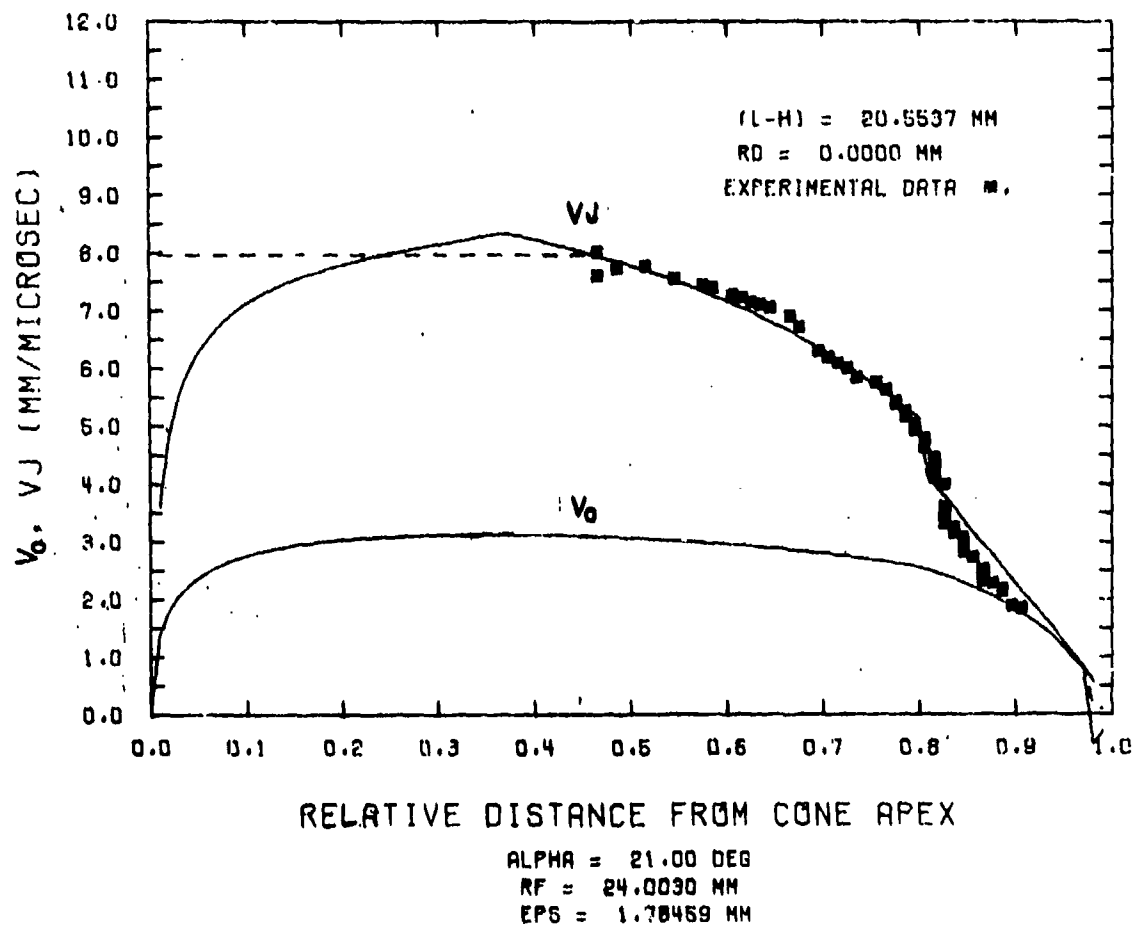


Figure 11. Comparison between BASC code and experimental results for scaled, heavily-confined, shaped charges (Reference 10). Jet and collapse velocities vs % of the distance from the apex of the cones are shown. Solid lines are BASC code results. The dashed line is the final, computed, compressed jet velocity profile. The jet tip for both experimental and BASC results is shown compressed into a position 48% from the apex of the cone.

forty-three is 10 grams heavier in the BASC code results. Also, the code results show 46% more total accumulated kinetic energy from the first forty-three jet particles than that obtained from experiment. Overall, this comparison is reasonable considering the material loss from such physical effects as erosion and possible errors in the data reduction measurements and calculations of the experimental results.

Table I. BASC Calculations and Measured, Experimental Radiographic Data<sup>16</sup> of Jet Velocity, Mass, and Kinetic Energy for the 3.2-in. (81.3-mm), BRL, Precision, Shaped Charge

Jet Particle No.	Velocity (km/s)		Accumulated Mass (g)		Kinetic Energy (kJ)	
	Measured	BASC	Measured	BASC	Measured	BASC
1	7.74	7.74	3.7	4.3	110.	128.
43	2.90	2.98	22.6	32.6	348.	507.

The last experimental comparison will be the characteristics of jets from small caliber shaped charges with copper and aluminum liners. In this study, only the liner designs referred to as 1½-inch (38.1-mm) liner base diameter charges in Reference 17 will be considered. The apex angle will include 20°, 40°, 60° and 90° cones for both the copper and aluminum liners. These are all confined with a steel casing around the charge. The pertinent dimensions for the shaped-charge designs considered can be found in Reference 17 on pages 9 and 10. The results from the BASC code calculations and the experiments are summarized on Table II. All of the results of the BASC code are within a 5% error boundary except the liner type with a 90° apex angle and aluminum cone which is 10%. The results of the calculations are slightly higher than the experiments for all of the aluminum liner types.

Table II. Summary of Results From the Small-Caliber, Shaped-Charge Study for the 38.1-mm (1½-in.), Base Diameter Design

Liner Type	Jet Tip Velocity, $V_{jo}$ (km/s)	
	Experiment	Predicted (BASC)
20° apex angle, Cu Cone	9.9	9.80
40° apex angle, Cu Cone	8.2	8.00
60° apex angle, Cu Cone	6.7	6.64
90° apex angle, Cu Cone	5.5	5.48

	<u>Liner Type</u>	<u>Experiment</u>	<u>Predicted (BASC)</u>
20°	apex angle, Al Cone	11.2	11.80
40°	apex angle, Al Cone	9.3	9.75
60°	apex angle, Al Cone	8.1	8.54
90°	apex angle, Al Cone	6.8	7.50

The final two sets of comparisons involve the two numerical techniques: BASC and BRLSC. First, Figure 12, is a comparison of the calculated flow field for the 105-mm, unconfined shaped charge reported<sup>2</sup> for the BRLSC code and the BASC code at 25  $\mu$ sec after initiation of the explosive. The radius of the jet and slug are in excellent agreement, but the tip of the jet is slightly behind in the BASC calculation. This is due to the fact that the BASC code takes very large time steps which will not allow the inverse velocity gradient to be equilibrated at this snap shot in time. The second set of comparisons between BASC and BRLSC codes is shown on Table III. The results shown are from an improved version of the BRLSC Code.<sup>20</sup> Table III is a tabulation of six calculated results from BASC and BRLSC as well as four experimental results, which can be used as a guide in this comparison. All of the calculations and experiments in this study involved identical copper liners (21°, 81.3-mm base diameter, 1.9-mm thickness). Calculations 1 through 4 employed TNT, Comp B, Octol and PBX, respectively as explosive filler.

Table III. Computational Matrix and Summary of Results

<u>Calculation Number</u>	<u>HE</u>	<u>Confinement</u>	<u>Total Jet Mass (g)</u>		<u>Jet Kinetic Energy (kJ)</u>	
	<u>Fill</u>		<u>BRLSC</u>	<u>BASC</u>	<u>BRLSC</u>	<u>BASC</u>
1	TNT	Light	19.66	28.42	196.	385.
2	Comp B	Light	22.56	30.73	287.	507.
3	Octol	Light	27.40	32.03	349.	570.
4	PBX	Light	27.97	33.49	381.	614.
5	Comp B	Heavy	33.63	60.41	475.	1262.
6	Octol	Heavy	35.17	60.16	360.	1406.

<sup>20</sup>R.T. Sedgwick, L.J. Walsh and M.S. Chawla, "Effects of High Explosive Parameters and Degree of Confinement on Jets from Lined Shaped Charges," BRL CR 245, Prepared by System, Science and Software, July 1975. (AD #B006987L)



<u>Calculation Number</u>	<u>Predicted Final Jet Tip Velocity (km/s)</u>		<u>Measured Jet Tip Velocity (km/s)</u>
	<u>BRLSC</u>	<u>BASC</u>	
1	5.95	6.79	6.8
2	7.15	7.74	7.7
3	7.40	8.17	8.1
4	7.80	8.44	8.3
5	7.44	8.86	---
6	7.78	9.36	---

In these four calculations, the charge was confined in an aluminum casing (treated as being unconfined in the BASC code). Calculations 5 and 6 employed Comp B and Octol as explosive fills which were confined in a steel casing with a thickness of 10mm. The calculational matrix is shown in Table III with the degree of confinement provided by the aluminum and steel casing is referred to as light and heavy, respectively. The summary of the results of both codes (BASC and BRLSC) as well as the results of the jet tip velocities measured from experiments<sup>21</sup> are also shown in Table III. The quantities summarized represent the total jet which is composed of jet material having a velocity greater than or equal to a velocity of 2.8 km/s. The results indicate the ratio of predicted total jet mass of BRLSC code to the BASC code is approximately 65% for all calculations.

The comparison of jet tip velocity data is illustrated on Figure 13. The theoretically predicted values for calculations 1 through 4 from both the BASC and BRLSC codes are compared with the experimental determined jet tip velocity data. Figure 13 is a plot of calculated jet velocities versus the measured jet tip velocities. The open triangles are the final jet tip velocities as predicted by the BRLSC code<sup>20</sup> and the open squares are those predicted by the BASC code. There is very good agreement between the BASC code and the observed results for those rounds considered and the BRLSC code's agreement is less than 10% except for TNT which is 12.5%.

In summary, we have demonstrated in these comparisons the predictive ability of the BASC code. In addition, several of its salient features has been shown. These features include the predictive characteristics of BASC with respect to variations in casing confinement, variations in cone apex angle, variations in liner density and thickness, variations in explosive fill, and the effects of scaling. We have, also, demonstrated its predictive ability with respect to other sophisticated numerical schemes.

<sup>21</sup>J. Simon, "The Effect of Explosive Detonation Characteristics on Shaped-Charge Performance," BRLMR-2414, August 1979. (AD #B000337L)

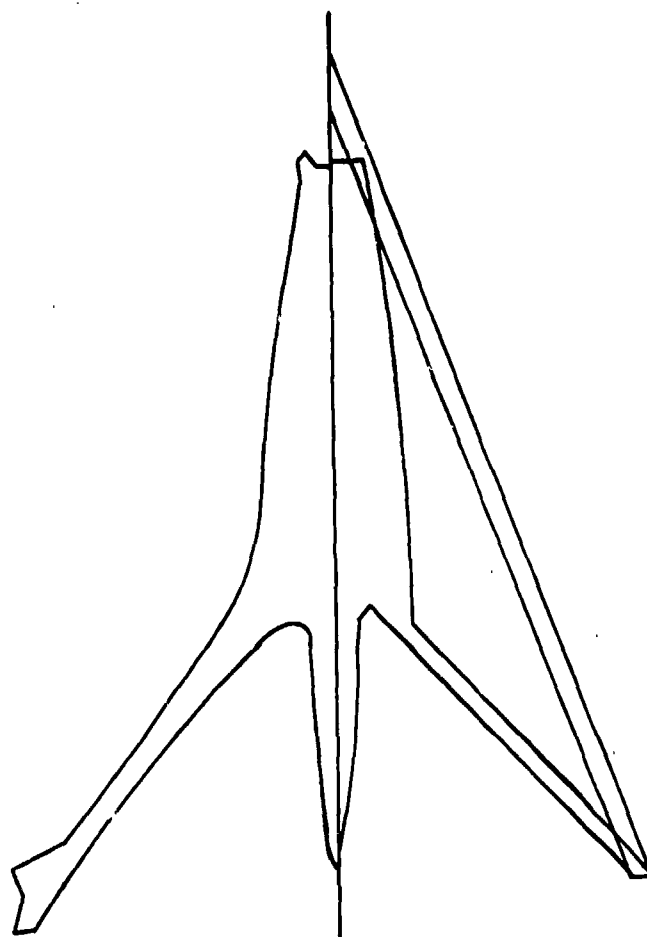


Figure 12. Comparison of calculations of the flow field between BASC code and a 2-D, Hydrodynamic Code (BRLSC) (Reference 2). These calculations are of the 105-mm, unconfined, shaped-charge shown at the time approximately 25 $\mu$ s after the initiation of the explosive. The BRLSC Code results are to the left of the axis of symmetry and the BASC code results are to the right.

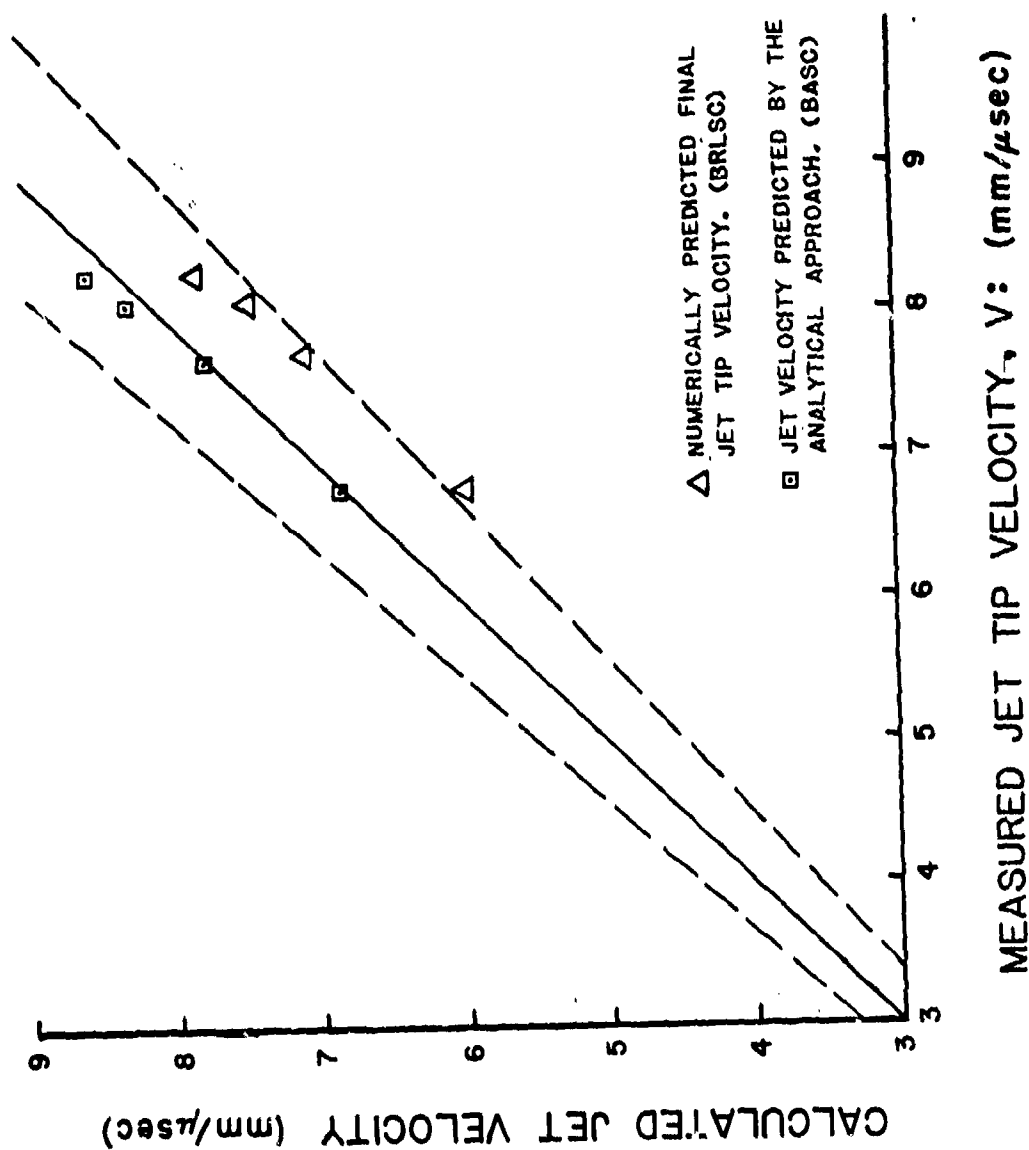


Figure 13. Comparison of Calculated and Measured Jet Tip Velocities from Shaped-Charge Designs 1 through 4, Shown on Table III.

Based upon these comparisons, it is concluded that, for many conical shaped-charge designs, the BASC code is a useful, predictive model for jet characteristics. Since it requires only fractions of a minute to do several shaped-charge design calculations on most computing systems (large or small), this calculational scheme is an economical approach for warhead design and analysis.

## V. LIMITATIONS OF BASC

The current version of BASC requires several empirically determined constants. The collapse model requires the determination of  $\phi$  and  $k$ ; their relationship to the angle of incidence,  $i$ ; the constant  $A$  for a confined casing; and finally,  $C_1$ , a constant calibrated to the 105-mm, unconfined, shaped charge for the time of collapse near the apex of the cone. Penetration requires a cutoff velocity,  $U_{min}$ , and a breakup time,  $T_1$ .

Breakup time, as used in the theory of DiPersio, et. al.<sup>7</sup>, is a quantity determined from a radiograph of the jet. It is treated as a constant, but only approximates that value for conical liners with uniform wall thicknesses. Work by Simon<sup>21</sup> shows that breakup time for a given geometry remains a constant over changes in the Chapman-Jouguet pressure of 186-380 kbars. The codes, as used, requires a predetermined value for  $T_1$  for a given shaped-charge geometry. The value of  $T_1$  scales with cone diameter. All of these limitations stem from the empirical nature of the equations of the code.

Geometrical considerations for conical collapse may or may not be a limitation to the BASC code. Research will have to be conducted in order to determine this fact. Since the initial cone half-angle varies as a function of distance along the axis,  $z$ , and since each zonal element has no interaction with other elements, all geometrical considerations should be solvable.

Current limitations will be corrected by future changes to the code. These changes are explained in the next section.

## VI. FUTURE MODIFICATIONS

There are several tasks intended to improve BASC. The major areas of research are the following: (1) modeling the jet formation for

several different materials and geometrical consideration, (2) modeling the viscous effects, as suggested by Walters,<sup>22</sup> in a nature compatible with BASC, (3) modeling the breakup of the jet with reference to a maximum strain to break, (4) modeling the threshold for a jet-no-jet criteria for many more materials.

As explained earlier, elements in the apex region of a conical liner reach the jet formation region before these elements are accelerated to their maximum attainable velocity. To account for this characteristic of conical, liner apex flow, which is so important in determining exact jet-tip velocities, transient acceleration is being modeled by an empirical constant determined from normalized, copper liner data<sup>13</sup>. This constant accounts for the number of shock reverberations that occur in the copper liner before it enters the flow of jet formation. We can improve upon this constant by obtaining experimental collapse data for a number of different materials.

Research conducted by Simon indicates that the breakup of the shaped-charge jet may be related to a maximum strain, which is a function of strain rate. These observations were made for copper liners with many different explosive fillers ranging from TNT (C-J pressure 186 kbars) to a high HMX explosive (C-J pressure 380 kbars), but with only one charge and liner geometry. Work is continuing by Chou, et. al.,<sup>23,24</sup> to define the critical condition for breakup models based on these results and will be added to the code where piecewise strains and strain rates will be calculated. We will continue to calculate penetration velocity,  $U$ , and will terminate penetration according to a minimum value of  $U$  as demonstrated by DiPersio, et. al.; but we will explore the use of  $v_{jmin}$  as the penetration cut-off criteria. Finite-difference codes may be applied either directly or in a simplified form to generate a library of parameters by computational experiments. This will assist in the research of some of these critical parameters utilized in the BASC code.

As certain elements of the shaped-charge collapse problem are accurately modelled, that section of the BASC code will be modified. The code is considered a "living" code, constantly being updated but applied within its limitations at all stages of its development.

<sup>22</sup>W. Walters, "Influence of Material Viscosity on the Theory of Shaped-Charge Penetration," ARBRL-MR-02041, August 1979. (AD #B041486L)

<sup>23</sup>P.C. Chou and J. Carleone, "Calculations of Shaped-Charge Jet Strain, Radius and Breakup Time," BRLCR-246, July 1976. (AD #B007240L)

<sup>24</sup>J. Carleone, P.C. Chou, and R. Ciccarelli, "Shaped-Charge Jet Stability and Penetration Calculations," BRLCR-351, September 1977. (AD #A050117)

## VII. SUMMARY

BASC is a simple, well-documented, shaped-charge code that may be applied to many problems to predict trends and gross effects. Difficulties in predicting the jet-tip characteristics still exist for some materials, but future modifications should correct this deficiency. The code is so structured that it can grow and become more widely applicable as modeling improvements are available.

## ACKNOWLEDGEMENT

The author is grateful to Robert DiPersio, now retired, and Robert R. Karpp, now of Los Alamos Scientific Laboratory, for their many contributions in the formulation of the BASC code.

The author is grateful to Mary Scarborough of the BRL for her artistic work in the preparation of the illustrations used in this report.

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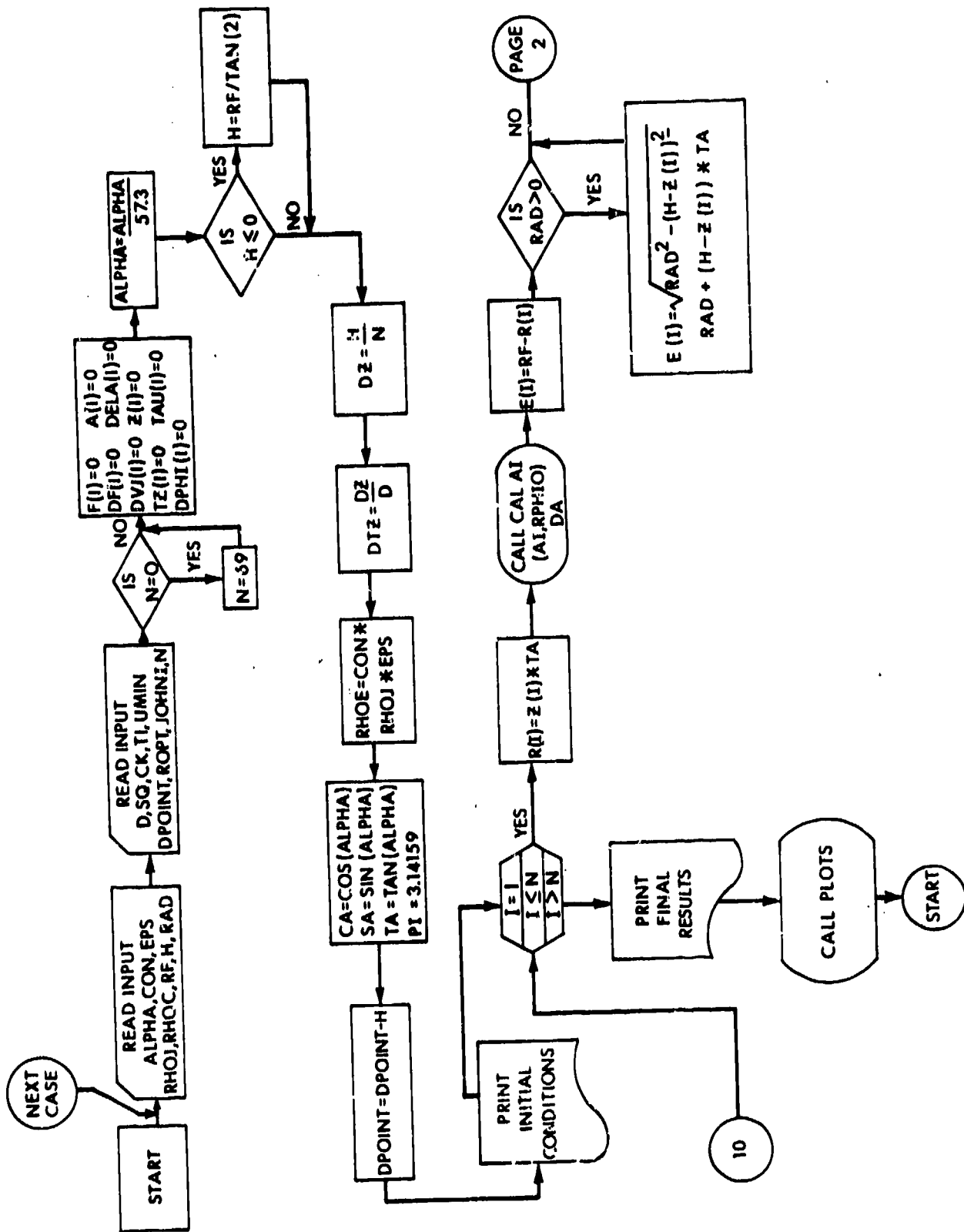
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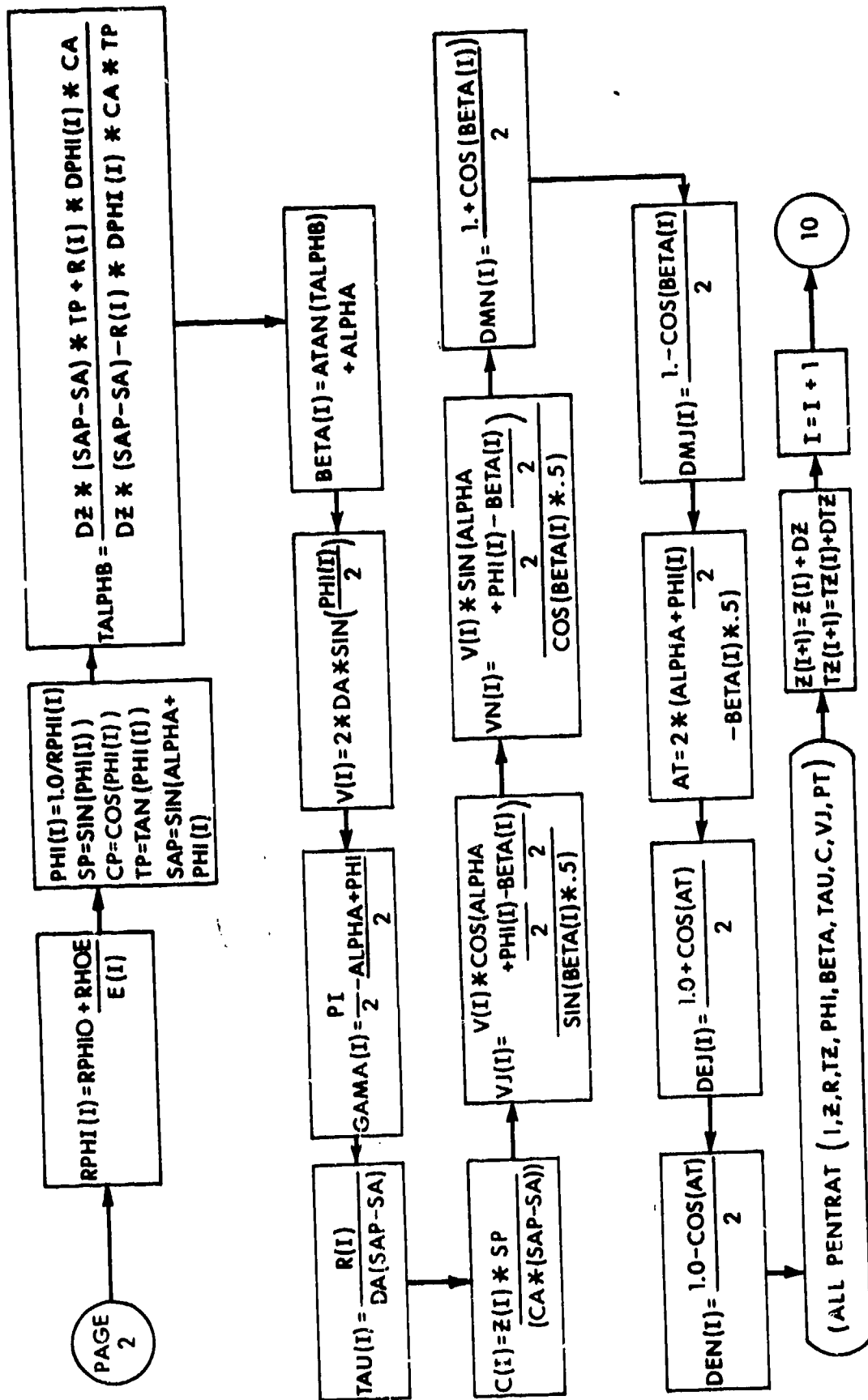
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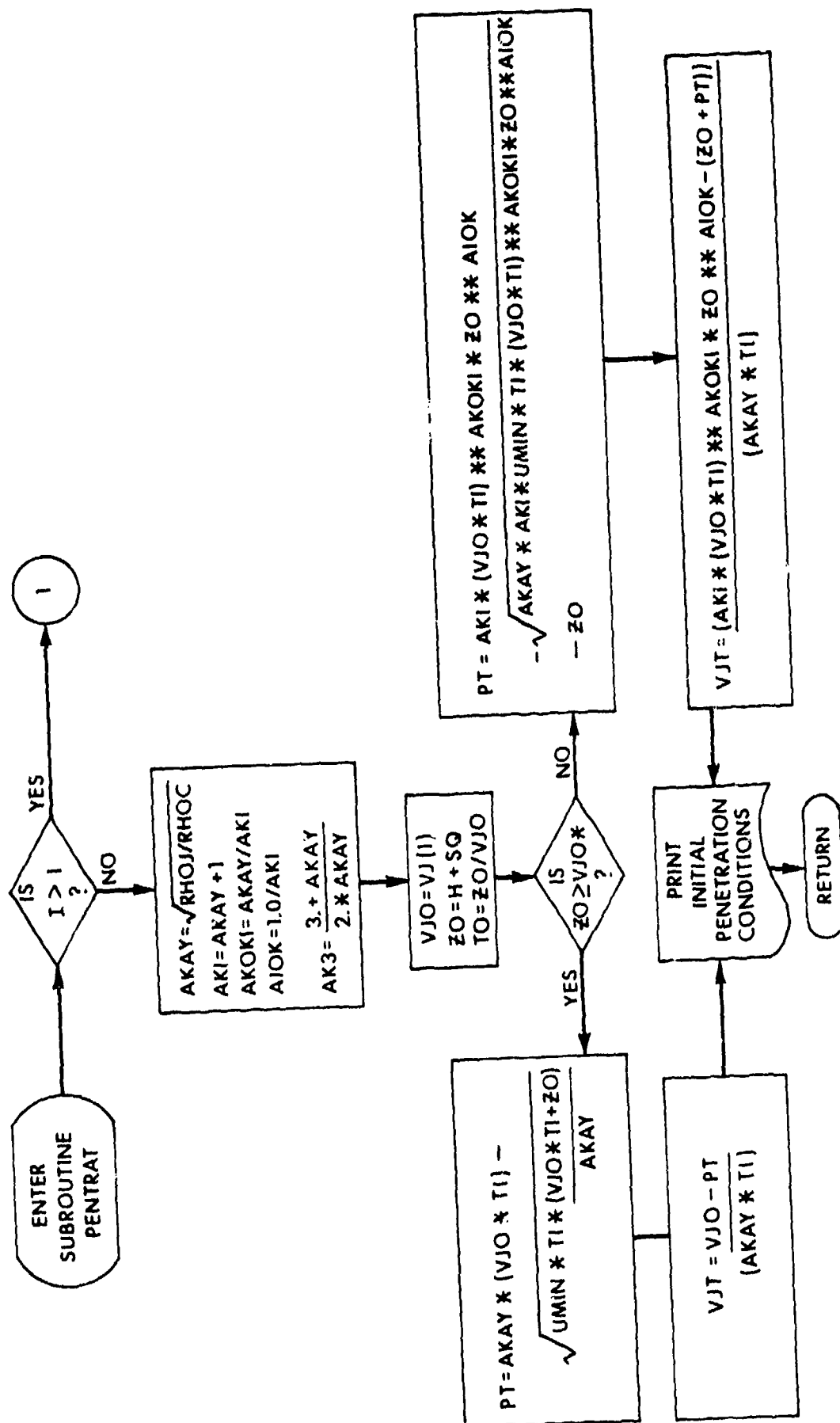
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24. J. Carleone, P.C. Chou, and R. Ciccarelli, "Shaped-Charge Jet Stability and Penetration Calculations," BRLCR-351, September 1977. (AD #A050117)

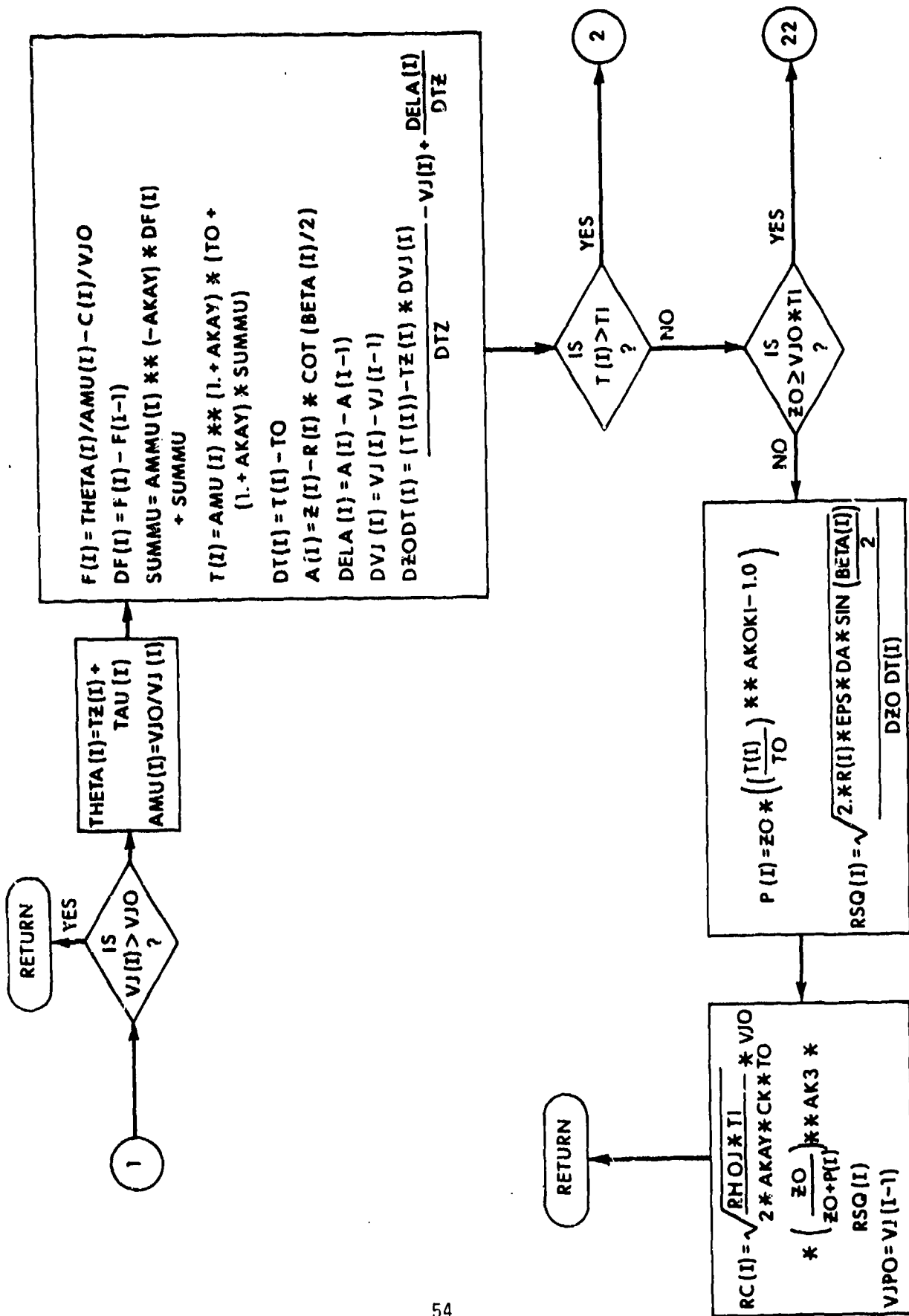


APPENDIX A  
FLOWCHART OF COMPUTER PROGRAM









22

$$P(I) = \frac{VJO * (T(I) - TO * TI)}{TI + T(I)} \cdot \frac{AKAY}{AKAY}$$

$$VJP = VJO * \left( TI + \frac{TO}{AKAY} \right) \cdot \frac{TI + T(I)}{AKAY}$$

$$DVJP = VJP - VJPO$$

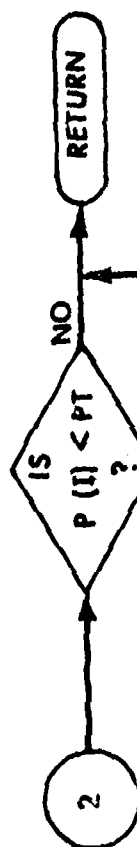
$$VJPO = VJP$$

$$DZODT(I) = (TI - T(I)) * DVJP - VJP + DELA(I)$$

$$RSQ(I) = \sqrt{\frac{2 * R(I) * EPS * DA * \sin\left(\frac{BETA(I)}{2}\right) ** 2}{DZODT(I)}}$$

$$RC(I) = \sqrt{\frac{RHOJ}{2 * AKAY * CK}} * RSQ(I) * VJO$$

RETURN

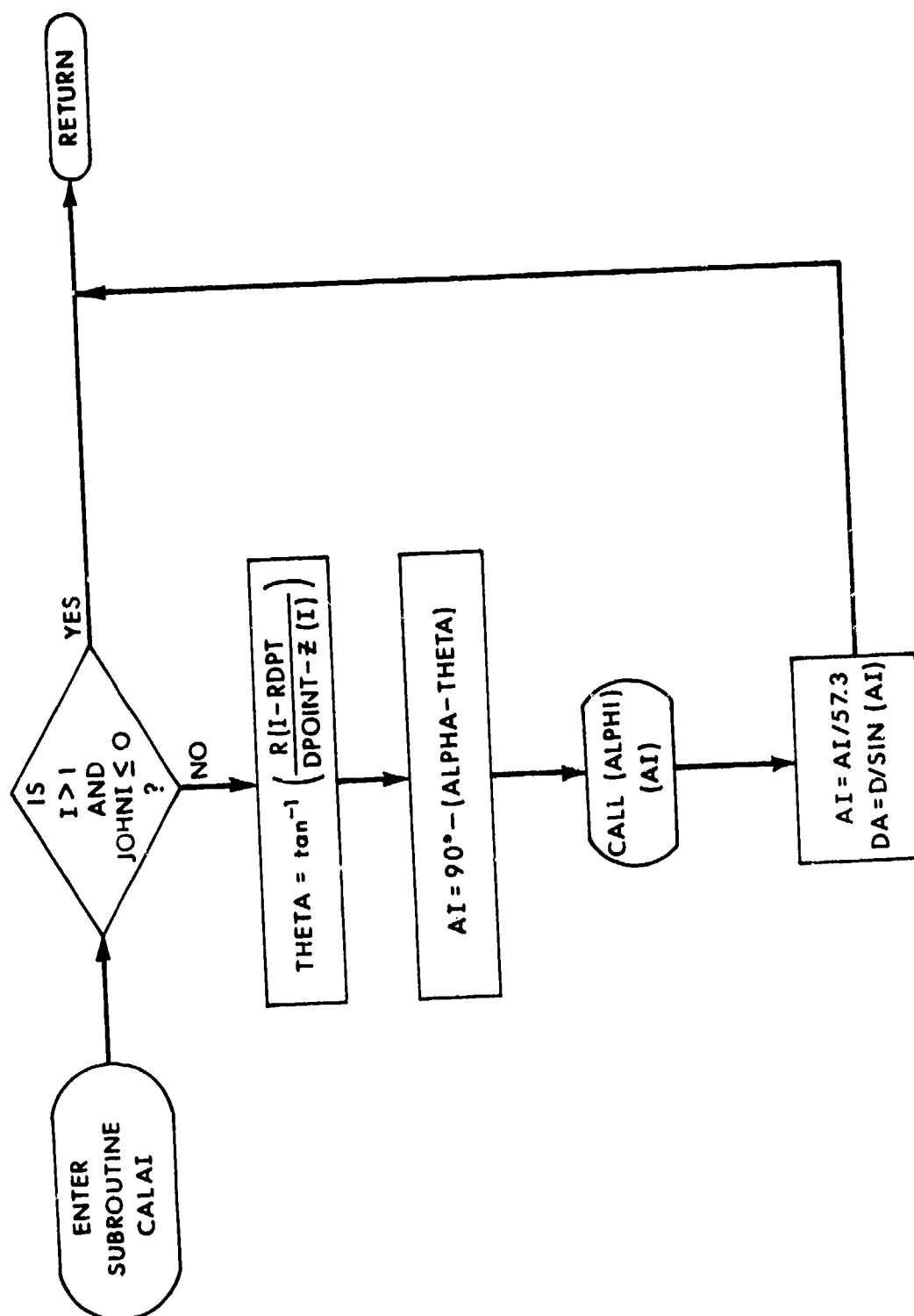


$$P(I) = \left( \frac{AKI * (T_I / T_O)}{T(I) + AKAY * T_I} - 1.0 \right) * ZO$$

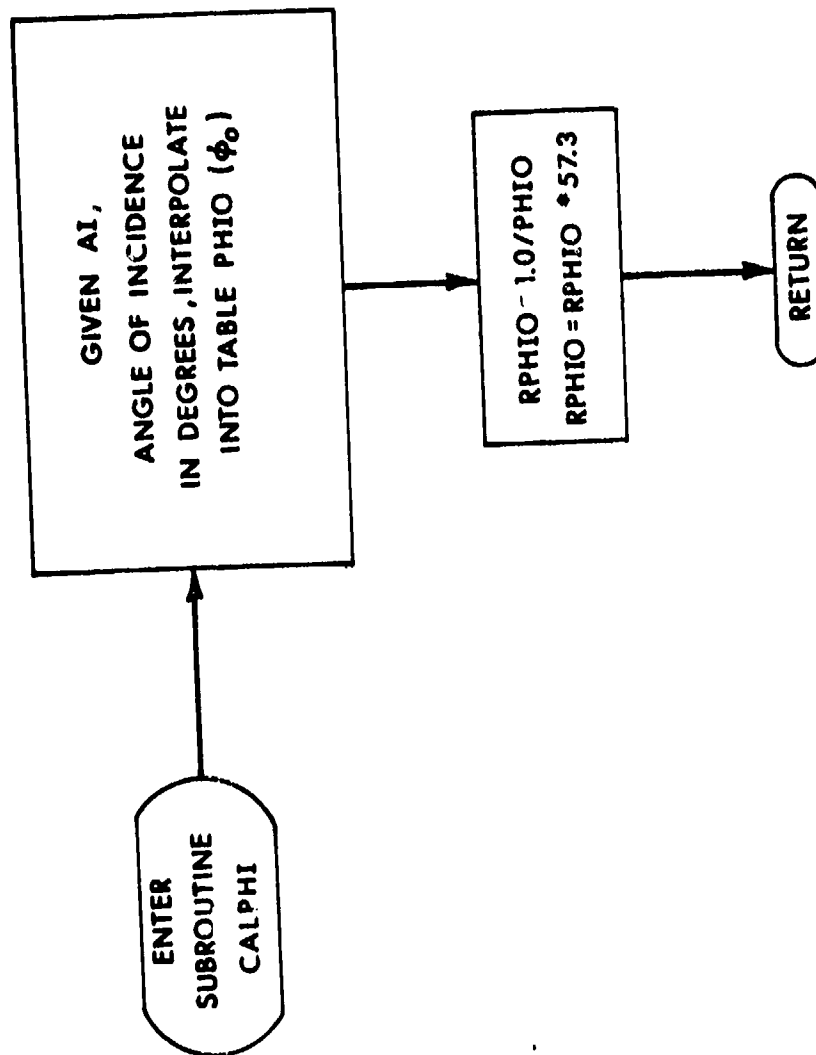
$$RSQ(I) = \sqrt{\frac{2 * R(I) * EPS * DA * \sin\left(\frac{BETA}{2}\right) ** 2}{DZODT(I)}}$$

$$RC(I) = \sqrt{\frac{RHOJ}{2 * AKAY * CK} * \frac{VJO}{AKAY} * (AKI *}$$

$$\left( \frac{ZO}{VJO * T_I} \right) ** AIOK - \frac{(ZO + P(I))}{VJO * T_I} * RSQ(I)$$







APPENDIX B  
FORTRAN IV  
COMPUTER PROGRAM LISTING, SAMPLE INPUT AND OUTPUT

C	PROGRAM RASC(INPUT,OUTPUT,TAPES=INPUT,TAPE6=OUTPUT)	MAIN	3
C	BRL ANALYTICAL SHAPED CHARGE (RASC) CODE	MAIN	4
C		MAIN	5
C	PROGRAMMED BY JOHN T. HARRISON	MAIN	6
C		MAIN	7
C		MAIN	8
C	BRL ANALYTICAL SHAPED CHARGE CODE	MAIN	9
C		MAIN	10
C	WRITTEN IN STANDARD FORTRAN IV	MAIN	11
C	SEMI-EMPIRICAL CODE BASED ON THE FORMULA-	MAIN	12
C		MAIN	13
C	$1/PHI=1/PHI0+K*RHOJ*EPS/E$	MAIN	14
C		MAIN	15
C	UNITS FOR THIS CODE ARE CENTAMETERS,GRAMS,AND MICROSECONDS.	MAIN	16
C		MAIN	17
C	THE CODE WILL CALCULATE COLLAPSE VELOCITY OF THE LINER AND	MAIN	18
C	VELOCITIES,MASSSES,AND ENERGIES OF BOTH THE JET AND SLUG.	MAIN	19
C		MAIN	20
C	THIS CODE WILL PREDICT PENETRATION AND HOLE PROFILE AND GIVE	MAIN	21
C	PENETRATION-STANDOFF CURVES AND TABLES.	MAIN	22
C		MAIN	23
C	LIST OF VARIABLES	MAIN	24
C		MAIN	25
C	ALRAD = ALPHA IN RADIANS	MAIN	26
C	CON = DETONATION PRODUCTS CONSTANT (K).	MAIN	27
C	DZ = DELTA Z DISTANCE ALONG THE LENGTH OF THE LINER	MAIN	28
C	DT7 = TIME THE DETONATION WAVE ARRIVES AT EACH Z POINT	MAIN	29
C	AI = ANGLE OF INCIDENCE	MAIN	30
C	Z(I) = DISTANCE ALONG LINER	MAIN	31
C	TZ(I) = TIME AT EACH POINT ON LINER	MAIN	32
C	GAMA(I) = ANGLE BETWEEN AXIS AND COLLAPSE DIRECTION.	MAIN	33
C	BETA(I) = COLLAPSE ANGLE	MAIN	34
C	V(I) = COLLAPSE VELOCITY	MAIN	35
C	TAU(I) = TIME FOR EACH ELEMENT TO REACH THE AXIS	MAIN	36
C	PHI(I) = ANGLE OF LINER BENDING.	MAIN	37
C	PHI0 = ANGLE OF EXPLOSIVE BENDING.	MAIN	38
C	RPHI0 = $1/PHI0$	MAIN	39
C	RPHI = $1/PHI$	MAIN	40
C	C(I) = POINT ON THE AXIS WHERE EACH PARTICLE WILL HIT	MAIN	41
C	VJ(I) = VELOCITY OF THE JET	MAIN	42
C	VN(I) = VELOCITY OF THE SLUG	MAIN	43
C	CO = SOUND SPEED OF LINER MATERIAL (COPPER CO=395 CM/MSEC)	MAIN	44
C	RV(I) = FLOW VELOCITY (VF). DIRECTED INTO STAGNATION POINT.	MAIN	44A
C	DFJ(I) = DELTA ENERGY IN EACH JET SEGMENT	MAIN	44P
C	DFN(I) = DELTA ENERGY IN EACH SLUG SEGMENT	MAIN	45
C	R(I) = RADIUS OF THE CONE AT EACH DELTA Z DISTANCE	MAIN	46
C	DA = APARENT DETONATION RATE.	MAIN	47
C	DML(I) = DELTA MASS OF THE LINER	MAIN	48
C	E(I) = AMOUNT OF EXPLOSIVE ABOVE EACH POSITION ON THE LINER	MAIN	49
C	DPOINT = INITIALLY IS THE TOTAL HEIGHT OF THE CHARGE. THEN BE-	MAIN	50
C	COMES THE INITIATION POINT ON THE Z DIRECTION.	MAIN	51
C	NRAD = RADIUS OF CURVATURE OF THE CONFINMENT.	MAIN	52
C	IF NRAD=0.-THEN THERE WILL BE A LINEAR THICKNESS OF EXPL	MAIN	53
C		MAIN	54
C	INPUT PARAMETERS	MAIN	55
C		MAIN	56

```

C
C
C CARD NUMBER 1
C
C
C IJOHN = 1, A PARAMETRIC STUDY, 2, A NEW CASE.
C
C HEAD = A HEADING CARD.
C
C
C CARD NUMBER 2
C
C ALPHA = HALF ANGLE OF THE CONE IN DEGREES.
C EPS = THICKNESS OF THE LINER, (0 IF UNKNOWN OR VARIABLE)
C RHOJ = DENSITY OF THE LINER
C RF = FINAL RADIUS OF THE CONE
C H = HEIGHT OF THE CONE. IF H=0., THEN H WILL BE COMPUTED.
C COF = CONFINEMENT FACTOR (0 FOR UNCONFINED, IF CONFINED = THICKNESS OF LINER)
C RHOCON = DENSITY OF THE CONFINEMENT
C NPLT = NUMBER OF PLOTS, 0=SKIP, 1=ALL, 2=VEL. + PENETRATION
C NPOS = NUMBER OF (R,Z) POSITIONS TO BE READ IN.
C IF NPOS IS 0 DO NOT READ IN (R,Z) COORDINATES.
C
C CARD NUMBER 3
C
C RHOC = DENSITY OF THE TARGET
C SO = STAND OFF DISTANCE
C IF SO= 0. PENETRATION STANDOFF CURVES WILL BE PLOTTED.
C CK = CONSTANT FOR DETERMINING HOLE VOLUME
C IF (CK =0.) CK WILL BE CALCULATED AND UMIN WILL THEN
C BECOME THE BMN TO BE READ IN.
C UMIN = VELOCITY MIN. USED IN THE PENETRATION THEORY
C IF (CK =0.) UMIN WILL BE THE BMN.
C T1 = BREAKUP TIME OF THE JET
C IF (T1 = 0.) T1 WILL BE CALCULATED.
C DPOINT = INITIALLY IS THE TOTAL HEIGHT OF THE CHARGE
C RDPT = RADIUS ABOVE AXIS WHERE EXPLOSIVE IS INITIATED.
C NEXPL = EQUATION OF STATE NUMBER FOR THE EXPLOSIVE.
C
C NEXPL EXPLOSIVE DENSITY DETONATION RATE
C 1 COMP B 1.72 .8 MAIN 94
C 2 OCTOL 1.82 .85 MAIN 95
C 3 LX 1.84 .8774 MAIN 96
C
C N = THE NUMBER OF ZONES Z-AXIS IS TO SUBDIVIDED. IF N=0 THE
C DEFAULT VALUE IS N=69.
C
C
C DIMENSION EXPLO(6), RHOM(6), DV(6)
C DIMENSION RI(100), IMAT(100)
C DIMENSION EPSI(100), DAI(100)
C DIMENSION RP(300), PLOT(3)
C DIMENSION PSO(50), PPT(50)
C DIMENSION CONF(4)
C DIMENSION UJK(200), DMJK(200)
C DIMENSION HEAD(6)
C DIMENSION DML(100), RV(100), DEL(100)
C DIMENSION Z(100), T2(100), TAU(100), E(100), PHI(100), DPHI(100),
1 BETA(100), GAMA(100), R(100), C(100), VJ(100), VN(100), DMJ(100)
2 , DMN(100), DEJ(100), DEN(100), RPHI(100), V(100)

```

C		MAIN114
C	COMMON	MAIN115
C	COMMON ALRAD, EPS, RHOJ, RHOC, RF, RPHI, DTZ, SO, CK, DA, H, U,	MAIN116
	1 PT, DEGTOR	MAIN117
	COMMON AMU(100), THETA(100), F(100), DF(100), T(100), DT(100),	MAIN118
	1 G(100), P(100), A(100), DELA(100), DVJ(100), DZOLT(100),	MAIN119
	2 RSO(100), RC(100), T1, UMIN	MAIN120
	COMMON /LPLOT/ Z, TZ, TAU, E, PH1, DPH1, BETA, GAMA, R, C, VJ, VN,	MAIN121
	1 DMJ, DMN, DEJ, DEN, RPH1, V	MAIN122
C		MAIN123
	DATA HEAD(6) /1H>/	MAIN124
	DATA CONF(1), CONF(2) /10MUNCONFINED, 1H>/	MAIN125D
	DATA CONF(3), CONF(4) /10MCONFINED, 1H>/	MAIN126D
	DATA PLOT(1), PLOT(2), PLOT(3) /4HSKIP, 3HALL, 9HVEL + FEN/	MAIN127D
	DATA EXPLO(1) /6HCOMP R/ , EXPLO(2) /5HOCTOL/	MAIN129D
	DATA EXPLO(3) /4H TNT/ , EXPLO(4) /8HPEX=9404/ , EXPLO(5) /4HLX14/	MAIN128D
	DATA EXPLO(6) /4H0280/	MAIN130D
	DATA RHOM(1), RHOM(2) /1.72, 1.82/	MAIN130E
	DATA RHOM(3), RHOM(4) /1.63, 1.84/ , RHOM(5) /1.81/ , RHOM(6) /1.74/	MAIN131D
	DATA DV(1), DV(2), DV(3), DV(4) / .794, .848, .690, .880/	MAIN132D
	DATA DV(5), DV(6) / .8749, .8477/	MAIN133D
C	*****	MAIN134D
C	READ IN HEADING ONE TIME FOR EACH PARAMETRIC STUDY	MAIN135
C		MAIN136
C	1 CONTINUE	MAIN137
	HEAD (5,37) IJOHN,HEAD	MAIN138
	IF (EOF(5)) 9A765,2	MAIN139H
	2 WRITE (6,36) IJOHN,HEAD	MAIN139A
C		MAIN140H
C	BEGIN LOOP FOR A PARAMETRIC STUDY	MAIN141
C		MAIN142
C		MAIN143
C		MAIN145
C	READ IN INPUT ( 2 CARDS. BOTH WITH 7E10.3,215 FORMAT.)	MAIN146
C		MAIN147
C		MAIN148
C	READ IN LINER AND CONFINEMENT INPUTS	MAIN149
C		MAIN150
C	READ (5,36) ALPHA,EPS,RHOJ,RF,H,COF,RHOCON,NPLT,N1(-5	MAIN151R
C		MAIN152
C	READ IN TARGET AND EXPLOSIVE INPUTS	MAIN153
C		MAIN154
C	READ (5,36) RHOC,SO,CK,UMIN,T1,DPOINT,RDPT,NEXPL,A	MAIN155R
	SH=M	MAIN156
	CO=.395	MAIN156A
	TMJJ=0.	MAIN156H
	JOHN1=1	MAIN157
	TECUT1=0.	MAIN157A
	TECUT2=0.	MAIN157B
	TECUT3=0.	MAIN157C
	TECUT4=0.	MAIN157D
	TECUT5=0.	MAIN157E
	TMJCUT=0.	MAIN157F
	TMCL1=0.	MAIN157G
	TMCL2=0.	MAIN157H
	TMCL3=0.	MAIN157I
	TMCL4=0.	MAIN157J

	TMCUTS=0.	MAIN157A
	RVMAX=0.	MAIN157L
	NRAD=0	MAIN158
	N=OV(NEXPL)	MAIN159
	ICOF=1	MAIN160
	IF (COF.GT.0.) ICOF=3	MAIN161
	CON=.78	MAIN162
	IS=2	MAIN163
	IF (N.EQ.0) N=69	MAIN164
	IERROR=0	MAIN165
	HAD=FLOAT(NRAD)	MAIN166
	VJO=0.	MAIN167
	TWLL=0.	MAIN167A
	TMJ=0	MAIN168
	TEJ=0.	MAIN169
	TKF=0.	MAIN170
	TMN=0.	MAIN171
	TML=0.	MAIN172
	SUMVJ=0.	MAIN173
	VJMAX=0.	MAIN174
	UJKMAX=0.	MAIN175
	TMASS=0.0	MAIN176
	DMASS=0.	MAIN177
	DMJTE=0.	MAIN178
	DO 3 J=1,1400	MAIN179
	Z(J)=0.	MAIN180
	3 AMU(J)=0.	MAIN181
C		MAIN182
C	CONVERT ANGLE ALPHA FROM DEGREES TO RADIANS.	MAIN183
C		MAIN184
	PI=3.141592653589793	MAIN185
	DEGTOR=57.2957795131	MAIN186
	ALRAD=ALPHA/DEGTOR	MAIN187
	TRRF=YAN(.5*ALRAD)	MAIN187A
C	COMPUTE THE HEIGHT OF THE LINER IF NO GIVEN IN INPUT.	MAIN188
C		MAIN189
	IF (H.LE.0.) H=RF/TAN(ALRAD)	MAIN190
C		MAIN191
	DZ=H/FLOAT(N)	MAIN192
	DTZ=DZ/D	MAIN193
C	SFT UP CONSTANTS TO BE IN LATER EQUATIONS.	MAIN194
	RHOE=CON*RHOJ*EPS	MAIN195
	SA=SIN(ALRAD)	MAIN196
	CA=COS(ALRAD)	MAIN197
	TA=TAN(ALRAD)	MAIN198
	IF (PF.LE.0.) RF=TA*H	MAIN199
	PJPS=PI/3.	MAIN200
	PIP =2.0*PI	MAIN200A
	SIGMA=RHOJ*EPS	MAIN201
C		MAIN202
C	DETONATION POINT IS THE HEIGHT OF THE CHARGE - HEIGHT OF THE CONE.	MAIN203
C		MAIN204
C	DPOINT=DPOINT-H	MAIN205
	IF (DPOINT. < 0.) DPOINT=2.*RF	MAIN205A
	IF (COF.GT.0.) CONFTE=COF	MAIN206
	DELCOF=CONFTE/20.	MAIN207
C		MAIN208

C	COF = CONFINEMENT FACTOR (INCREASE IN EXPLOSIVE THICKNESS).	MAIN204
C		MAIN210
	COF=1.+.64*RHOCN*CONFTH*8.9+.177R/(7.R*.60198*RHOU*EPS)*2.4003*HF	MAIN211
	1 / (DPCINT*2.0553)	MAIN212
	WRITE (6,28) ALPHA,CON,EPS,RHOU,CONF(ICOF),CONFTH,RHOCN,COF,H,RF	MAIN213
	1 NZ,DTZ,N	MAIN214
	WRITE (6,29) EXPLO(NEXPL),D,RHOHE(NEXPL),DPOINT,HDPT,RHUC,SO,	MAIN215
	1 (WIN,PLOT(NPLT+1))	MAIN216
	RRF=PF	MAIN217
	PHI(1)=3.0/DEGTOR	MAIN218
C	INITIAL TIME OF COLLAPSE TW=TAU.	MAIN219
	TW=.1	MAIN220
C	CONTANT FOR THE ACCELERATION ROUTINE.	MAIN221
	C1=2.+.CO*EPS*RHOU/(.392+.269*8.9)	MAIN221A
	FTERM=.000001	MAIN222
C		MAIN223
C	SETUP ROUTINE	MAIN224
C		MAIN225
C	SETUP INITIAL POSITIONS (R,Z)	MAIN226
C		MAIN227
C	RI= INSIDE RADIUS OF THE CONE. R = OUTSIDE RADIUS OF THE CONE.	MAIN228
C		MAIN229
	X=EPS/CA	MAIN230
C		MAIN231
C	THIS IS A SETUP ROUTINE FOR A SHAPED CHARGE LINEAR WITH CONSTANT EPS	MAIN232
C		MAIN233
	IF (NPOS.EQ.0) NPOS=1	MAIN234
	NPOS=NPOS+1	MAIN235
	WRITE (6,27)	MAIN236
	IF (NPOS.GE.N) GO TO 5	MAIN237
	DO 4 J=NPOS,N	MAIN238
	Z(J)=Z(J-1)+DZ	MAIN239
	IMAT(J)=1	MAIN239A
	R(J)=Z(J)*TA	MAIN240
	RI(J)=R(J)-X	MAIN241
	IF (RI(J).LE.0.) RI(J)=0.	MAIN242
	IF (RI(J).LE.0.) IS=J+1	MAIN242A
	IF (ABS(RI(J)-0.0).LE..0001) RI(J)=0.0	MAIN242B
	4 CONTINUE	MAIN243
	5 CONTINUE	MAIN244
	WRITE (6,31)	MAIN245
	WRITE (6,34) (I,Z(I),I=1,N)	MAIN246
	WRITE (6,32)	MAIN247
	WRITE (6,34) (I,R(I),I=1,N)	MAIN248
	WRITE (6,33)	MAIN249
	WRITE (6,34) (I,RI(I),I=1,N)	MAIN250
C		MAIN251
C		MAIN252
C	E D OF SETUP ROUTINE	MAIN253
C		MAIN254
	TZ(1)=DPOINT/D	MAIN255
	EPSI(1)=EPS	MAIN255A
C		MAIN256
C	START SPACE + TIME ITERATION.	MAIN257
C		MAIN258
	DO 10 I=2,N	MAIN259
	V(I)=0.	MAIN260

	VN(I)=0.	MAIN261
	VJ(I)=0.	MAIN262
	UJK(I)=0.	MAIN263
C	FIND HALF APEX ANGLE-ALPHA	MAIN266
	DZ=Z(I)-Z(I-1)	MAIN267
	DR=R(I)-R(I-1)	MAIN268
	ALRAD=ATAN2(DR,DZ)	MAIN269
	SA=SIN(ALRAD)	MAIN270
	CA=COS(ALRAD)	MAIN271
	TA=TAN(ALRAD)	MAIN272
	IF (IPAT(I).EQ.1) CO=.392	MAIN273
	IF (IHAT(I).EQ.1) KHOJ=8.9	MAIN273A
	EPS=(R(I)-RI(I))*CA	MAIN274
C		MAIN275
C	GO TO SUBROUTINE CALAI TO CALCULATE INCLINATION ANGLE (AI),	MAIN276
C	(DA)-DETONATION COMPONENT TO VELOCITY, AND (RPHIO)=1./PHIO.	MAIN277
C		MAIN278
	CALL CALAI (I,Z,R,DPOINT,RDPT,JOHNI,AI)	MAIN279
	DAI(I)=DA	MAIN280
C		MAIN281
C	COF IS FACTOR FOR CONFINEMENT.	MAIN282
C		MAIN283
C		MAIN284
C	IF RAD IS GIVEN NO EQ. 0. CLACULATE CURVED THICKNESS OF EXPLOSIVE.	MAIN285
C		MAIN286
C	IF (RAD.GT.0.) E(I)=SQRT(ABS(RAD**2-(H-Z(I))*2))-RAD*(H-Z(I))*TA	MAIN287
C		MAIN288
C	CALCULATE EXPLOSIVE THICKNESS	MAIN289
C	E(I)=COF*(RF-R(I))*(COF-1.)*RHOCON*CONF TK/7.8	MAIN290
C	CALCULATE 1./PHI	MAIN291
C		MAIN292
	TZ(I)=Z(I)/(DA*CA)	MAIN293
	TZ(I)=TZ(I)+TZ(I)	MAIN294
	THE90=.5*PI-AI	MAIN295
	THE=ALRAD-THE90	MAIN296
	EPSI(I)=EPS/COS(THE90)	MAIN297
	E(I)=E(I)/COS(THE)	MAIN298
	RHOE=CON*RHOJ*EPSI(I)	MAIN299
	CONE=RHOE*CI	MAIN300
	RPHI(I)=RPHIO+RHOE/E(I)	MAIN301
C		MAIN302
C	SUBROUTINE TO CALCULATE THE ACCELERATION OF THE LINER.	MAIN303
C		MAIN304
C	ROUTINE WILL ITERATE ON TAU AND PHI.	MAIN305
C		MAIN306
	DO 6 J=1,300	MAIN307
	CT=EXP((SQRT(CONE/(E(I)*TW))))	MAIN308
	RP(J)=RPHI(I)*CT	MAIN309
	TW=R(I)/(DA*(SIN(ALRAD+1.0/RP(J))-SA))	MAIN310
	IF (J.EQ.1) GO TO 6	MAIN311
	IF (ABS(RP(J)-RP(J-1)).LE.ETERM) GO TO 7	MAIN312
	6 CONTINUE	MAIN315
C		MAIN316
C	END ITERATIVE PROCESS	MAIN317
C		MAIN318
	J=J-1	MAIN319
	7 CONTINUE	MAIN320



C		MAIN321
C	THE END OF THE INVERSE VELOCITY GRADIENT SUBROUTINE.	MAIN322
C		MAIN323
	TAU(I)=TW	MAIN324
	RPHI(I)=RP(J)	MAIN325
C		MAIN326
	PHI(I)=1.0/RPHI(I)	MAIN327
	DPHI(I)=ABS(PHI(I-1)-PHI(I))	MAIN328
	IF(RI(I).LE.0.0) GO TO 10	MAIN328A
	SP=SIN(PHI(I))	MAIN329
	TP=TAN(PHI(I))	MAIN330
	SAP=SIN(ALRAD+PHI(I))	MAIN331
C	CALCULATE TAN(BETA-ALPHA)	MAIN332
C		MAIN333
	TALPHAB=(D7*(SAP-SA)*TP+R(I)*DPHI(I)*CA)/(DZ*(SAP-SA)-R(I)*DPHI(I)	MAIN334
	1 *CA*TP)	MAIN335
C		MAIN336
C	SOLVE FOR THE ANGLE BETA IN RADIAN.	MAIN337
C		MAIN338
	BETA(I)=ATAN(TALPHAB)+ALRAD	MAIN339
	IF (I.EQ.1) BETA(I)=ALRAD+PHI(I)	MAIN340
C		MAIN341
C	SOLVE FOR THE POINT OF COLLAPSE ON THE AXIS, STAGNATION POINT +SP.	MAIN342
	C(I)=Z(I)*SP/(CA*(SAP-SA))	MAIN343
C	SOLVE FOR COLLAPSE VELOCITY	MAIN344
	V(I)=2.0*DA*SIN(PHI(I)*.5)	MAIN345
	GAMA(I)=.5*PI-(ALRAD+.5*PHI(I))	MAIN346
	VJ(I)=V(I)*COS(ALRAD+.5*PHI(I)-.5*BETA(I))/SIN(.5*BETA(I))	MAIN347
	VN(I)=V(I)*SIN(ALRAD+.5*PHI(I)-.5*BETA(I))/COS(.5*BETA(I))	MAIN348
	DML(I)=PI*EPS*R(I)*DZ/CA	MAIN349
	DML(I)=DML(I)*RHOJ	MAIN349A
	DMASS=DML(I)	MAIN349B
	DMN(I)=.5*(1.+COS(BETA(I)))	MAIN350
	DMJ(I)=.5*(1.-COS(BETA(I)))	MAIN351
	AT=2.*(ALRAD+.5*PHI(I)-.5*BETA(I))	MAIN352
	DEJ(I)=.5*(1.+COS(AT))	MAIN353
	DEN(I)=.5*(1.-COS(AT))	MAIN354
C		MAIN355
C	COMPUTE THE RELATIVE VELOCITY OF THE JET AND SLUG.	MAIN356
C	THIS IS THE FLOW VELOCITY. (VF).	MAIN357
	RV(I)=.5*(VJ(I)-VN(I))	MAIN358
	IF(RV(I).GT.RVMAX) RVMAX=RV(I)	MAIN358A
C		MAIN359
	TML=TML+DML(I)	MAIN360
	TMJ=TMJ+DMJ(I)*DML(I)	MAIN361
C	IF THE FLOW VELOCITY IS GREATER THAN 1.23*CO, REMOVE MASS FROM JET	MAIN361A
	IF(RV(I)/CO.GT.1.23) TMJ=0.	MAIN361B
C	ALSO REMOVE KINETIC ENERGY FOR JET.	MAIN361C
	IF(RV(I)/CO.GT.1.23) TEJ=0.	MAIN361D
	TMN=TMN+DMN(I)*DML(I)	MAIN362
	DMJK(I)=DMJ(I)*DML(I)	MAIN363
	DEL(I)=.5*DML(I)*V(I)**2	MAIN364
	TKE=TKE+DEL(I)	MAIN365
	TEJ=TEJ+DEJ(I)*DEL(I)	MAIN366
	UJK(I)=VJ(I)	MAIN367
	IF (VJ(I).LE.VJMAX) GO TO 9	MAIN368
	VJMAX=VJ(I)	MAIN369

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DMJTE=DMJTE+DMJK(I)
SUMVJ=SUMVJ+VJ(I)*DMJK(I)
VJTE=SUMVJ/DMJTE
MM=J
9 CONTINUE
10 CONTINUE
END OF TIME AND SPACE INCREMENT ALONG THE HEIGHT OF THE LINER.

CALL SUBROUTINE TO COMPRESS THE JET AND CALCULATE NEW JET VELOCITY
CALL VELGR (N,UJK,DMJK,IS)
TEJCUT=0.
DO 12 I=IS,N
  IF(UJK(I).GE..5) TMCUT1=TMCUT1+DMJK(I)
  IF(UJK(I).GE..4) TMCUT2=TMCUT2+DMJK(I)
  IF(UJK(I).GE..3) TMCUT3=TMCUT3+DMJK(I)
  IF(UJK(I).GE..2) TMCUT4=TMCUT4+DMJK(I)
  IF(UJK(I).GE..1) TMCUT5=TMCUT5+DMJK(I)
  IF(UJK(I).GE..25) TMJCUT=TMJCUT+DMJK(I)
  IF(UJK(I).GE..25) TEJCUT=TEJCUT+DEJ(I)*DEL(I)
  IF(UJK(I).GE..5) TECUT1=TECUT1+DEJ(I)*DEL(I)
  IF(UJK(I).GE..4) TECUT2=TECUT2+DEJ(I)*DEL(I)
  IF(UJK(I).GE..3) TECUT3=TECUT3+DEJ(I)*DEL(I)
  IF(UJK(I).GE..2) TECUT4=TECUT4+DEJ(I)*DEL(I)
  IF(UJK(I).GE..1) TECUT5=TECUT5+DEJ(I)*DEL(I)
  IF(UJK(I).LE.UJKMAX) GO TO 11
  IF(RV(I)/CO.GT.1.23) TEJCUT=0.
  IF(RV(I)/CO.GT.1.23) TMJCUT=0.
  IF(RV(I)/CO.GT.1.23) TECUT1=0.
  IF(RV(I)/CO.GT.1.23) TECUT2=0.
  IF(RV(I)/CO.GT.1.23) TECUT3=0.
  IF(RV(I)/CO.GT.1.23) TECUT4=0.
  IF(RV(I)/CO.GT.1.23) TECUT5=0.
  IF(RV(I)/CO.GT.1.23) TMCUT5=0.
  IF(RV(I)/CO.GT.1.23) TMCUT4=0.
  IF(RV(I)/CO.GT.1.23) TMCUT3=0.
  IF(RV(I)/CO.GT.1.23) TMCUT2=0.
  IF(RV(I)/CO.GT.1.23) TMCUT1=0.
  TMASS=TMASS+DMJK(I)
  UJKMAX=UJK(I)
  MN=I
11 CONTINUE
  IF(RV(I)/CO.GT.1.23) TMASS=DMJK(I+1)
  IF(RV(I)/CO.GT.1.23) UJKMAX=UJK(I+1)
  IF(RV(I)/CO.GT.1.23) MN=I+1
12 CONTINUE
  N1=1
  N2=N
  IF(N.GE.57) N2=57
13 CONTINUE
  WRITE(6,23)
  WRITE(6,24)
  DO 14 I=N1,N2
    PHI(I)=PHI(I)*DEGTOR
    BETA(I)=BETA(I)*DEGTOR
    GAMA(I)=GAMA(I)*DEGTOR

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MAIN370
MAIN371
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MAIN381
MAIN382
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MAIN383A
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MAIN383C
MAIN383D
MAIN383E
MAIN383F
MAIN384
MAIN384A
MAIN384B
MAIN384C
MAIN384D
MAIN384E
MAIN385
MAIN385A
MAIN385B
MAIN385C
MAIN385D
MAIN385E
MAIN385F
MAIN385G
MAIN385H
MAIN385I
MAIN385J
MAIN385K
MAIN385L
MAIN386
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MAIN389A
MAIN389B
MAIN389C
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MAIN397
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MAIN399
MAIN400

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DPHI(I)=DPHI(I)*DEGTOR	MAIN401
Z(I)=Z(I)/H	MAIN402
WRITE (6,23) I,Z(I),TZ(I),EPSI(I),E(I),PHI(I),BETA(I),DPHI(I),	MAIN403
1 RPHI(I),V(I),W(I),YAU(I),C(I)	MAIN404
BETA(I)=BETA(I)/DEGTOR	MAIN405
14 CONTINUE	MAIN406
IF (N.LE.N2) GO TO 15	MAIN407
N1=N2+1	MAIN408
N2=N	MAIN409
GO TO 13	MAIN410
15 CONTINUE	MAIN411
TMJ=0.0	MAIN411A
N1=1	MAIN412
N2=N	MAIN413
IF (N.GE.57) N2=57	MAIN414
16 CONTINUE	MAIN415
WRITE (6,23)	MAIN416
WRITE (6,25)	MAIN417
DO 17 I=N1,N2	MAIN418
TMJ=TMJ+DMJK(I)	MAIN418A
IF (RV(I)/CO.GT.1.23) TMJ=0.	MAIN418B
WRITE (6,23) I,Z(I),VJ(I),VN(I),RV(I),DMJ(I),DMN(I),DEL(I),DML(I),	MAIN419
1DMJK(I),TMJ	MAIN420
C Z(I)=Z(I)*H	MAIN421
17 CONTINUE	MAIN421A
IF (N.LE.N2) GO TO 18	MAIN422
N1=N2+1	MAIN423
N2=N	MAIN424
GO TO 16	MAIN425
18 CONTINUE	MAIN426
IF (IERROR.EQ.1) GO TO 20	MAIN427
IF (RHOC.EQ.0.0) GO TO 20	MAIN427A
C CALL PENETRATION SUBROUTINE CALCULATES JET RADIUS,DEPTH OF	MAIN428
C PENETRATION, HOLE RADIUS,,AND ALSO PENETRATION STANDOFF CURVES.	MAIN429
C CALL PENTRAT (N,Z,R,TZ,DMJ,BETA,TAU,C,UJK,RV,VJO,DMASS,MN,HEAD,PSO	MAIN430
1,PPT,NCD,EPSI,DAI)	MAIN431
C	MAIN432
C	MAIN433
C	MAIN434
C	MAIN435
DO 19 I=1,N	MAIN436
C CONVERT VELOCITIES TO (MM/MICROSEC) FOR PLOTTING AND REPORTS.	MAIN437
C ALSO CONVERT DISTANCES TO MM.	MAIN438
C	MAIN439
BETA(I)=BETA(I)*DEGTOR	MAIN440
VJ(I)=VJ(I)*10.	MAIN441
V(I)=V(I)*10.	MAIN442
C(I)=C(I)*10.	MAIN443
UJK(I)=10.*UJK(I)	MAIN444
Z(I)=Z(I)/H	MAIN444A
19 CONTINUE	MAIN445
IF (NPLT.EQ.0) GO TO 20	MAIN446
CALL PLOTS (N,UJK,DPOINT,RDPT,HEAD,PSO,PPT,NCD,NPLT,IS)	MAIN447
20 CONTINUE	MAIN448
UJKMAX=UJKMAX*10.	MAIN449
C PRINT SUMMARY OF RESULTS	MAIN450
C	MAIN451

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WRITE (6,39) HEAD
WRITE (6,39) TML,TMJ,TMN,TKE,TEJ,TEJCUT,TMJCUT
AMMACH=RVMAX/CO
WRITE (6,41) TECUT1,TMCUT1,TECUT2,TMCUT2,TECUT3,TMCUT3
1TECUT4,TMCUT4,TECUT5,TMCUT5,AMMACH
WRITE (6,35) MN,2(MN),UJCMAX,TMASS
NPOS=0
IF (IJOHN.GT.1) GO TO 1
GO TO 2
9H765 WRITE (6,40)
WRITE (6,23) I,TW,RP(J),RPHI(I),CT,BETA(1)
IEROR=1
STOP

C
C
C
FORMATS

23 FORMAT (14,2X,1P12E10,3)
24 FORMAT (2H11,5X,2H Z,8X,2HTZ,8X,3HEPS,7X,2H E,6X,3HPI,7X,4HBETA,6X,2H1N467
1X,4HDPHI,6X,4HDPHI,6X,4H V,6X,4H R,6X,4H TAU,6X,4H C/) MAIN468
25 FORMAT (2H11,6X,1HZ,8X,3H VJ,8X,2HVN,7X,3H RV,7X,3HDMJ,7X,3HDMN,7X,MAIN469
1,3HDEL,7X,3HDL,7X,4HDMJK,9X,3HTMJ/) MAIN470
26 FORMAT (8E10,4) MAIN471
27 FORMAT(1H1/' INITIAL POSITION'/) MAIN472
28 FORMAT( ///50X,'INPUT PARAMETERS'///5X,' ALPHA = ',F6.3,7X,'K COMAIN473
INSTANT = ',F6.3,7X,'THICKNESS OF LINER = ',F6.3,7X,' DENSITY = ',F6.3,7X,' FACMAIN474
2F6.3//5X,A10,5X,' THICKNESS = ',F6.3,7X,' DENSITY = ',F6.3,5X,' FACMAIN475
3TOP = ',F6.4//5X,' LINER HEIGHT = ',F10.4,7X,' LINER RADIUS = ',F10.4,MAIN476
4,///5X,' INCREMENT OF Z (DZ) = ',F10.4,7X,' INCREMENT OF TIME (DT) MAIN477
5 = ',F10.4,7X,' NUMBER OF ELEMENTS (N) = ',14,///50X,' DETONATI(N'///) MAIN478
29 FORMAT(5X,'EXPLOSIVE = ',5X,A10,3X,' DETONATION VELOCITY(D) = ', MAIN479
1F10.5,5X,'EXPLOSIVE DENSITY = ',F6.3//5X,' POINT OF DETONATION FROM MAIN480
2CONE APEX(DPOINT) = ',F10.4,7X,' DET. RADIUS (RDPT) = ',F10.4//50X,MAIN481
3'TARGET INPUTS'///5X,' TARGET DENSITY = ',F7.3,7X,' STAND-OFF DISTANCEMAIN482
4E = ',F7.3,7X,' CK CONSTANT(BHN) = ',F10.6//50X,' MISC. INPUTS'///5X, MAIN483
5A10,' PLOTS') MAIN484
30 FORMAT(//////////,30X,' SUMMARY OF RESULTS'///////// LINER MASS = 'MAIN485
1F10.4,' GM 1,5X,' JET MASS = ',F10.4,' GM 1,5X,' SLUG MASS = ',MAIN486
2F10.4,' GM 1/// TOTAL KINETIC ENERGY = ',F10.4,' ERGS*1.0E12'///MAIN487
3 TOTAL JET KINETIC ENERGY = ',F10.4,' ERGS*1.0E12'/// TOTAL JET K1MAIN488
4NETIC ENERGY ABOVE JET VELOCITY .25 CM/MSEC = ',F10.4,' ERGS*1.0E1MAIN489
52'/// TOTAL JET MASS ABOVE JET VELOCITY .25 CM/MSEC = ',F10.4,' GMMAIN490
6 '///) MAIN491
31 FORMAT (///7(3H 1,6X,2H Z,7X)) MAIN492
32 FORMAT (///7(3H 1,6X,2H R,7X)) MAIN493
33 FORMAT (///7(3H 1,6X,2HPI,7X)) MAIN494
34 FORMAT (7(14,4X,F7.4,3X)) MAIN495
35 FORMAT(1 JET TIP AT 1 = '1,5,' RELATIVE POSITION Z/H = ',F10.8//MAIN496
1' MEASURED JET TIP VELOCITY = ',F10.8,' MM/MICROSEC'/// MEASURED MAIN497
2ASS OF JET TIP = ',F10.4,' GM 1/1H)) MAIN498
36 FORMAT (7E10,3,215) MAIN499
37 FORMAT (11,8A10) MAIN500
38 FORMAT (11,30X,8A10) MAIN501
39 FORMAT (1H1,30X,8A10) MAIN502
40 FORMAT (1ERROR) MAIN503
41 FORMAT(// KINETIC ENERGY ABOVE .5 CM/MSEC = ',F10.4,
1' AND JET MASS = ',F10.4,
2' KINETIC ENERGY ABOVE .4 CM/MSEC = ',F10.4,
3' AND JET MASS = ',F10.4,
4' KINETIC ENERGY ABOVE .3 CM/MSEC = ',F10.4,
5' AND JET MASS = ',F10.4,
6' KINETIC ENERGY ABOVE .2 CM/MSEC = ',F10.4,
7' AND JET MASS = ',F10.4,
8' KINETIC ENERGY ABOVE .1 CM/MSEC = ',F10.4,
9' AND JET MASS = ',F10.4//
A' MAXIMUM RELATIVE MACH NUMBER = ',F10.4//)
END

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	SUBROUTINE PENTRAT (N,Z,P,TZ,DMJ,BETA,TAU,C,VJ,RV,VJO,DMASS,MN,	* 511* 1
	1 HEAD,PSO,PPT,NCD,EPSI,DAI)	PENTT 2
C		PENTT 3
C	COMPUTE PENETRATION,RADIUS OF JET,RADIUS OF THE HOLE.	PENTT 4
C		PENTT 5
	DIMENSION HEAD(8)	PENTT 6
	DIMENSION EPSI(100), DAI(100)	PENTT 7
	DIMENSION PPT(50), PSO(50), PVJT(50)	PENTT 8
	DIMENSION Z(100), TZ(100), DMJ(100), BETA(100), TAU(100), C(100),	PENTT 9
	1 VJ(100), R(100), RV(100)	PENTT10
C		PENTT11
C	COMMON	PENTT12
C		PENTT13
	COMMON ALHAD, EPS, RHOJ, RHOC, FF, RPHIO, DTZ, SO, CK, DA, H, U,	PENTT14
	1 PT, DERTOP	PENTT15
	COMMON AMU(100), THETA(100), F(100), DF(100), T(100), DT(100),	PENTT16
	1 G(100), P(100), A(100), DELA(100), DVJ(100), ZUDT(100),	PENTT17
	2 RSG(100), RC(100), T1, UMIN	PENTT18
C	PENETRATION PHASE	PENTT19
	VJO=0.0	PENTT19A
	DO 11 I=MN,N	PENTT20
	EPS=EPSI(I)	PENTT21
	DA=DAI(I)	PENTT22
	IF (1.LE.MN) VJ(I)=VJ(MN)	PENTT23
	IF (1.EQ.1) RETURN	PENTT24
	IF (VJO.GT.0.0) GO TO 4	PENTT25
	VJO=VJ(MN)	PENTT26
	JC1=1	PENTT27
	AKAY=SQRT(RHOJ/RHOC)	PENTT28
	AK1=AKAY+1.	PENTT29
	AK3=(3.+AKAY)/(2.*AKAY)	PENTT30
	AKOK1=AKAY/AK1	PENTT31
	A1OK=1./AK1	PENTT32
	BHN=320.	PENTT32A
	IF (UMIN.GT.1.0) BHN=UMIN	PENTT32B
	IF (CK.EQ.0.) CK=(2250.+4.20*BHN)*1.0E-5	PENTT33
	IF (UMIN.GT.1.0) UMIN=(CK*1.0E5-1350.)/20400.	PENTT34
C	T1=T1*(.7511-.28)/(VJO-.28)	PENTT35
	M=MN	PENTT36
	A(I)=0.	PENTT37
	DELA(I)=0.	PENTT38
	DF(I)=0.	PENTT39
	F(I)=0.	PENTT40
	DVJ(I)=0.	PENTT41
	SUMMU=0.	PENTT42
	K=1-1	PENTT42A
	T1(K)=TZ(K)-TZ(1)	PENTT42B
	DMASS=DMJ(K)	PENTT42C
	IF (VJ(K).LE.0.0) VJ(K)=1.0E-02	PENTT42D
	AMU(K)=V. J(K)	PENTT42E
	THETA(K)=Y. TAU(K)	PENTT42F
	F(K)=THETA(K)/AMU(K)-C(K)/VJO	PENTT42G
	DF(K)=F(K)	PENTT42H
	SUMMU=AMU(K)*(-AKAY)*DF(K)+SUMMU	PENTT42I
	T(K)=AMU(K)**(1.+AKAY)*(TO-(1.+AKAY)*SUMMU)	PENTT42J
	DT(K)=T(K)-TO	PENTT42K
	G(K)=VJO*(T(K)/AMU(K)-F(K))	PENTT42L

	A(K)=Z(K)-R(K)*COT(ETA(K)/2.)	PENTT42M
	DELA(K)=A(K)	PENTT42N
	VJ(K)=VJ(2)	PENTT42O
	NCD=1	PENTT43
	IF (SO.EQ.0.) NCD=25	PENTT44
	DO 3 J=1,NCD	PENTT45
	IF (NCD.GT.1) SO=FLOAT(J)*RF*2.	PENTT46
	ZO=H+SO	PENTT47
	TO=ZO/VJO	PENTT48
	IF (ZO.GE.VJO*T1) GO TO 1	PENTT49
	PT=AK1*(VJO*T1)**AKOK1*ZO**A1OK-SQRT(AKAY*AK1*UMIN*T1*(VJO*T1)*	PENTT50
	1 *AKOK1*ZO**A1OK)-ZO	PENTT51
	VJT=(AK1*(VJO*T1)**AKOK1*ZO**A1OK-(ZO+PT))/(AKAY*T1)	PENTT52
	GO TO 2	PENTT53
1	CONTINUE	PENTT54
	PT=AKAY*(VJO*T1-SQRT(UMIN*T1*(VJO*T1+ZO/AKAY)))	PENTT55
	VJT=VJO-PT/(AKAY*T1)	PENTT56
2	CONTINUE	PENTT57
	PSO(J)=FLOAT(J)	PENTT58
	POT(J)=PT	PENTT59
	PPT(J)=PPT(J)*10.	PENTT60
	PVJT(J)=VJT	PENTT61
3	CONTINUE	PENTT63
C	IF PLOTTING PENETRATION- STANDOFF CURVE USE 2. CD TO PLOT HOLE PRO	PENTT64
	IF (NCD.GT.1) NCDI=3	PENTT64A
	IF (NCD.GT.1) SO=FLOAT(NCDI)*RF*2.	PENTT65
	IF (NCD.GT.1) PT=PPT(NCDI)*.1	PENTT66
	IF (NCD.GT.1) VJT=PVJT(NCDI)*.1	PENTT67
	ZO=SO+H	PENTT68
	TO=ZO/VJO	PENTT69
C	RETURN	PENTT70
4	CONTINUE	PENTT71
	TZ(I)=TZ(I)-TZ(1)	PENTT72
	DMASS=DMJ(I)	PENTT75
	IF (VJ(I).LE.0.0) VJ(I)=1.0E-02	PENTT76
	AMU(I)=VJO/VJ(I)	PENTT77
	IF (AMU(I).LT.0.) RETURN	PENTT78
	THETA(I)=TZ(I)+TAU(I)	PENTT79
	F(I)=THETA(I)/AMU(I)-C(I)/VJO	PENTT79A
	DF(I)=F(I)-F(I-1)	PENTT80
	SUMMU=AMU(I)**(-AKAY)*DF(I)+SUMMU	PENTT81
	T(I)=AMU(I)**(1.+AKAY)*(TO+(1.+AKAY)*SUMMU)	PENTT82
	DT(I)=T(I)-TO	PENTT83
	G(I)=VJO*(T(I)/AMU(I)-F(I))	PENTT84
	A(I)=Z(I)-R(I)*COT(BETA(I)/2.)	PENTT85
	DELA(I)=A(I)-A(I-1)	PENTT86
	DVJ(I)=VJ(I)-VJ(I-1)	PENTT87
	DZODT(I)=(T(I)-TZ(1))*DVJ(I)/DTZ-VJ(I)+DELA(I)/DTZ	PENTT88
C	DZODT(I)=VJ(I)-((T(I)-TZ(1))*DVJ(I)/DTZ+DELA(I)/DTZ)	PENTT89
	IF (T(I).GT.T1) GO TO 6	PENTT90
	IF (ZO.GE.VJO*T1) GO TO 5	PENTT91
C		PENTT92
C	CALCULATE DEPTH OF PENETRATION	PENTT93
C		PENTT94
	P(I)=ZO*((T(I)/TO)**AKOK1-1.0)	PENTT95
	RSQ(I)=2.*R(I)*EPS*DA*DMASS/DZODT(I)	PENTT96
	RSQ(I)=SQRT(ABS(P(I)))	PENTT97

C		PENT198
C	CALCULATE RADIUS OF THE JET	PENT199
C		PENT100
	RC(I)=SQRT(RHOJ/(2.*AKAY*CK))*SQRT(T1/TO)*VJO*(ZO/(ZO+P(I)))*AK3	PENT101
	1 *RSQ(I)	PENT102
	GO TO 9	PENT103
5	CONTINUE	PENT104
	P(I)=VJO*(T(I)-TO)*T1/(T1+T(I)/AKAY)	PENT105
	VJP=VJO*(T1+TO/AKAY)/(T1+T(I)/AKAY)	PENT106
	DVJP=VJP-VJPO	PENT107
	DZODT(I)=(T1-T(I))*DVJP-VJP*DELA(I)	PENT108
	DZODT(I)=-DZODT(I)	PENT109
	RSQ(I)=2.*R(I)*EPS*DA*DMASS/DZODT(I)	PENT110
	RSQ(I)=SQRT(RSQ(I))	PENT111
	RC(I)=SQRT(RHOJ/(2.*AKAY*CK))*RSQ(I)*VJO*(1.-P(I)/(AKAY*VJO*T1))	PENT112
	VJPO=VJP	PENT113
	GO TO 10	PENT114
6	CONTINUE	PENT115
	IF (P(I).LT.PT) GO TO 8	PENT116
	IF (JC1.NE.1) GO TO 7	PENT117
	JC1=2	PENT118
	TP=T(I)	PENT119
7	CONTINUE	PENT120
	IF (T(I).GT.TP) RETURN	PENT121
8	CONTINUE	PENT122
	DZODT(I)=(T1-TZ(I))*DVJ(I)/DTZ-VJ(I)*DELA(I)/DTZ	PENT123
	P(I)=(AK1*(T1/TO)*AKOK)*T(I)/(T(I)+AKAY*T1)-1.0)*ZO	PENT124
	RSQ(I)=2.*R(I)*EPS*DA*DMASS/DZODT(I)	PENT125
	RSQ(I)=SQRT(RSQ(I))	PENT126
	RC(I)=SQRT(RHOJ/(2.*AKAY*CK))*(VJO/AKAY)*(AK1*(ZO/(VJO*T1))*A1OK	PENT127
	1 -(ZO+P(I))/(VJO*T1))*RSQ(I)	PENT128
9	CONTINUE	PENT129
	VJPO=VJ(I-1)	PENT130
10	CONTINUE	PENT131
11	CONTINUE	PENT132
	WRITE (6,31) HEAD	PENT133W
	WRITE (6,22)	PENT134W
	WRITE (6,29) T1,TO,ZO,PT,VJT,VJO,AKAY,UMIN,CK	PENT135W
	N1=1	PENT136
	N2=N	PENT137
	IF (N.GE.57) N2=55	PENT138
	WRITE (6,30)	PENT139W
12	CONTINUE	PENT140
	WRITE (6,23)	PENT141W
	WRITE (6,27)	PENT142W
	II=0	PENT143
	DO 13 I=N1,N2	PENT144
	WRITE (6,23) I,AMU(I),THETA(I),F(I),DF(I),T(I),DT(I),G(I),RSQ(I),	PENT145W
	1 A(I),DELA(I),DVJ(I),DZODT(I)	PENT146W
13	CONTINUE	PENT147
	IF (N.LE.N2) GO TO 14	PENT148
	N1=N2+1	PENT149
	N2=N	PENT150
	GO TO 12	PENT151
14	CONTINUE	PENT152
	DO 15 I=MN,N	PENT153
15	CONTINUE	PENT154

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N1=1
N2=N
IF (N.GE.50) N2=50
WRITE (6,31) HEAD
SO=SO/(2.*RF)
WRITE (6,28) SO
WRITE (6,23)
WRITE (6,25)
16 CONTINUE
IF (N1.GT.25) WRITE (6,26)
DO 19 I=N1,N2
IF (P(I).GT.PT) GO TO 20
IF (P(I).EQ.0.0) II=II+1
IF (I.GT.25) GO TO 17
WRITE (6,24) I,P(I),RC(I),PPT(I),PSO(I)
GO TO 18
17 CONTINUE
WRITE (6,24) I,P(I),RC(I)
18 CONTINUE
19 CONTINUE
IF (N.LE.N2) GO TO 20
N1=N2+1
N2=N
GO TO 16
20 CONTINUE
II=II+1
DO 21 I=1,II
RC(I)=RC(II)
21 CONTINUE
RETURN
C
C
C
PENETRATION FORMATS
22 FORMAT(40X,'INITIAL CONDITIONS FOR PENETRATION'//)
23 FORMAT (14,2X,1P12E10.3)
24 FORMAT (13,7X,2(F10.5,10X,F10.5,20X))
25 FORMAT (3H 1,12X,'P(CM)',15X,'PC(CM)',23X,'PT(MM)',15X,'SO(CD)')
26 FORMAT (1H1,3H 1,13X,'P',18X,'RC')
27 FORMAT (2H 1,6X,2PMU,8X,5MTHETA,5X,2H F,7X,3H DF,7X,2H T,8X,2HDT,8P
1X,2H G,8X,2HRSQ,8X,1HA,8X,4HDELA,7X,2HGV,6X,4HVDZ1/)
28 FORMAT(10X,'HOLE PROFILE SO = ',F6.2,3X CD,25X,'PENTRATION STAN
1DOFF'//)
29 FORMAT(5X,'T1 = ',F10.4,5X,'T0 = ',F10.4,5X,' Z0 = ',F10.4,4X,
1' PT = ',F10.4,5X,'VJT = ',F10.4,5X,' VJO = ',F10.4,4X,' AKAY = '
2,F10.6,5X,' UMIN = ',F10.5,5X,' CK = ',F10.5//)
30 FORMAT(1H1,' PENETRATION PHASE'//)
31 FORMAT(1H1,30X,RA10//)
END

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	SUBROUTINE CALAI (I,Z,R,DPOINT,RDPT,JOHNI,AI)	* 701* 2
C		CALAI 2
C	SUBROUTINE WILL CALCULATE INCLINATION ANGLE (AI), DETONATION	CALAI 3
C	COMPONENT OF THE COLLAPSE VELOCITY (DA), AND WILL CALL CALPHI	CALAI 4
C	SUBROUTINE WHICH WILL CALCULATE 1./PHI0 (RPHI0) GIVEN (AI).	CALAI 5
C		CALAI 6
	DIMENSION Z(100), R(100)	CALAI 7
	COMMON ALRAD, EPS, RHOJ, RHOC, RF, RPHI0, DTZ, SO, CK, DA, H, D.	CALAI 8
	1 PT, DEGTOR	CALAI 9
	ALPHA=ALRAD*DEGTOR	CALAI 10
	THETA=ATAN((R(I)-RDPT)/(DPOINT+Z(I)))	CALAI 11
	THETA=THETA*DEGTOR	CALAI 12
	AI=90.-(ALPHA-THETA)	CALAI 13
	IF (AI.LT.0.) AI=180.-(ALPHA-THETA)	CALAI 14
	CALL CALPHI (AI)	CALAI 15
	AI=AI/DEGTOR	CALAI 16
	DA=D/SIN(AI)	CALAI 17
	RETURN	CALAI 18
	END	CALAI 19-

	SUBROUTINE CALPHI (AI)	* 720* 3
C		CALPI 2
C	GIVEN AN INCLINATION ANGLE (AI) INTERPOLATE INTO TABLES AND FIND	CALPI 3
C	AN ANGLE PHIO. THEN CALCULATE 1./PHIO (RPHIO) AND RETURN.	CALPI 4
C		CALPI 5
	DIMENSION TPHI(37), TAI(37)	CALPI 6
	COMMON ALRAD, EPS, RHOJ, RMOC, RF, RPHIO, DTZ, SO, CK, DA, H, D,	CALPI 7
	1 PT, DEGTOR	CALPI 8
	DATA TPHI /0.0, .08, .25, .59, 1.02, 1.43, 1.76, 2.03, 2.47, 2.86,	CALPI 9D
	1 3.30, 3.70, 4.14, 4.51, 5.08, 5.65, 5.98, 6.37, 6.70, 7.01, 7.22,	CALPI10U
	2 7.42, 7.58, 7.73, 7.90, 8.04, 8.17, 8.23, 8.26, 8.30, 8.26, 8.16,	CALPI11D
	3 8.06, 7.99, 7.85, 7.75, 7.56/	CALPI12D
	DATA TAI /0.0, .02, .05, .15, .23, .32, .42, .48, .60, .72, .84, .	CALPI13D
	194, 1.06, 1.14, 1.29, 1.44, 1.54, 1.67, 1.79, 1.92, 2.03, 2.16, 2.	CALPI14D
	228, 2.41, 2.61, 2.79, 3.01, 3.14, 3.25, 3.38, 3.57, 3.72, 3.88, 4.	CALPI15D
	300, 4.16, 4.29, 4.47/	CALPI16D
C	TAI SCALE FACTOR 100 DEG = 5.01 COUNTS (SFI=5.01/100.)	CALPI17
C	TPHI SCALE FACTOR 40 DEG = 8.08 COUNTS (SFPHI=8.08/40.)	CALPI18
	SFI=.0501	CALPI19
	SFPHI=.207	CALPI20
	X=SFI*AI	CALPI21
	DO 1 J=1,37	CALPI22
	IF (TAI(J).GT.X) GO TO 2	CALPI23
1	CONTINUE	CALPI24
2	CONTINUE	CALPI25
	DXI=X-TAI(J-1)	CALPI26
	RDX=DXI/(TAI(J)-TAI(J-1))	CALPI27
	Y=RD*(TPHI(J)-TPHI(J-1))+TPHI(J-1)	CALPI28
	IF (PHIO.LE..0001) PHIO=0.	CALPI29
	PHIO=Y/SFPHI	CALPI30
	RPHIO=1.0/PHIO	CALPI31
	RPHIO=RPHIO*DEGTOR	CALPI32
	RETURN	CALPI33
	END	CALPI34-

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      SUBROUTINE VELGW (IK,UJK,DMJK,IS)
      DIMENSION UJK(200), DMJK(200)
      E=0.
      IJK=IK-1
      1 KP=1
      C      COMPRESSION OF THE JET
      DO 2 M=IS,IJK
      IF (UJK(M).GE.UJK(M+1)) GO TO 2
      IF (ABS(UJK(M)-UJK(M+1)).LE.1.E-5) GO TO 2
      T1=DMJK(M)*UJK(M)+DMJK(M+1)*UJK(M+1)
      T2=2*DMJK(M+1)*(UJK(M)-UJK(M+1))
      DEM=1.0/(DMJK(M)+DMJK(M+1))
      UJK(M)=(T1-T2)*DEM
      UJK(M+1)=(T1+T2)*DEM
      KP=0
      2 CONTINUE
      IF (KP.EQ.0) GO TO 1
      RETURN
      END

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      * 754 * 4
      VELGW 2
      VELGW 3
      VELGW 4
      VELGW 5
      VELGW 6
      VELGW 7
      VELGW 8
      VELGW 9
      VELGW 10
      VELGW 11
      VELGW 12
      VELGW 13
      VELGW 14
      VELGW 15
      VELGW 16
      VELGW 17
      VELGW 18
      VELGW 19-

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SURROUTINE PLOTS (N,UJK,DPOINT,RDPT,HEAD,PSO,PPT,NCD,NPLT,ISTART) * 773* 5
DIMENSION PSO(50), PPT(50), S(50) PLOTS 2
DIMENSION R(5000), HEAD(R) PLOTS 3
DIMENSION UJK(200), DMJK(200) PLOTS 4
DIMENSION RD(4), CV(4), AB(4), TC(4), PC(4), DMNJ(4), VVC(4) PLOTS 5
DIMENSION Z(100), TZ(100), TAU(100), E(100), PHI(100), DPHI(100), PLOTS 6
1 BETA(100), GAMA(100), R(100), C(100), VJ(100), VN(100), DMJ(100) PLOTS 7
2 , DMN(100), DEJ(100), DEN(100), RPHI(100), V(100) PLOTS 8
DIMENSION F(4), G(4) PLOTS 9
COMMON ALRAD, EPS, RHOJ, RHOC, RF, RPHIC, DTZ, SO, CK, DA, H, D, PLOTS 10
1 PT, DEGTOR PLOTS 11
COMMON /LPLOT/ Z, TZ, TAU, E, PHI, DPHI, BETA, GAMA, R, C, VJ, VN, PLOTS 12
1 DMJ, DMN, DEJ, DEN, RPHI, V PLOTS 13
DATA DMNJ(1), DMNJ(2) /10HRELATIVE M, 4HASS>/ PLOTS 14D
DATA RD(1), RD(2), RD(3) /10HRELATIVE D, 10HISTANCE FR, 10HOM CONE PLOTS 15D
1 AP/ PLOTS 16D
DATA RD(4) /3HEX>/ PLOTS 17D
DATA CV(1), CV(2), CV(3) /10HCOLLAPSE V, 10HELOCITY IN, 10H CM/MIC PLOTS 18D
1KOS/ PLOTS 19D
DATA CV(4) /3HEC>/ PLOTS 20D
DATA AB(1), AB(2), AB(3) /10HANGLE BETA, 10H (DEGREE), 1H>/ PLOTS 21D
DATA TC(1), TC(2), TC(3) /10HTIME OF CO, 10HLLAPSE (MI, 10HCRG SEC PLOTS 22D
1)>/ PLOTS 23D
DATA PC(1), PC(2), PC(3) /10HPOINT OF C, 10HOLLAPSE (M, 3HM)>/ PLOTS 24D
DATA VVC(1), VVC(2), VVC(3) /10H V, VJ (MM, 10H/MICROSEC), 1H>/ PLOTS 25D
DATA TJET /4HJET>/ PLOTS 26D
DATA TSLUG /5HSLUG>/ PLOTS 27D
XPAGE=14 PLOTS 28
CALL PLTCCR (XPAGE,1,R(1),B(5000)) PLOTS 29
II=N-ISTART PLOTS 30
YMAX=12. PLOTS 31
CALL PLOTS1 (3.0,3.0,1.0,YMAX,.1,.50,RD,VVC,16,10,4) PLOTS 32
CALL PLTCCD (1.0,Z(ISTART),V(ISTART),II) PLOTS 33
CALL PLTCCD (1.0,Z(ISTART),VJ(ISTART),II) PLOTS 34
CALL PLTCCD (4.0,Z(ISTART),UJK(ISTART),II) PLOTS 35
CALL PLCCSP (XS,YS,UFAC) PLOTS 36
XX=.25 PLOTS 37
YY=VJ(35)+.04*.5 PLOTS 38
YY=V(35)+.04*.5 PLOTS 39
PLOTS 40
C DO EXPERIMENTAL DATA POINTS FOR THE 105MM UNCONFINED SHAPED CHARGE. PLOTS 41
C PLOTS 42
C CALL PLTEXP(N) PLOTS 43
C PLOTS 44
DPOINT=DPOINT*10. PLOTS 45
ENCODE (21,2,S(1)) DPOINT PLOTS 46*
RDPT=RDPT*10. PLOTS 47
ENCODE (17,3,S(4)) RDPT PLOTS 48*
ALPHA=ALRAD*DEGTOR PLOTS 49
ENCODE (20,4,S(8)) ALPHA PLOTS 50*
RF=RF*10. PLOTS 51
ENCODE (14,5,S(12)) RF PLOTS 52*
EPS=EPS*10. PLOTS 53
ENCODE (19,6,S(16)) EPS PLOTS 54*
YY=YMAX+.2*YS PLOTS 54A
XMAX=1.0 PLOTS 54B
XX=.5*XMAX-.2*XS*7. PLOTS 54C

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XX=0.0	PLOT554D
CALL PLTCCT (.2,HEAD(1),0.,1.,XX,YY)	PLOT554E
YY=.85*YMAX	PLOT555
XX=.6	PLOT556
CALL PLTCCT (.,5(1),0.,1.,XX,YY)	PLOT557
YY=.80*YMAX	PLOT558
CALL PLTCCT (.,5(4),0.,1.,XX,YY)	PLOT559
XA=.5-.1*XS*10.	PLOT560
YA=-1.1*YS	PLOT561
CALL PLTCCT (.,5(8),0.,1.,XA,YA)	PLOT562
XA=.5-.1*XS*9.	PLOT563
YA=-1.3*YS	PLOT564
CALL PLTCCT (.,5(12),0.,1.,XA,YA)	PLOT565
YA=-1.5*YS	PLOT566
CALL PLTCCT (.,5(16),0.,1.,XA,YA)	PLOT567
C THIS NEXT STATEMENT BYPASSES SOME PLOTTING ROUTINES	PLOT568
C IF (NPLT.GT.0) GO TO 1	PLOT569
CALL PLOTS1 (3.,16.,1.,160.,1.,20.,RD,PC,16,11,4)	PLOT570
CALL PLTCCD (1.0,Z(ISTART),C(ISTART),11)	PLOT571
CALL PLTCCP	PLOT572
CALL PLOTS1 (3.0,3.0,1.0,1.0.,1.,1.,RD,DMNJ,16,10,4)	PLOT573
XX=.35	PLOT574
YY=DMN(35)=.1	PLOT575
CALL PLTCCT (.2,TSLUG,0.,1.,XX,YY)	PLOT576
YY=DMJ(35)=.1	PLOT577
CALL PLTCCT (.2,TJET,0.,1.,XX,YY)	PLOT578
CALL PLTCCD (1.0,Z(ISTART),DMJ(ISTART),11)	PLOT579
CALL PLTCCD (1.0,Z(ISTART),DMN(ISTART),11)	PLOT580
CALL PLOTS1 (3.0,16.,1.0,180.,1.,20.,RD,AB,16,10,4)	PLOT581
CALL PLTCCD (1.0,Z(ISTART),HETA(ISTART),11)	PLOT582
1 CONTINUE	PLOT584
CALL PLTCCP	PLOT585
C THIS NEXT STATEMENT PART OF PLOTTING BYPASS	PLOT586
C	PLOT587
C DO PENETRATION PLOTS( HOLE PROFILES).	PLOT588
C	PLOT589
CALL PLOTPEN (N,S,PSO,PFT,NCD,HEAD)	PLOT590
RETURN	PLOT591
C	PLOT592
2 FORMAT (9H (L-H) = ,F8.4,4H MM>)	PLOT593
3 FORMAT (5H RD = ,F8.4,4H MM>)	PLOT594
4 FORMAT (9H ALPHA = ,F6.2,5H DEG>)	PLOT595
5 FORMAT (6H RF = ,F8.4,4H MM>)	PLOT596
6 FORMAT (7H EPS = ,F8.5,4H MM>)	PLOT597
END	PLOT598

SUBROUTINE PLOTS1 (XBAR,YBAR,XMAX,YMAX,DX,DY,PTX,PTY,NNCX,NNCY,	* 871* 6
1 NCD)	PLOT1 2
DIMENSION PTX(4), PTY(4)	PLOT1 3
CNY=NNCY	PLOT1 4
CNX=NNCX	PLOT1 5
YS=YMAX/5.	PLOT1 6
XS=XMAX/7.	PLOT1 7
CALL PLTCCS (XBAR,YBAR,XMIN,YMIN,XS,YS)	PLOT1 8
HT=.15	PLOT1 9
YY=-.8*YS	PLOT110
XX=.5*XMAX-HT*CNX*XS	PLOT111
CALL PLTCCT (HT,PTX(1),0.,1.,XX,YY)	PLOT112
XX=-.8*XS	PLOT113
YY=.5*YMAX-HT*CNY*YS	PLOT114
CALL PLTCCT (HT,PTY(1),1.,0.,XX,YY)	PLOT115
CALL PLTCCA (DX,DY,XMIN,XMAX,YMIN,YMAX,4)	PLOT116
CALL LABELA (DX,DY,XMIN,XMAX,YMIN,YMAX,1.0,1.0)	PLOT117
RETURN	PLOT118
END	PLOT119-

SUBROUTINE PLOTPEN (N,S,PSU,PPT,NCD,HEAD)	* 840* 7
DIMENSION TSU(4), TPT(4),HEAD(8)	PLOTN 2
DIMENSION S(50), PSO(50), PPT(50)	PLOTN 3
DIMENSION HOL(2), HCN(100), PEN(3), PENT(3), RAD(2)	PLOTN 4
COMMON ALRAD, EPS, RHOJ, RMOC, WF, RPHIO, DTZ, SU, CK, DA, N, D,	PLOTN 5
1 PT, DEGTOR	PLOTN 6
COMMON AMU(100), THETA(100), F(100), DF(100), T(100), DT(100),	PLOTN 7
1 G(100), P(100), A(100), DELA(100), DVJ(100), DZDET(100),	PLOTN 8
2 RSG(100), RC(100), T1, UMIN	PLOTN 9
DATA HOL(1), HOL(2) /10MHOLE PROF1, 3HLE>/	PLOTN10D
DATA PEN(1), PEN(2), PEN(3) /10HRADIUS-VS-, 10MDEPTH OF P, RMEN.(MPL	PLOTN11D
1M)>/	PLOTN12D
DATA PENT(1), PENT(2), PENT(3) /10MTIME-VS-DE, 10MPTH OF PEN, 6M.(	PLOTN13D
1CM)>/	PLOTN14D
DATA RAD(1), RAD(2) /10HRADIUS (MM, 2M)>/	PLOTN15D
DATA TSU(1), TSU(2) /10MSTANDOFF -, 5M(CD)>/	PLOTN16D
DATA TPT(1), TPT(2) /10MPENETRATIO, 7MM (MM)>/	PLOTN17D
RM=0.	PLOTN18
DO 1 I=1,N	PLOTN19
IF (P(I),GE,PT) GO TO 2	PLOTN20
P(I)=-P(I)	PLOTN21
RCN(I)=-RC(I)	PLOTN22
IF (RC(I),GT,RM) RM=RC(I)	PLOTN23
1 CONTINUE	PLOTN24
P(N)=-P(N)	PLOTN25
2 CONTINUE	PLOTN26
N=1	PLOTN27
PM=-P(N)	PLOTN28
XMAX=10.	PLOTN29
XMIN=-XMAX	PLOTN30
YMIN=-PM	PLOTN31
IF (PM,LT,100.) YMIN=-100.	PLOTN32
IF (PM,LT,60.) YMIN=-60.	PLOTN33
YMAX=0.	PLOTN34
YS=(YMAX-YMIN)/R.	PLOTN35
XS=YS	PLOTN36
XBAR=3.0	PLOTN37
YBAR=15.0	PLOTN38
CALL PLTCCS (XBAR,YBAR,XMIN,YMIN,XS,YS)	PLOTN39
XA=-.2*XS*.5*12.	PLOTN40
YA=YMAX+.3*YS	PLOTN41
CALL PLTCCT (.20,HOL(1),0.,1.,XA,YA)	PLOTN42
XA=-.1*XS*5.	PLOTN43
YA=YMIN-.5*YS	PLOTN44
CALL PLTCCT (.1,RAD(1),0.,1.,XA,YA)	PLOTN45
XA=-.1*XS*9.	PLOTN46
YA=YMIN-.75*YS	PLOTN47
CALL PLTCCT (.1,S(1),0.,1.,XA,YA)	PLOTN48
XA=-.1*XS*7.	PLOTN49
YA=YMIN-YS	PLOTN50
CALL PLTCCT (.1,S(4),0.,1.,XA,YA)	PLOTN51
ENCODE (1A,5,T(1) )SO	PLOTN52*
XA=-.1*XS*8.	PLOTN53
YA=YS*1.2	PLOTN54
CALL PLTCCT (.1,T(1),0.,1.,XA,YA)	PLOTN55
XA=XMIN-.8*XS	PLOTN56
YA=.5*YMIN-.1*YS*.5*1R.	PLOTN57

CALL PLTCCT (.1,PEN(2),1.,0.,XA,YA)	PLOTN58
XA=-.2*XS*15.	PLOTN58A
YA=YS*1.4	PLOTN58B
CALL PLTCCT (.2,HEAD(1),0.,1.,XA,YA)	PLOTN58C
DX=1.	PLOTN59
DY=1.	PLOTN60
CALL PLTCCA (DX,DY,XMIN,XMAX,YMIN,YMAX,4)	PLOTN61
DY=2.*DY	PLOTN62
DX=5.0	PLOTN63
XMIN=-XMAX	PLOTN64
XMAX=0.	PLOTN65
CALL LABELA (DX,DY,XVIN,XMAX,YMIN,YMAX,-10.0,-10.0)	PLOTN66
DY=0.	PLOTN67
XMAX=-XVIN	PLOTN68
XVIN=0.	PLOTN69
CALL LABELA (DX,DY,XVIN,XMAX,YMIN,YMAX,10.0,0.0)	PLOTN70
RC(N)=0.	PLOTN71
P(N)=PM	PLOTN72
CALL PLTCCD (1.0,RC(1),P(1),N)	PLOTN73
RCN(N)=0.	PLOTN74
CALL PLTCCD (1.0,RCN(1),P(1),N)	PLOTN75
XA=-.1*XS*10.	PLOTN76
YA=.P*YS	PLOTN77
CALL PLTCCT (.1,S(8),0.,1.,XA,YA)	PLOTN78
XA=-.1*XS*9.	PLOTN79
YA=YS	PLOTN80
CALL PLTCCT (.1,S(12),0.,1.,XA,YA)	PLOTN81
YA=.6*YS	PLOTN82
CALL PLTCCT (.1,S(16),0.,1.,XA,YA)	PLOTN83
IF (NCD.GT.1) GO TO 3	PLOTN84
CALL PLTCCP	PLOTN85
RETURN	PLOTN86
3 CONTINUE	PLOTN87
XMAX=PSO(NCD)	PLOTN88
YMAX=0.	PLOTN89
DO 4 I=1,NCD	PLOTN90
IF (PPT(I).GT.YMAX) YMAX=PPT(I)	PLOTN91
4 CONTINUE	PLOTN92
YMAX=.01*YMAX	PLOTN93
YMAX=AIN(YMAX)+1.0	PLOTN94
YMAX=YMAX*100.	PLOTN95
DY=.2*YMAX	PLOTN96
CALL PLOTS1 (3.0,3.0,XMAX,YMAX,1.,DY,TSO,TPT,7.8,2)	PLOTN97
CALL PLCCSP (XS,YS,UFAC)	PLOTN98
XA=.5*XMAX-.2*XS*7.	PLOTN98A
XA=0.	PLOTN98B
YA=YMAX+.2*YS	PLOTN98C
CALL PLTCCT (.2,HEAD(1),0.,1.,XA,YA)	PLOTN98D
YA=-1.1*YS	PLOTN99
XA=.5*XMAX-.1*XS*10.	PLOT100
CALL PLTCCT (.1,S(8),0.,1.,XA,YA)	PLOT101
XA=.5*XMAX-.1*XS*9.	PLOT102
YA=-1.3*YS	PLOT103
CALL PLTCCT (.1,S(12),0.,1.,XA,YA)	PLOT104
YA=-1.5*YS	PLOT105
CALL PLTCCT (.1,S(16),0.,1.,XA,YA)	PLOT106
CALL PLTCCD (1.0,PSO(1),PPT(1),NCD)	PLOT107
CALL PLTCCP	PLOT108
RETURN	PLOT109
C 5 FORMAT (6F 50 = .F6.2,4F CD>)	PLOT110
END	PLOT111
	PLOT112-



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SURROUTINE PLTEP (N)
DIMENSION PEXX(60), PEXVJ(60), PEXVO(60), PEXV(4), TEXP(3).
1 PEXPX(4)
DIMENSION BASC(3)
COMMON ALHAD, EPS, RHOJ, RHOC, RF, RPHIO, DTZ, SO, CK, DA, H, U
DATA TEXP(1), TEXP(2), TEXP(3) /10HEXPERIMENT, 10HAL DATA , 2H, >
1/
DATA (PEXX(I), I=1,9) /4.0, 4.5, 5.0, 5.5, 6.0, 6.5, 7.0, 7.5, 8.0/
DATA (PEXVJ(I), I=1,9) /7.01, 6.68, 6.35, 6.03, 5.62, 5.01, 4.17, 3.19, 2.25/
DATA (PEXVO(I), I=1,9) /2.050, 2.069, 2.065, 2.046, 2.013, 1.956, 1.870, 1.724, 1.509/
DATA (PEXPX(I), I=1,3) /5.68, 6.94, 8.21/
DATA (PEXV(I), I=1,3) /2.1, 1.78, 1.42/
M=10
ICOUNT=ICOUNT+1
IF (ICOUNT.LE.1) GO TO 1
READ (5,5) M, (PEXX(I), PEXVJ(I), I=1, M)
M=M+1
PEXX(M)=.875
PEXVJ(M)=9.13
GO TO 4
1 CONTINUE
PEXX(10)=.875
PEXVJ(10)=9.13
PEXPX(4)=.910
PEXV(4)=9.13
FAC=.5*.7506
DO 2 I=1,9
PEXX(I)=PEXX(I)+FAC
PEXV(I)=PEXV(I)/H
2 CONTINUE
DO 3 I=1,3
PEXPX(I)=PEXPX(I)*FAC
PEXPX(I)=PEXPX(I)/H
3 CONTINUE
CALL PLTCCD (2,5,PEXX(1),PEXVO(1),9)
CALL PLTCCD (2,4,PEXPX(1),PEXV(1),4)
4 CONTINUE
YY=9.
XX=.6
CALL PLTCCT (.1,TEXP(1),0.,1.,XX,YY)
YY=8.
C CALL PLTCCT (.1,BASC(1),0.,1.,XX,YY)
CALL PLTCCD (2,5,PEXX(1),PEXVJ(1),M)
RETURN
C
5 FORMAT (I5/(HE10,1))
END

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*1002* 8
PLTEP 2
PLTEP 3
PLTEP 4
PLTEP 5
PLTEP 6D
PLTEP 7U
PLTEP 8D
PLTEP 9D
PLTEP 10U
PLTEP 11D
PLTEP 12U
PLTEP 13D
PLTEP 14U
PLTEP 15
PLTEP 16
PLTEP 17
PLTEP 18M
PLTEP 19
PLTEP 20
PLTEP 21
PLTEP 22
PLTEP 23
PLTEP 24
PLTEP 25
PLTEP 26
PLTEP 27
PLTEP 28
PLTEP 29
PLTEP 30
PLTEP 31
PLTEP 32
PLTEP 33
PLTEP 34
PLTEP 35
PLTEP 36
PLTEP 37
PLTEP 38
PLTEP 39
PLTEP 40
PLTEP 41
PLTEP 42
PLTEP 43
PLTEP 44
PLTEP 45
PLTEP 46
PLTEP 47
PLTEP 48
PLTEP 49-

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# SAMPLE INPUT

CARD NO. 1

HEADING CARD

1 105-MM SHAPED CHARGE SAMPLE CASE

CARD NO. 2

			LINER PARAMETERS		CONF TKS	RHO CONF.	NPLT	NPOS
ALPHA	FPS	RHOJ	HF	H				
21.	.269	8.9	4.3185	0.	0.	0.	0	0

CARD NO. 3

			TARGET PARAMETERS		DPOINT	WDPT	NEXP	N
RHOC	SO	CK	BHN	T1				
7.8	0.	0.	300.	112.	3.9899	0.	1	100

# 105-MM SHAPED CHARGE SAMPLE CASE

## INPUT PARAMETERS

ALPHA = 21.000 K CONSTANT = .180 THICKNESS OF LINER = .269 DENSITY = 8.900  
 UNCONFINED THICKNESS = 0.000 DENSITY = 0.000 FACTOR = 1.0000  
 LINER HEIGHT = 11.2561 LINER RADIUS = 4.3185  
 INCREMENT OF Z (DZ) = .1125 INCREMENT OF TIME (DTZ) = .1410 NUMBER OF ELEMENTS (N) = 100

## DETONATION

EXPLOSIVE = CORD 2 DETONATION VELOCITY(U) = .79000 EXPLOSIVE DENSITY = 1.720  
 POINT OF DETONATION FROM CONE APEX(DPOINT) = 3.9000 DET. RADIUS (RDPT) = 0.0000

## TARGET INPUTS

TARGET DENSITY = 7.000 STAND-OFF DISTANCE = 9.000 CK CONSTANT(BHN) = 300.000000

## MISC. INPUTS

SKIP PLOTS

# INITIAL POSITION

I	1	Z	0.0000	I	2	Z	.1125	I	3	Z	.2250	I	4	Z	.3375	I	5	Z	.4500	I	6	Z	.5625	I	7	Z	.6750
R	1	R	.7875	R	9	R	.9000	R	10	R	1.0125	R	11	R	1.1250	R	12	R	1.2375	R	13	R	1.3500	R	14	R	1.4625
15	15	15	1.5750	16	16	16	1.6875	17	17	17	1.8000	18	18	18	1.9125	19	19	19	2.0250	20	20	2.1375	21	21	21	2.2500	
22	22	22	2.3625	23	23	23	2.4750	24	24	24	2.5875	25	25	25	2.7000	26	26	26	2.8125	27	27	2.9250	28	28	28	3.0375	
29	29	29	3.1500	30	30	30	3.2625	31	31	31	3.3750	32	32	32	3.4875	33	33	33	3.6000	34	34	3.7125	35	35	35	3.8250	
36	36	36	3.9375	37	37	37	4.0500	38	38	38	4.1625	39	39	39	4.2750	40	40	40	4.3875	41	41	4.5000	42	42	42	4.6125	
43	43	43	4.7250	44	44	44	4.8375	45	45	45	4.9500	46	46	46	5.0625	47	47	47	5.1750	48	48	5.2875	49	49	49	5.4000	
50	50	50	5.5125	51	51	51	5.6250	52	52	52	5.7375	53	53	53	5.8500	54	54	54	5.9625	55	55	6.0750	56	56	56	6.1875	
57	57	57	6.3000	58	58	58	6.4125	59	59	59	6.5250	60	60	60	6.6375	61	61	61	6.7500	62	62	6.8625	63	63	63	6.9750	
64	64	64	7.0875	65	65	65	7.2000	66	66	66	7.3125	67	67	67	7.4250	68	68	68	7.5375	69	69	7.6500	70	70	70	7.7625	
71	71	71	7.8750	72	72	72	7.9875	73	73	73	8.1000	74	74	74	8.2125	75	75	75	8.3250	76	76	8.4375	77	77	77	8.5500	
78	78	78	8.6625	79	79	79	8.7750	80	80	80	8.8875	81	81	81	9.0000	82	82	82	9.1125	83	83	9.2250	84	84	84	9.3375	
85	85	85	9.4500	86	86	86	9.5625	87	87	87	9.6750	88	88	88	9.7875	89	89	89	9.9000	90	90	10.0125	91	91	91	10.1250	
92	92	92	10.2375	93	93	93	10.3500	94	94	94	10.4625	95	95	95	10.5750	96	96	96	10.6875	97	97	10.8000	98	98	98	10.9125	
99	99	99	11.0250	100	100	100	11.1375																				

I	1	R	0.0000	I	2	P	.0432	I	3	R	.0864	I	4	R	.1296	I	5	R	.1727	I	6	R	.2159	I	7	R	.2591
R	1	R	.3023	R	9	R	.3455	R	10	R	.3887	R	11	R	.4319	R	12	R	.4750	R	13	R	.5182	R	14	R	.5614
15	15	15	.6046	16	16	16	.6478	17	17	17	.6910	18	18	18	.7341	19	19	19	.7773	20	20	.8205	21	21	21	.8637	
22	22	22	.9069	23	23	23	.9501	24	24	24	.9933	25	25	25	1.0364	26	26	26	1.0796	27	27	1.1228	28	28	28	1.1660	
29	29	29	1.2092	30	30	30	1.2524	31	31	31	1.2956	32	32	32	1.3387	33	33	33	1.3819	34	34	1.4251	35	35	35	1.4683	
36	36	36	1.5115	37	37	37	1.5547	38	38	38	1.5978	39	39	39	1.6410	40	40	40	1.6842	41	41	1.7274	42	42	42	1.7706	
43	43	43	1.6138	44	44	44	1.6570	45	45	45	1.7001	46	46	46	1.7433	47	47	47	1.7865	48	48	1.8297	49	49	49	1.8729	
50	50	50	2.0161	51	51	51	2.0593	52	52	52	2.1024	53	53	53	2.1456	54	54	54	2.1888	55	55	2.2320	56	56	56	2.2752	
57	57	57	2.4184	58	58	58	2.4615	59	59	59	2.5047	60	60	60	2.5479	61	61	61	2.5911	62	62	2.6343	63	63	63	2.6775	
64	64	64	2.7070	65	65	65	2.7502	66	66	66	2.7934	67	67	67	2.8366	68	68	68	2.8798	69	69	2.9230	70	70	70	2.9662	
71	71	71	3.0230	72	72	72	3.0661	73	73	73	3.1093	74	74	74	3.1525	75	75	75	3.1957	76	76	3.2389	77	77	77	3.2821	
78	78	78	3.2525	79	79	79	3.2957	80	80	80	3.3389	81	81	81	3.3821	82	82	82	3.4253	83	83	3.4685	84	84	84	3.5117	
85	85	85	3.6275	86	86	86	3.6707	87	87	87	3.7139	88	88	88	3.7571	89	89	89	3.8003	90	90	3.8435	91	91	91	3.8867	
92	92	92	3.9298	93	93	93	3.9730	94	94	94	4.0162	95	95	95	4.0594	96	96	96	4.1026	97	97	4.1458	98	98	98	4.1890	
99	99	99	4.2321	100	100	100	4.2753																				

I	1	PI	0.0000	I	2	PI	.0513	I	3	PI	.1026	I	4	PI	.1539	I	5	PI	.2052	I	6	PI	.2565	I	7	PI	.3078
R	1	R	.0142	R	9	R	.0573	R	10	R	.1085	R	11	R	.1598	R	12	R	.2111	R	13	R	.2624	R	14	R	.3137
15	15	15	.3165	16	16	16	.3596	17	17	17	.4028	18	18	18	.4459	19	19	19	.4892	20	20	.5324	21	21	21	.5756	
22	22	22	.6187	23	23	23	.6619	24	24	24	.7051	25	25	25	.7483	26	26	26	.7915	27	27	.8347	28	28	28	.8779	
29	29	29	.9210	30	30	30	.9642	31	31	31	1.0074	32	32	32	1.0506	33	33	33	1.0938	34	34	1.1370	35	35	35	1.1802	
36	36	36	1.2233	37	37	37	1.2665	38	38	38	1.3097	39	39	39	1.3529	40	40	40	1.3961	41	41	1.4393	42	42	42	1.4824	
43	43	43	1.5256	44	44	44	1.5688	45	45	45	1.6120	46	46	46	1.6552	47	47	47	1.6984	48	48	1.7416	49	49	49	1.7847	
50	50	50	1.6219	51	51	51	1.6651	52	52	52	1.7083	53	53	53	1.7515	54	54	54	1.7947	55	55	1.8379	56	56	56	1.8811	
57	57	57	2.1302	58	58	58	2.1734	59	59	59	2.2166	60	60	60	2.2598	61	61	61	2.3030	62	62	2.3461	63	63	63	2.3893	
64	64	64	2.4325	65	65	65	2.4757	66	66	66	2.5189	67	67	67	2.5621	68	68	68	2.6053	69	69	2.6484	70	70	70	2.6916	
71	71	71	2.7348	72	72	72	2.7780	73	73	73	2.8212	74	74	74	2.8644	75	75	75	2.9076	76	76	2.9507	77	77	77	2.9939	
78	78	78	3.0371	79	79	79	3.0803	80	80	80	3.1235	81	81	81	3.1667	82	82	82	3.2099	83	83	3.2530	84	84	84	3.2962	
85	85	85	3.3394	86	86	86	3.3826	87	87	87	3.4258	88	88	88	3.4690	89	89	89	3.5121	90	90	3.5553	91	91	91	3.5985	
92	92	92	3.6417	93	93	93	3.6849	94	94	94	3.7281	95	95	95	3.7713	96	96	96	3.8144	97	97	3.8576	98	98	98	3.9008	
99	99	99	3.9440	100	100	100	3.9872																				

I	1	PI	0.0000	I	2	PI	.0142	I	3	PI	.0284	I	4	PI	.0426	I	5	PI	.0568	I	6	PI	.0710	I	7	PI	.0852
R	1	R	.3165	R	9	R	.3596	R	10	R	.4028	R	11	R	.4459	R	12	R	.4892	R	13	R	.5324	R	14	R	.5756
15	15	15	.6187	16	16	16	.6619	17	17	17	.7051	18	18	18	.7483	19	19	19	.7915	20	20	.8347	21	21	21	.8779	
22	22	22	.9210	23	23	23	.9642	24	24	24	1.0074	25	25	25	1.0506	26	26	26	1.0938	27	27	1.1370	28	28	28	1.1802	
29	29	29	1.2233	30	30	30	1.2665	31	31	31	1.3097	32	32	32	1.3529	33	33	33	1.3961	34	34	1.4393	35	35	35	1.4824	
36	36	36	1.5256	37	37	37	1.5688	38	38	38	1.6120	39	39	39	1.6552	40	40	40	1.6984	41	41	1.7416	42	42	42	1.7847	
43	43	43	1.8219	44	44	44	1.8651	45	45	45	1.9083	46	46	46	1.9515	47	47	47	1.9947	48	48	2.0379	49	49	49	2.0811	
50	50	50	2.1302	51	51	51	2.1734	52	52	52	2.2166	53	53	53	2.2598	54	54	54	2.3030	55	55	2.3461	56	56	56	2.3893	
57	57	57	2.4325	58	58	58	2.4757	59	59	59	2.5189	60	60	60	2.5621	61	61	61	2.6053	62	62	2.6484	63	63	63	2.6916	
64	64	64	2.7348	65	65	65	2.7780	66	66	66	2.8212	67	67	67	2.8644	68	68	68	2.9076	69	69	2.9507	70	70	70	2.9939	
71	71	71	3.0371	72	72	72	3.0803	73	73	73	3.1235	74	74	74	3.1667	75	75	75	3.2099	76	76	3.2530	77	77	77	3.2962	
78	78	78	3.3394	79	79	79	3.3826	80	80	80	3.4258	81	81	81	3.4690	82	82	82	3.5121	83	83	3.5553	84	84	84	3.5985	
85	85	85	3.6417	86	86	86	3.6849	87	87	87	3.7281	88	88	88	3.7713	89	89	89	3.8144	90	90	3.8576	91	91	91	3.9008	
92	92	92	3.9440	93	93	93	3.9872	94	94	94	4.0304	95	95	95	4.0736												

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7	TZ	EPS	E	PHI	NETA	DPHI	MPHI	V	H	TAU	C
5.700E-01	1.353E+01	2.714E-01	1.908E+00	1.572E+01	4.812E+01	1.355E-01	3.645E+00	2.202E-01	2.462E+00	1.276E+01	7.769E+00
5.800E-01	1.368E+01	2.714E-01	1.908E+00	1.572E+01	4.812E+01	1.355E-01	3.645E+00	2.202E-01	2.462E+00	1.276E+01	7.769E+00
5.900E-01	1.383E+01	2.713E-01	1.821E+00	1.543E+01	4.937E+01	1.467E-01	3.713E+00	2.141E-01	2.548E+00	1.344E+01	8.033E+00
6.000E-01	1.398E+01	2.713E-01	1.777E+00	1.524E+01	5.005E+01	1.524E-01	3.750E+00	2.140E-01	2.591E+00	1.380E+01	8.165E+00
6.100E-01	1.414E+01	2.712E-01	1.733E+00	1.512E+01	5.077E+01	1.584E-01	3.789E+00	2.117E-01	2.634E+00	1.417E+01	8.296E+00
6.200E-01	1.429E+01	2.712E-01	1.689E+00	1.496E+01	5.153E+01	1.645E-01	3.831E+00	2.094E-01	2.677E+00	1.455E+01	8.427E+00
6.300E-01	1.444E+01	2.711E-01	1.645E+00	1.478E+01	5.234E+01	1.707E-01	3.875E+00	2.070E-01	2.721E+00	1.494E+01	8.558E+00
6.400E-01	1.459E+01	2.711E-01	1.601E+00	1.461E+01	5.319E+01	1.772E-01	3.922E+00	2.045E-01	2.764E+00	1.535E+01	8.688E+00
6.500E-01	1.474E+01	2.711E-01	1.557E+00	1.442E+01	5.409E+01	1.839E-01	3.972E+00	2.019E-01	2.807E+00	1.578E+01	8.818E+00
6.600E-01	1.489E+01	2.710E-01	1.513E+00	1.423E+01	5.505E+01	1.908E-01	4.026E+00	1.992E-01	2.850E+00	1.622E+01	8.948E+00
6.700E-01	1.504E+01	2.710E-01	1.469E+00	1.404E+01	5.604E+01	1.974E-01	4.082E+00	1.964E-01	2.893E+00	1.669E+01	9.077E+00
6.800E-01	1.520E+01	2.709E-01	1.425E+00	1.383E+01	5.711E+01	2.050E-01	4.143E+00	1.935E-01	2.937E+00	1.717E+01	9.206E+00
6.900E-01	1.535E+01	2.709E-01	1.381E+00	1.362E+01	5.825E+01	2.127E-01	4.207E+00	1.905E-01	2.980E+00	1.768E+01	9.334E+00
7.000E-01	1.550E+01	2.709E-01	1.337E+00	1.340E+01	5.946E+01	2.206E-01	4.277E+00	1.874E-01	3.023E+00	1.821E+01	9.462E+00
7.100E-01	1.565E+01	2.708E-01	1.293E+00	1.317E+01	6.074E+01	2.289E-01	4.351E+00	1.842E-01	3.066E+00	1.878E+01	9.592E+00
7.200E-01	1.580E+01	2.708E-01	1.249E+00	1.293E+01	6.210E+01	2.376E-01	4.431E+00	1.809E-01	3.109E+00	1.937E+01	9.716E+00
7.300E-01	1.595E+01	2.707E-01	1.204E+00	1.268E+01	6.353E+01	2.466E-01	4.517E+00	1.774E-01	3.153E+00	2.000E+01	9.843E+00
7.400E-01	1.610E+01	2.707E-01	1.160E+00	1.243E+01	6.503E+01	2.559E-01	4.610E+00	1.739E-01	3.196E+00	2.067E+01	9.968E+00
7.500E-01	1.626E+01	2.707E-01	1.116E+00	1.216E+01	6.662E+01	2.657E-01	4.711E+00	1.701E-01	3.239E+00	2.138E+01	1.009E+01
7.600E-01	1.641E+01	2.707E-01	1.071E+00	1.189E+01	6.829E+01	2.760E-01	4.820E+00	1.663E-01	3.282E+00	2.214E+01	1.022E+01
7.700E-01	1.656E+01	2.706E-01	1.027E+00	1.160E+01	7.003E+01	2.867E-01	4.939E+00	1.623E-01	3.325E+00	2.296E+01	1.034E+01
7.800E-01	1.671E+01	2.706E-01	9.826E-01	1.130E+01	7.186E+01	2.980E-01	5.070E+00	1.581E-01	3.368E+00	2.384E+01	1.047E+01
7.900E-01	1.686E+01	2.706E-01	9.382E-01	1.099E+01	7.376E+01	3.098E-01	5.213E+00	1.537E-01	3.412E+00	2.479E+01	1.059E+01
8.000E-01	1.701E+01	2.705E-01	8.937E-01	1.067E+01	7.572E+01	3.222E-01	5.370E+00	1.492E-01	3.455E+00	2.583E+01	1.071E+01
8.100E-01	1.716E+01	2.705E-01	8.493E-01	1.033E+01	7.775E+01	3.353E-01	5.544E+00	1.445E-01	3.498E+00	2.696E+01	1.083E+01
8.200E-01	1.731E+01	2.705E-01	8.048E-01	9.985E-01	7.984E+01	3.491E-01	5.738E+00	1.397E-01	3.541E+00	2.821E+01	1.095E+01
8.300E-01	1.747E+01	2.705E-01	7.602E-01	9.622E-01	8.197E+01	3.637E-01	5.955E+00	1.346E-01	3.584E+00	2.959E+01	1.107E+01
8.400E-01	1.762E+01	2.704E-01	7.157E-01	9.242E-01	8.413E+01	3.791E-01	6.199E+00	1.293E-01	3.627E+00	3.112E+01	1.119E+01
8.500E-01	1.777E+01	2.704E-01	6.711E-01	8.847E+00	8.631E+01	3.953E-01	6.476E+00	1.237E-01	3.671E+00	3.285E+01	1.131E+01
8.600E-01	1.792E+01	2.704E-01	6.265E-01	8.435E+00	8.849E+01	4.126E-01	6.793E+00	1.180E-01	3.714E+00	3.480E+01	1.142E+01
8.700E-01	1.807E+01	2.704E-01	5.819E-01	8.004E+00	9.066E+01	4.310E-01	7.159E+00	1.119E-01	3.757E+00	3.703E+01	1.154E+01
8.800E-01	1.822E+01	2.703E-01	5.373E-01	7.553E+00	9.280E+01	4.505E-01	7.586E+00	1.056E-01	3.800E+00	3.942E+01	1.165E+01
8.900E-01	1.837E+01	2.703E-01	4.924E-01	7.082E+00	9.489E+01	4.712E-01	8.090E+00	9.905E-02	3.843E+00	4.266E+01	1.177E+01
9.000E-01	1.853E+01	2.703E-01	4.479E-01	6.548E+00	9.693E+01	4.934E-01	8.696E+00	9.215E-02	3.887E+00	4.627E+01	1.188E+01
9.100E-01	1.868E+01	2.703E-01	4.032E-01	6.071E+00	9.888E+01	5.171E-01	9.437E+00	8.492E-02	3.930E+00	5.067E+01	1.199E+01
9.200E-01	1.883E+01	2.702E-01	3.585E-01	5.529E+00	1.007E+02	5.425E-01	1.036E+01	7.733E-02	3.973E+00	5.614E+01	1.210E+01
9.300E-01	1.898E+01	2.702E-01	3.138E-01	4.959E+00	1.025E+02	5.698E-01	1.155E+01	6.936E-02	4.016E+00	6.313E+01	1.221E+01
9.400E-01	1.913E+01	2.702E-01	2.690E-01	4.360E+00	1.041E+02	5.992E-01	1.314E+01	6.098E-02	4.059E+00	7.241E+01	1.231E+01
9.500E-01	1.928E+01	2.702E-01	2.242E-01	3.729E+00	1.056E+02	6.308E-01	1.536E+01	5.216E-02	4.103E+00	8.536E+01	1.242E+01
9.600E-01	1.943E+01	2.702E-01	1.794E-01	3.064E+00	1.070E+02	6.659E-01	1.870E+01	4.286E-02	4.146E+00	1.047E+02	1.252E+01
9.700E-01	1.958E+01	2.702E-01	1.346E-01	2.362E+00	1.082E+02	7.020E-01	2.426E+01	3.304E-02	4.189E+00	1.369E+02	1.262E+01
9.800E-01	1.973E+01	2.701E-01	8.974E-02	1.620E+00	1.093E+02	7.422E-01	3.537E+01	2.266E-02	4.232E+00	2.012E+02	1.272E+01
9.900E-01	1.988E+01	2.701E-01	4.484E-02	8.338E-01	1.102E+02	7.840E-01	6.871E+01	1.166E-02	4.275E+00	3.038E+02	1.281E+01



7	VJ	VN	KV	DMJ	DMN	DEL	DPL	DMJK	TMJ
5.700F-01	5.392E-01	2.017E-02	2.590E-01	1.662E-01	6.338E-01	1.042E-01	4.462E+00	7.417E-01	1.640E+01
5.800F-01	5.273E-01	1.847E-02	2.544E-01	1.702E-01	6.298E-01	1.041E-01	4.540E+00	7.728E-01	1.758E+01
5.900F-01	5.162E-01	1.671E-02	2.497E-01	1.744E-01	6.256E-01	1.079E-01	4.619E+00	8.057E-01	1.836E+01
6.000F-01	5.048F-01	1.488E-02	2.449E-01	1.790E-01	6.210E-01	1.075E-01	4.697E+00	8.406E-01	1.922E+01
6.100E-01	4.931E-01	1.298E-02	2.401E-01	1.838E-01	6.162E-01	1.070E-01	4.775E+00	8.777E-01	2.010E+01
6.200E-01	4.812E-01	1.100E-02	2.351E-01	1.890E-01	6.110E-01	1.064E-01	4.853E+00	9.171E-01	2.102E+01
6.300F-01	4.690E-01	8.950E-03	2.300E-01	1.945E-01	6.055E-01	1.056E-01	4.932E+00	9.592E-01	2.198E+01
6.400E-01	4.568E-01	6.821E-03	2.249E-01	2.004E-01	5.996E-01	1.047E-01	5.010E+00	1.004E+00	2.298E+01
6.500E-01	4.439E-01	4.612E-03	2.196E-01	2.068E-01	5.932E-01	1.037E-01	5.088E+00	1.052E+00	2.403E+01
6.600F-01	4.310F-01	2.320E-03	2.144E-01	2.136E-01	5.864E-01	1.025E-01	5.167E+00	1.103E+00	2.514E+01
6.700F-01	4.181E-01	9.977E-06	2.091E-01	2.207E-01	5.793E-01	1.012E-01	5.245E+00	1.158E+00	2.629E+01
6.800F-01	4.049E-01	7.444E-03	2.037E-01	2.285E-01	5.715E-01	9.969E-02	5.323E+00	1.216E+00	2.751E+01
6.900F-01	3.914E-01	5.014E-03	1.982E-01	2.369E-01	5.631E-01	9.806E-02	5.401E+00	1.280E+00	2.879E+01
7.000F-01	3.777E-01	2.656E-03	1.927E-01	2.454E-01	5.541E-01	9.627E-02	5.480E+00	1.348E+00	3.014E+01
7.100F-01	3.640E-01	1.034E-02	1.872F-01	2.536E-01	5.444E-01	9.433E-02	5.558E+00	1.421E+00	3.156E+01
7.200F-01	3.501E-01	1.319E-02	1.816E-01	2.620E-01	5.340E-01	9.223E-02	5.636E+00	1.499E+00	3.306E+01
7.300F-01	3.361E-01	1.609E-02	1.761E-01	2.711E-01	5.229E-01	9.097E-02	5.715E+00	1.583E+00	3.464E+01
7.400F-01	3.220E-01	1.905E-02	1.706E-01	2.800E-01	5.110E-01	8.755E-02	5.793E+00	1.674E+00	3.631E+01
7.500F-01	3.080E-01	2.209E-02	1.650F-01	3.016E-01	4.984E-01	8.497E-02	5.871E+00	1.771E+00	3.809E+01
7.600F-01	2.939E-01	2.518E-02	1.595E-01	3.150E-01	4.850E-01	8.224E-02	5.949E+00	1.874E+00	3.996E+01
7.700E-01	2.799E-01	2.831E-02	1.541E-01	3.293E-01	4.707E-01	7.934E-02	6.028E+00	1.985E+00	4.194E+01
7.800F-01	2.659E-01	3.147E-02	1.487E-01	3.433E-01	4.557E-01	7.629E-02	6.106E+00	2.102E+00	4.405E+01
7.900F-01	2.520E-01	3.444E-02	1.433E-01	3.576E-01	4.403E-01	7.309E-02	6.184E+00	2.227E+00	4.627E+01
8.000F-01	2.382E-01	3.777E-02	1.380E-01	3.767E-01	4.233E-01	6.974E-02	6.263E+00	2.359E+00	4.863E+01
8.100F-01	2.246E-01	4.085E-02	1.328E-01	3.939E-01	4.061E-01	6.624E-02	6.341E+00	2.498E+00	5.113E+01
8.200E-01	2.112E-01	4.363E-02	1.275E-01	4.114E-01	3.882E-01	6.240E-02	6.419E+00	2.643E+00	5.377E+01
8.300E-01	1.980E-01	4.646E-02	1.223E-01	4.301E-01	3.699E-01	5.843E-02	6.497E+00	2.795E+00	5.657E+01
8.400F-01	1.851E-01	4.929E-02	1.172E-01	4.489E-01	3.511E-01	5.494E-02	6.576E+00	2.952E+00	5.952E+01
8.500F-01	1.723E-01	5.166E-02	1.120E-01	4.678E-01	3.322E-01	5.094E-02	6.654E+00	3.113E+00	6.263E+01
8.600F-01	1.598E-01	5.369E-02	1.068E-01	4.868E-01	3.132E-01	4.685E-02	6.732E+00	3.277E+00	6.591E+01
8.700F-01	1.476E-01	5.532E-02	1.015E-01	5.054E-01	2.942E-01	4.267E-02	6.810E+00	3.444E+00	6.935E+01
8.800F-01	1.356E-01	5.645E-02	9.603E-02	5.244E-01	2.756E-01	3.844E-02	6.889E+00	3.613E+00	7.297E+01
8.900F-01	1.239E-01	5.701E-02	9.044E-02	5.427E-01	2.573E-01	3.418E-02	6.967E+00	3.781E+00	7.675E+01
9.000F-01	1.125E-01	5.690E-02	8.461E-02	5.603E-01	2.397E-01	2.991E-02	7.045E+00	3.947E+00	8.070E+01
9.100F-01	1.010E-01	5.602E-02	7.850E-02	5.772E-01	2.228E-01	2.569E-02	7.124E+00	4.112E+00	8.491E+01
9.200F-01	8.978E-02	5.430E-02	7.204E-02	5.932E-01	2.068E-01	2.153E-02	7.202E+00	4.272E+00	8.904E+01
9.300F-01	7.870E-02	5.142E-02	6.516F-02	6.082E-01	1.918E-01	1.751E-02	7.280E+00	4.424E+00	9.351E+01
9.400F-01	6.770F-02	4.792E-02	5.781E-02	6.221E-01	1.779E-01	1.368E-02	7.358E+00	4.577E+00	9.804E+01
9.500E-01	5.672E-02	4.310E-02	4.991E-02	6.344E-01	1.652E-01	1.012E-02	7.437E+00	4.721E+00	1.028E+02
9.600F-01	4.571E-02	3.709E-02	4.140E-02	6.463E-01	1.537E-01	6.901E-03	7.515E+00	4.857E+00	1.077E+02
9.700F-01	3.440E-02	2.682E-02	3.221F-02	6.565E-01	1.435F-01	4.144E-02	7.593E+00	4.985E+00	1.126E+02
9.800F-01	2.333F-02	1.724F-02	2.229E-02	6.654E-01	1.346E-01	1.949E-03	7.672E+00	5.105E+00	1.178E+02
9.900F-01	1.143E-02	1.132E-02	1.157F-02	6.730E-01	1.270E-01	5.270E-04	7.750E+00	5.215E+00	1.230E+02



# 105-MM SHAPED CHARGE SAMPLE CASE

## INITIAL CONDITIONS FOR PENETRATION

U1 =	112.0000	U0 =	53.0030	Z0 =	37.1411
U2 =	36.0439	VJT =	.0316	VJ0 =	.7011
APAY =	1.042148	UMIN =	.10588	CK =	.03510

37343 (20) ILVGL 47.4.0

MLT	TIME	F	DF	T	UT	G	H-SO	A	DELA	UV	U/UT
1	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
2	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
3	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
4	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
5	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
6	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
7	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
8	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
9	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
10	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
11	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
12	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
13	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
14	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
15	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
16	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
17	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
18	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
19	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
20	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
21	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
22	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
23	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
24	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
25	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
26	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
27	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
28	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
29	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
30	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
31	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
32	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
33	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
34	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
35	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
36	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
37	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
38	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
39	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
40	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
41	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
42	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
43	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
44	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
45	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
46	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
47	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
48	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
49	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
50	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
51	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
52	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
53	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
54	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
55	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
56	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
57	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
58	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
59	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
60	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
61	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
62	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
63	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
64	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
65	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
66	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
67	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
68	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
69	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
70	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
71	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
72	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
73	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
74	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
75	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
76	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
77	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
78	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
79	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
80	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
81	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
82	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
83	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
84	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
85	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
86	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
87	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
88	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
89	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
90	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
91	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
92	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
93	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
94	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
95	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
96	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
97	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
98	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
99	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.
100	00	0.	0.	0.	0.	0.	0.	0.	0.	0.	0.

54	1.254E+00	2.035E+01	5.629E+00	-1.108E-01	1.031E+02	5.014E+01	5.341E+01	1.559E+01	7.246E+01	4.056E+02	-1.011E-02	-6.703E+00
55	1.277E+00	2.028E+01	5.604E+00	-1.246E-01	1.049E+02	5.304E+01	5.404E+01	1.546E+01	7.246E+01	4.056E+02	-1.011E-02	-6.703E+00
56	1.303E+00	2.128E+01	5.265E+00	-1.303E-01	1.104E+02	5.702E+01	5.601E+01	1.532E+01	8.997E+01	6.994E+02	-1.063E-02	-7.622E+00
57	1.330E+00	2.174E+01	5.111E+00	-1.511E-01	1.133E+02	6.224E+01	5.722E+01	1.542E+01	9.493E+01	7.493E+02	-1.089E-02	-7.638E+00
58	1.354E+00	2.224E+01	4.942E+00	-1.604E-01	1.200E+02	6.703E+01	5.849E+01	1.561E+01	1.095E+00	1.001E+01	-1.116E-02	-7.970E+00
59	1.384E+00	2.274E+01	4.756E+00	-1.653E-01	1.252E+02	7.234E+01	5.985E+01	1.531E+01	1.200E+00	1.055E+01	-1.142E-02	-8.009E+00
60	1.422E+00	2.330E+01	4.555E+00	-2.015E-01	1.304E+02	7.779E+01	6.130E+01	1.603E+01	1.311E+00	1.110E+01	-1.168E-02	-8.224E+00
61	1.455E+00	2.384E+01	4.337E+00	-2.181E-01	1.349E+02	8.303E+01	6.285E+01	1.626E+01	1.424E+00	1.168E+01	-1.193E-02	-8.345E+00
62	1.494E+00	2.444E+01	4.102E+00	-2.330E-01	1.373E+02	8.866E+01	6.450E+01	1.652E+01	1.551E+00	1.227E+01	-1.218E-02	-8.459E+00
63	1.534E+00	2.504E+01	3.850E+00	-2.522E-01	1.511E+02	9.606E+01	6.627E+01	1.679E+01	1.679E+00	1.287E+01	-1.243E-02	-8.567E+00
64	1.579E+00	2.552E+01	3.590E+00	-2.607E-01	1.592E+02	1.062E+02	6.817E+01	1.709E+01	1.814E+00	1.349E+01	-1.265E-02	-8.666E+00
65	1.627E+00	2.612E+01	3.293E+00	-2.647E-01	1.683E+02	1.153E+02	7.021E+01	1.740E+01	1.955E+00	1.412E+01	-1.289E-02	-8.757E+00
66	1.677E+00	2.673E+01	2.994E+00	-2.947E-01	1.782E+02	1.251E+02	7.239E+01	1.785E+01	2.101E+00	1.449E+01	-1.292E-02	-8.735E+00
67	1.727E+00	2.737E+01	2.672E+00	-3.215E-01	1.893E+02	1.363E+02	7.475E+01	1.813E+01	2.254E+00	1.536E+01	-1.326E-02	-8.848E+00
68	1.794E+00	2.803E+01	2.332E+00	-3.407E-01	2.017E+02	1.487E+02	7.732E+01	1.851E+01	2.415E+00	1.671E+01	-1.368E-02	-9.011E+00
69	1.864E+00	2.871E+01	1.974E+00	-3.591E-01	2.157E+02	1.627E+02	8.010E+01	1.895E+01	2.582E+00	1.735E+01	-1.378E-02	-9.037E+00
70	1.924E+00	2.943E+01	1.504E+00	-3.751E-01	2.315E+02	1.784E+02	8.312E+01	1.943E+01	2.755E+00	1.796E+01	-1.395E-02	-9.040E+00
71	2.004E+00	3.017E+01	1.207E+00	-3.916E-01	2.492E+02	1.962E+02	8.640E+01	1.994E+01	2.935E+00	1.856E+01	-1.398E-02	-9.040E+00
72	2.084E+00	3.095E+01	7.992E+01	-4.076E-01	2.694E+02	2.164E+02	8.998E+01	2.050E+01	3.121E+00	1.913E+01	-1.404E-02	-9.014E+00
73	2.174E+00	3.177E+01	7.755E+01	-4.227E-01	2.923E+02	2.383E+02	9.389E+01	2.110E+01	3.312E+00	1.967E+01	-1.407E-02	-8.995E+00
74	2.274E+00	3.264E+01	7.606E+01	-4.370E-01	3.186E+02	2.656E+02	9.816E+01	2.176E+01	3.510E+00	2.017E+01	-1.407E-02	-8.995E+00
75	2.384E+00	3.354E+01	7.455E+01	-4.501E-01	3.487E+02	2.957E+02	1.024E+02	2.246E+01	3.710E+00	2.062E+01	-1.404E-02	-8.922E+00
76	2.504E+00	3.454E+01	7.277E+01	-4.621E-01	3.835E+02	3.305E+02	1.080E+02	2.322E+01	3.917E+00	2.101E+01	-1.398E-02	-8.719E+00
77	2.624E+00	3.564E+01	7.077E+01	-4.727E-01	4.238E+02	3.704E+02	1.137E+02	2.404E+01	4.127E+00	2.134E+01	-1.388E-02	-8.509E+00
78	2.744E+00	3.684E+01	6.847E+01	-4.814E-01	4.707E+02	4.177E+02	1.200E+02	2.491E+01	4.340E+00	2.160E+01	-1.375E-02	-8.261E+00
79	2.874E+00	3.814E+01	6.593E+01	-4.893E-01	5.257E+02	4.727E+02	1.269E+02	2.584E+01	4.556E+00	2.180E+01	-1.360E-02	-8.008E+00
80	3.014E+00	3.954E+01	6.313E+01	-4.962E-01	5.904E+02	5.374E+02	1.347E+02	2.684E+01	4.774E+00	2.199E+01	-1.341E-02	-7.742E+00
81	3.164E+00	4.094E+01	6.011E+01	-5.023E-01	6.671E+02	6.101E+02	1.433E+02	2.788E+01	4.993E+00	2.218E+01	-1.315E-02	-7.462E+00
82	3.324E+00	4.244E+01	5.693E+01	-5.077E-01	7.587E+02	7.057E+02	1.530E+02	2.898E+01	5.212E+00	2.237E+01	-1.280E-02	-7.126E+00
83	3.494E+00	4.394E+01	5.360E+01	-5.124E-01	8.690E+02	8.160E+02	1.639E+02	3.013E+01	5.430E+00	2.256E+01	-1.239E-02	-6.743E+00
84	3.674E+00	4.544E+01	4.917E+01	-5.161E-01	9.988E+02	9.408E+02	1.763E+02	3.131E+01	5.647E+00	2.275E+01	-1.193E-02	-6.307E+00
85	3.864E+00	4.694E+01	4.416E+01	-5.190E-01	1.147E+03	1.114E+03	1.903E+02	3.253E+01	5.862E+00	2.294E+01	-1.146E-02	-5.845E+00
86	4.064E+00	4.844E+01	3.895E+01	-5.211E-01	1.371E+03	1.318E+03	2.065E+02	3.375E+01	6.074E+00	2.313E+01	-1.095E-02	-5.369E+00
87	4.274E+00	5.004E+01	3.345E+01	-5.224E-01	1.624E+03	1.575E+03	2.253E+02	3.498E+01	6.281E+00	2.332E+01	-1.039E-02	-4.845E+00
88	4.494E+00	5.164E+01	2.774E+01	-5.230E-01	1.904E+03	1.906E+03	2.474E+02	3.618E+01	6.484E+00	2.351E+01	-1.075E-02	-4.297E+00
89	4.724E+00	5.324E+01	2.194E+01	-5.230E-01	2.202E+03	2.399E+03	2.734E+02	3.734E+01	6.682E+00	2.370E+01	-1.154E-02	-3.677E+00
90	4.964E+00	5.484E+01	1.614E+01	-5.230E-01	2.527E+03	3.058E+03	3.058E+02	3.843E+01	6.874E+00	2.389E+01	-1.135E-02	-3.055E+00
91	5.214E+00	5.644E+01	1.034E+01	-5.230E-01	2.874E+03	3.456E+03	3.456E+02	3.943E+01	7.060E+00	2.408E+01	-1.119E-02	-2.435E+00
92	5.474E+00	5.804E+01	4.44E+00	-5.230E-01	3.245E+03	3.732E+03	3.966E+02	4.032E+01	7.246E+00	2.427E+01	-1.108E-02	-1.815E+00
93	5.744E+00	5.964E+01	3.765E+00	-5.230E-01	3.641E+03	4.008E+03	4.644E+02	4.107E+01	7.431E+00	2.446E+01	-1.100E-02	-1.200E+00
94	6.024E+00	6.124E+01	3.085E+00	-5.230E-01	4.061E+03	4.690E+03	5.503E+02	4.167E+01	7.616E+00	2.465E+01	-1.098E-02	-6.497E+00
95	6.314E+00	6.284E+01	2.405E+00	-5.230E-01	4.581E+03	5.690E+03	6.601E+02	4.209E+01	7.801E+00	2.484E+01	-1.101E-02	-6.451E+00
96	6.614E+00	6.444E+01	1.725E+00	-5.230E-01	5.151E+03	6.890E+03	7.923E+02	4.232E+01	8.024E+00	2.503E+01	-1.111E-02	-6.449E+00
97	6.924E+00	6.604E+01	1.045E+00	-5.230E-01	5.770E+03	7.203E+03	9.427E+02	4.232E+01	8.202E+00	2.522E+01	-1.127E-02	-6.704E+00
98	7.244E+00	6.764E+01	4.65E+00	-5.230E-01	6.444E+03	8.009E+03	1.043E+03	4.232E+01	8.380E+00	2.541E+01	-1.151E-02	-6.933E+00
99	7.574E+00	6.924E+01	3.974E+00	-5.230E-01	7.177E+03	8.951E+03	1.195E+03	4.232E+01	8.557E+00	2.560E+01	-1.181E-02	-7.162E+00
100	7.914E+00	7.084E+01	3.304E+00	-5.230E-01	7.964E+03	9.951E+03	1.395E+03	4.232E+01	8.734E+00	2.579E+01	-1.216E-02	-7.393E+00

# 105-MM SHAPED CHARGE SAMPLE CASE

HOLE PROFILE

SD = 3.00 CD

PENTRATION STANDOFF

I	P(CM)	HC(CM)	PT(MM)	SO(CM)
1	0.00000	0.00000	311.70736	1.00000
2	0.00000	0.00000	354.96580	2.00000
3	0.00000	0.00000	380.93877	3.00000
4	0.00000	0.00000	395.21635	4.00000
5	0.00000	0.00000	400.88556	5.00000
6	0.00000	0.00000	399.86846	6.00000
7	0.00000	0.00000	393.45955	7.00000
8	0.00000	0.00000	382.67360	8.00000
9	0.00000	0.00000	370.83340	9.00000
10	0.00000	0.00000	359.28549	10.00000
11	0.00000	0.00000	348.00923	11.00000
12	0.00000	0.00000	336.98630	12.00000
13	0.00000	0.00000	326.20037	13.00000
14	0.00000	0.00000	315.63678	14.00000
15	0.00000	0.00000	305.28231	15.00000
16	0.00000	0.00000	295.12501	16.00000
17	0.00000	0.00000	285.15404	17.00000
18	0.00000	0.00000	275.35450	18.00000
19	0.00000	0.00000	265.73235	19.00000
20	0.00000	0.00000	256.26429	20.00000
21	0.00000	0.00000	246.94767	21.00000
22	0.00000	0.00000	237.77546	22.00000
23	0.00000	0.00000	228.74115	23.00000
24	0.00000	0.00000	219.83869	24.00000
25	0.00000	0.00000	211.06247	25.00000
26	0.00000	0.00000		
27	0.00000	0.00000		
28	0.00000	0.00000		
29	0.00000	0.00000		
30	0.00000	0.00000		
31	0.00000	0.00000		
32	0.00000	0.00000		
33	0.00000	0.00000		
34	0.00000	0.00000		
35	0.00000	0.00000		
36	0.00000	0.00000		
37	0.00000	0.00000		
38	4.08966	.15347		
39	4.34083	2.09919		
40	4.79126	1.46420		
41	5.25781	1.43460		
42	5.74158	1.40390		
43	6.24376	1.37219		
44	6.76563	1.33959		
45	7.30875	1.30593		
46	7.87454	1.27163		
47	8.46474	1.23655		
48	9.08121	1.20080		
49	9.72599	1.16446		
50	10.40131	1.12760		

I	P	PC
11	11.10963	1.09029
12	11.85363	1.05261
13	12.63628	1.01462
14	13.46083	.97641
15	14.33090	.93804
16	15.25045	.89958
17	16.22389	.86112
18	17.25609	.82272
19	18.34664	.80865
20	19.48506	.80245
21	20.67288	.79583
22	21.91220	.78874
23	23.20498	.78115
24	24.55304	.77303
25	25.95794	.76432
26	27.42096	.75498
27	28.94292	.74494
28	30.50486	.73845
29	32.14110	.72330
30	33.84023	.70978
31	35.59544	.69633
32	37.40375	.68182

# 105-MM SHAPED CHARGE SAMPLE CASE

## SUMMARY OF RESULTS

LINER MASS = 345.4492 GM      JET MASS = 122.9672 GM      SLUG MASS = 262.8820 GM  
 TOTAL KINETIC ENERGY = 6.3894 ERGS\*1.0E12  
 TOTAL JET KINETIC ENERGY = 6.2273 ERGS\*1.0E12  
 TOTAL JET KINETIC ENERGY ABOVE JET VELOCITY .25 CM/MSEC = 5.6405 ERGS\*1.0E12  
 TOTAL JET MASS ABOVE JET VELOCITY .25 CM/MSEC = 46.2734 GM

KINETIC ENERGY ABOVE .5 CM/MSEC =	3.8664	AND JET MASS =	19.2223
KINETIC ENERGY ABOVE .4 CM/MSEC =	4.4964	AND JET MASS =	27.5096
KINETIC ENERGY ABOVE .3 CM/MSEC =	5.3368	AND JET MASS =	38.0850
KINETIC ENERGY ABOVE .2 CM/MSEC =	5.8294	AND JET MASS =	53.7738
KINETIC ENERGY ABOVE .1 CM/MSEC =	6.1685	AND JET MASS =	84.4073

MAXIMUM RELATIVE MACH NUMBER = .8821

JET TIP AT I = 38 RELATIVE POSITION Z/H = .37000000

MEASURED JET TIP VELOCITY = 7.01100675 MM/MICROSEC

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